

4 FEB 09

SOLUZIONI

(I)

Le autofunzioni normalizzate della buca sono

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a} \quad \text{con} \quad E_n = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$$

quindi $\psi_1(x) = \sqrt{\frac{2}{a}} \sin \frac{\pi x}{a}$, $\psi_2 = \sqrt{\frac{2}{a}} \sin \frac{2\pi x}{a} = \sqrt{\frac{2}{a}} 2 \sin \frac{\pi x}{a} \cos \frac{\pi x}{a}$

$$\psi(x,0) = C \left(\psi_1(x) + \frac{\psi_2(x)}{2} \right) \xrightarrow{\text{normalizzato}} \psi(x,0) = \frac{2\psi_1(x) + \psi_2(x)}{\sqrt{5}}$$

$$\bullet \langle E \rangle = \sum E_n P(E_n) = \frac{1}{5} \cdot \frac{\pi^2 \hbar^2}{2ma^2} + \frac{1}{5} \frac{4\pi^2 \hbar^2}{2ma^2} = \frac{4\pi^2 \hbar^2}{5ma^2}$$

\uparrow $P(E_1)$ \uparrow $P(E_2)$

$$\bullet \psi(x,t) = \frac{2 e^{-\frac{i\pi^2 \hbar t}{2ma^2}} \psi_1(x) + e^{-\frac{4\pi^2 \hbar t}{2ma^2}} \psi_2(x)}{\sqrt{5}}$$

Lo stato finale è $\psi_f(x) = \frac{\psi_1(x) - 2\psi_2(x)}{\sqrt{5}}$

e la probabilità

$$P = |\langle \psi_f | \psi(t) \rangle|^2 = \left| \frac{2}{5} e^{-\frac{i\pi^2 \hbar t}{2ma^2}} - \frac{2}{5} e^{-\frac{4i\pi^2 \hbar t}{2ma^2}} \right|^2$$
$$= \frac{16}{25} \sin^2 \frac{3\pi^2 \hbar t}{4ma^2}$$

La probabilità che la particella sia in $a \leq x \leq 2a$ sarà zero se la buca è raddoppiata istantaneamente

$$\tilde{\psi}(x,0) = \begin{cases} \frac{2\psi_1(x) + \psi_2(x)}{\sqrt{5}} & 0 \leq x \leq a \\ 0 & a \leq x \leq 2a \end{cases}$$

Le autofunzioni del sistema finale sono $\psi_1^{FIN}(x) = \sqrt{\frac{2}{2a}} \sin \frac{n\pi x}{2a}$ e il primo stato eccitato è $\psi_2^{FIN}(x) = \sqrt{\frac{1}{a}} \sin \frac{\pi x}{a}$ che coincide con $\frac{1}{\sqrt{2}} \psi_1(x)$ in $0 \leq x \leq a$

$$P = |\langle \psi_2^{\text{FIP}} | \tilde{\psi} \rangle|^2 = \left| \int_0^{2a} \psi_2^{\text{FIP}*}(x) \tilde{\psi}(x) dx \right|^2 = \left| \int_0^a \psi_2^{\text{FIP}*}(x) \tilde{\psi}(x) dx \right|^2$$

$$= \left| \int_0^a \frac{1}{\sqrt{2}} \psi_1^*(x) \frac{2\psi_1(x) + \psi_2(x)}{\sqrt{5}} dx \right|^2 = \left| \frac{\sqrt{2}}{\sqrt{5}} \right|^2 = \frac{2}{5}$$

II

$H = \hbar I + \hbar \varepsilon \sigma_x$ che ha autovalori ed autovettori (base diagonalizzante σ_x) $E_{\pm} = \hbar \pm \hbar \varepsilon$ $|+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$ $|-\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$

• $\psi(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{|+\rangle + |-\rangle}{\sqrt{2}} \rightarrow \psi(t) = e^{-it} \frac{e^{-i\varepsilon t} |+\rangle + e^{i\varepsilon t} |-\rangle}{\sqrt{2}} = e^{-it} \begin{pmatrix} \cos \varepsilon t \\ -i \sin \varepsilon t \end{pmatrix}$

$P(|-\rangle) = \left| (0, 1) \cdot e^{-it} \begin{pmatrix} \cos \varepsilon t \\ -i \sin \varepsilon t \end{pmatrix} \right|^2 = \sin^2 \varepsilon t$

• $\langle C \rangle = \langle \psi(t) | C | \psi(t) \rangle = e^{it} (\cos \varepsilon t, i \sin \varepsilon t) i \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \cos \varepsilon t \\ -i \sin \varepsilon t \end{pmatrix} e^{-it}$
 $= (\cos \varepsilon t, i \sin \varepsilon t) \begin{pmatrix} -\sin \varepsilon t \\ i \cos \varepsilon t \end{pmatrix} = -2 \sin \varepsilon t \cos \varepsilon t = -\sin 2\varepsilon t$

~~Il σ_x ha autovalori ± 1 e autovettori $|c=1\rangle$ e $|c=-1\rangle$~~

* $C = \sigma_z$ ha autovettori $|c=1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$, $|c=-1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$

$P(c=1) = |\langle c=1 | \psi(t) \rangle|^2 = \left| \frac{1}{\sqrt{2}} (1, -i) \cdot e^{-it} \begin{pmatrix} \cos \varepsilon t \\ -i \sin \varepsilon t \end{pmatrix} \right|^2$
 $= \left| \frac{\cos \varepsilon t + \sin \varepsilon t}{\sqrt{2}} \right|^2 = \frac{1 + \sin 2\varepsilon t}{2}$

$P(c=-1) = |\langle c=-1 | \psi(t) \rangle|^2 = \left| \frac{\cos \varepsilon t - \sin \varepsilon t}{\sqrt{2}} \right|^2 = \frac{1 - \sin 2\varepsilon t}{2}$

• Teorema di Ehrenfest: $\frac{d}{dt} \langle C \rangle = \left\langle \frac{\partial C}{\partial t} \right\rangle + \frac{1}{i\hbar} \langle [C, H] \rangle$

$\left\{ \begin{array}{l} [C, H] = [\sigma_z, \hbar I + \hbar \varepsilon \sigma_x] = -i\hbar \varepsilon \sigma_y \\ \frac{\partial C}{\partial t} = 0 \end{array} \right.$

$\langle [C, H] \rangle = -i\hbar \varepsilon (\cos \varepsilon t, i \sin \varepsilon t) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} \cos \varepsilon t \\ -i \sin \varepsilon t \end{pmatrix}$
 $= -i\hbar \varepsilon (\cos^2 \varepsilon t - \sin^2 \varepsilon t)$
 $= -i\hbar \varepsilon \cos 2\varepsilon t$

e quindi $\frac{d\langle C \rangle}{dt} = -2 \cos 2\epsilon t = \frac{1}{i\hbar} \langle [C, H] \rangle$

III) $\psi = f(r) r \left(1 + i \frac{y}{r} \right)$
 solo e parso angolare
 serve per studiare L^2, L_x, L_z

Ora $Y_{00} = \frac{1}{\sqrt{4\pi}}$, $Y_{1, \pm 1} = \mp \sqrt{\frac{3}{8\pi}} \frac{x \pm iy}{r} \rightarrow \frac{iy}{r} = -\frac{Y_{11} + Y_{1,-1}}{2} \sqrt{\frac{8\pi}{3}}$

$\psi(x) \cong C \left(\sqrt{4\pi} Y_{00} - \sqrt{4\pi} \frac{Y_{11} + Y_{1,-1}}{\sqrt{6}} \right) \xrightarrow{\text{normalization}} \psi \sim \frac{\sqrt{6} Y_{00} - Y_{11} - Y_{1,-1}}{\sqrt{8}}$

$\psi = \sqrt{\frac{6}{8}} |00\rangle - \frac{1}{\sqrt{8}} |11\rangle - \frac{1}{\sqrt{8}} |1,-1\rangle$

• $L^2 = \hbar^2 \ell(\ell+1)$ e $\ell=0$ oppure $\ell=1$

$P(\ell=0) = 6/8$

$P(\ell=1) = 1/8 + 1/8 = 2/8$

$L_z = \hbar m$ e $m=0, m=1, m=-1$

$P(m=0) = 6/8$ $P(m=1) = 1/8$ $P(m=-1) = 1/8$

$L_x = 0, \pm \hbar$ con autofunzioni $|00\rangle_x = |00\rangle$

$|1 \pm i\rangle_x = \frac{|11\rangle \pm \sqrt{2}|10\rangle + |1,-1\rangle}{2} \rightarrow \frac{1}{2} \begin{pmatrix} 1 \\ \pm\sqrt{2} \\ 1 \end{pmatrix}$

$|10\rangle_x = \frac{|11\rangle - |1,-1\rangle}{\sqrt{2}} \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}$

$\psi = \sqrt{\frac{6}{8}} |00\rangle_x - \frac{1}{\sqrt{8}} (|1,-1\rangle_x + |11\rangle_x)$

$P(L_x=0) = 6/8$

$P(L_x=\hbar) = 1/8$

$P(L_x=-\hbar) = 1/8$

• $\langle L_z \rangle = \sum P(L_z) \hbar m = \frac{6}{8} \cdot 0 + \frac{1}{8} \hbar + \frac{1}{8} (-\hbar) = 0$

$\langle L_x \rangle = \sum P(L_x) \ell_x = \frac{6}{8} \cdot 0 + \frac{1}{8} \hbar + \frac{1}{8} (-\hbar) = 0$

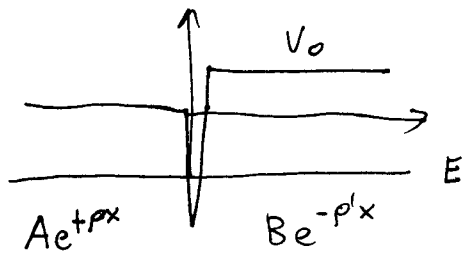
• $\psi(t) = \sqrt{\frac{6}{8}} |00\rangle_x - \frac{1}{\sqrt{8}} \left(|1,-1\rangle_x e^{i\mu B t} + |11\rangle_x e^{-i\mu B t} \right)$

oppure, se uno vuole mischiarla:

$$\psi(t) = \sqrt{\frac{6}{8}} |00\rangle - \frac{1}{\sqrt{8}} (|11\rangle \cos \mu B t + |1-1\rangle \cos \mu B t - i\sqrt{2} \sin \mu B t |10\rangle)$$

oppure: $\psi(x,t) = f(x) (v + iy \cos \mu B t + iz \sin \mu B t)$

IV

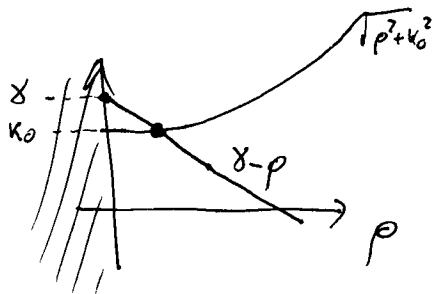


$\frac{\hbar^2 p^2}{2m} = -E$ $\frac{\hbar^2 p'^2}{2m} = V_0 - E \Rightarrow p'^2 = p^2 + \frac{2mV_0}{\hbar^2} \equiv p^2 + k_0^2$

Condizioni di raccordo $\begin{cases} \psi(0+) = \psi(0-) \\ \psi'(0+) - \psi'(0-) = -\gamma \psi(0) \end{cases} \begin{cases} A = B \\ -p'B - pA = -\gamma A \end{cases}$

\Downarrow
 $p + p' = \gamma$

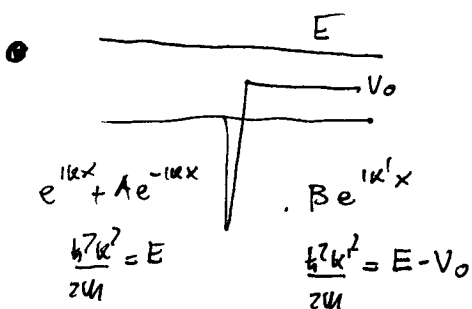
$$\sqrt{p^2 + k_0^2} = \gamma - p$$



stato legato se $\gamma > k_0 \Rightarrow k_0^2 \leq \gamma^2 \Rightarrow V_0 \leq \frac{\hbar^2 \gamma^2}{2m}$

$\int_{-\infty}^{+\infty} |\varphi(x)|^2 dx = \frac{|A|^2}{2p} + \frac{|B|^2}{2p'} = |A|^2 \frac{\gamma}{2pp'} \equiv 1 \Rightarrow A = \sqrt{\frac{2pp'}{\gamma}}$

$P(x < 0) = \int_{-\infty}^0 |\varphi(x)|^2 dx = \frac{2pp'}{\gamma} \int_{-\infty}^0 e^{2px} dx = \frac{p'}{\gamma}$



Raccordo: $1 + A = B$
 $ik'B - ik(1 - A) = -\gamma B$
 \Downarrow
 $(\gamma + ik)(1 + A) - ik(1 - A) = 0$

$$A = \frac{-\gamma + i(k - k')}{\gamma + i(k + k')}$$

$\Rightarrow R = |A|^2 \quad T = 1 - |A|^2 \Rightarrow R = \frac{\gamma^2 + (k - k')^2}{\gamma^2 + (k + k')^2} \quad T = \frac{4kk'}{\gamma^2 + (k + k')^2}$