The Shower Monte Carlo picture of hard high energy reactions

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Outline

Introduction:

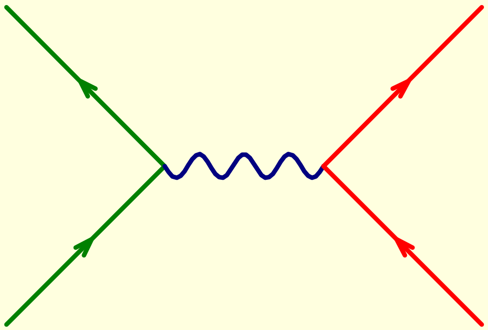
- High energy collisions and QCD
- The shower picture of hard interactions
- Shower Monte Carlo programs

Three topics on Shower improvements

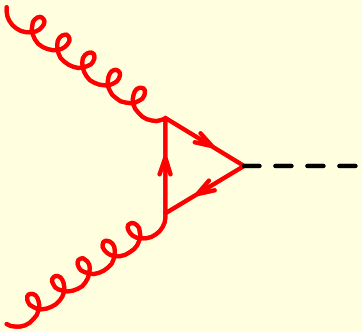
- NLO and showers: POWHEG
- POWHEG for general processes
- Truncated showers in angular ordered framework

High Energy Collisions

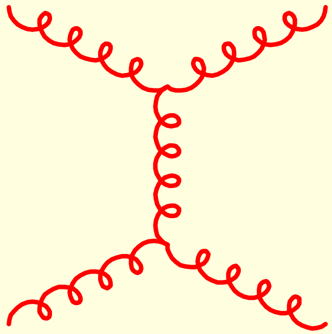
Frontier research in particle physics studies scattering and production of elementary constituents:



$e^+ e^- \to q\,\bar{q}$



$gg \to \mathrm{Higgs}$

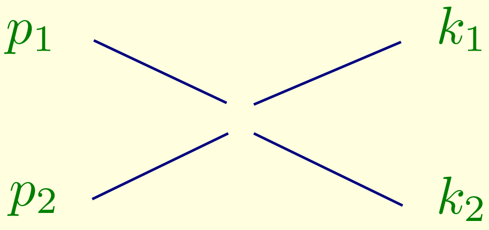


$gg \to gg$

Ideally, one needs elementary constituents as projectiles and targets, (i.e. a collider for leptons, gluons and quarks) and a final state detector of leptons, gluons and quarks. Not obvious for quarks and gluons:

- At short distance: asymptotic freedom, quarks and gluons behave as free particles
- At long distance: infrared slavery, very strong interactions hide the simplicity of constituent description

What is the characteristic distance?

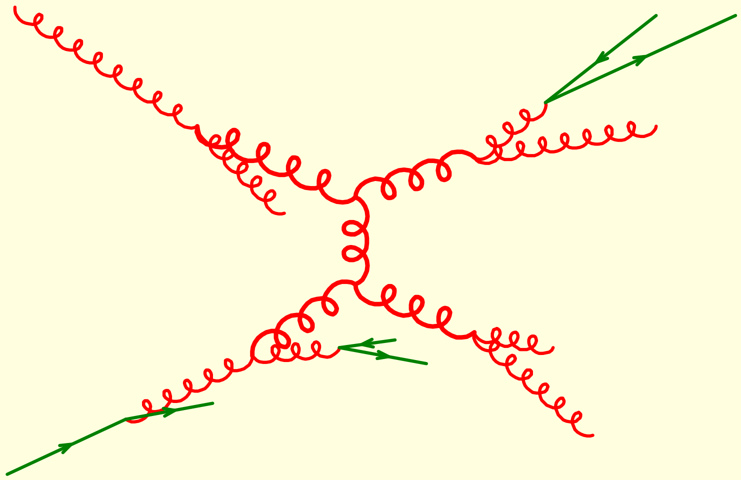


If the CM energy is large, and the momentum transferred are large, the reaction is well defined even if we assume an indeterminacy on the momentum of the particles, as long as it is less than the large scales:

$$\delta p_1 \approx \delta p_2 \approx \delta k_1 \approx \delta k_2 \lessapprox \sqrt{s} \approx \sqrt{t} \approx \sqrt{u}$$

If δp, δk are larger than the characteristic scale Λ_{QCD} (when strong interactions become strong) the QCD description of the process is adequate thanks to asymptotic freedom. The distances involved are of order $h/\delta p < h/\Lambda_{\mathrm{QCD}}$.
If our collider and experiment were much smaller than h/Λ_{QCD} the process would be fully calculable using perturbation theory.

Dominant corrections



Collinear splitting processes in the initial and final state (always with transverse momenta $> \Lambda_{\mathrm{QCD}}$) are strongly enhanced. This is due to the fact that in perturbation theory the energy denominators are small. There are algorithms to evaluate all these enhanced contributions: The

Shower algorithms

Shower algorithms give a description of a hard collision up to distances of order h/Λ_{QCD}. At larger distances, theory is of little help:
Perturbation theory breaks down, need to resort to non-perturbative methods (i.e. lattice calculations). However, these methods can be applied only to symple systems.
The only viable alternative is to use models of hadron formation.

Shower Monte Carlo programs

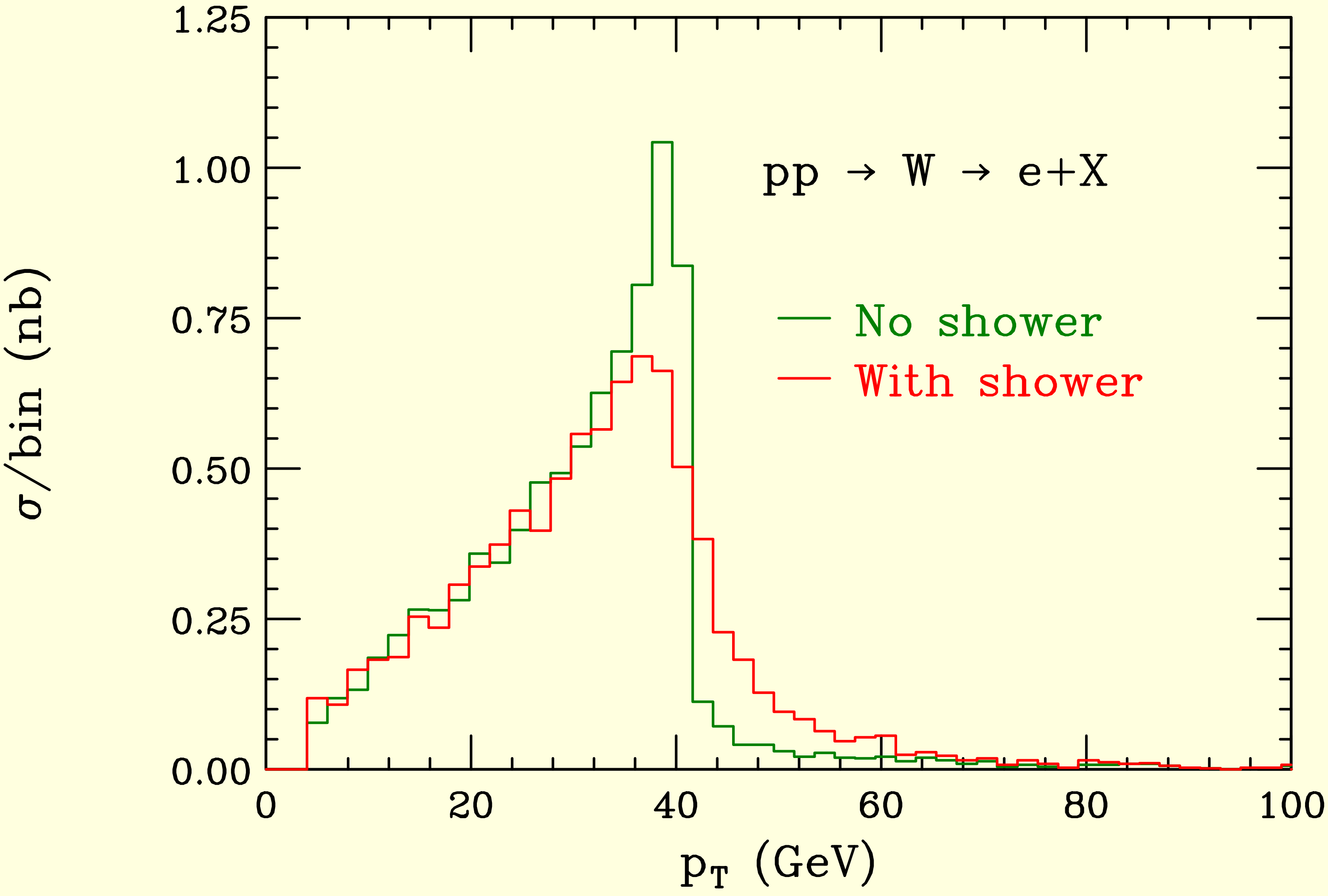
Capabilities

- Large library of hard events cross sections (SM and BSM)
- Dress hard events with QCD radiation
- Models for hadron formation
- Models for underlying event, multi-parton collisions, minimum bias
- Library for (spacetime) decays of unstable particles

The name SHOWER from item 2.
The hope (and the experience) is: the "Models" part is the same at all energies, and process independent

Once tuned at some energy, the SMC is predictive for all other energies.

An example: (half an our of work)



Detailed description of the final state for each generated event:

IHEP	ID	IDPDG	IST	MO1	MO2	DA1	DA2	P-X	P-Y	P-Z	ENERGY	MASS	V-X	V-Y	V-Z	V-C*T
30	NU_E	12	1	28	23	0	0	64.30	25.12	-1194.4	1196.4	0.00	0.000E+00	0.000E+00	0.000E+00	0.000E+00
31	E+	-11	1	29	23	0	0	-22.36	6.19	-234.2	235.4	0.00	0.000E+00	0.000E+00	0.000E+00	0.000E+00
230	PI0	111	1	155	24	0	0	0.31	0.38	0.9	1.0	0.13	4.209E-11	6.148E-11	-3.341E-11	5.192E-10
231	RHO+	213	197	155	24	317	318	-0.06	0.07	0.1	0.8	0.77	4.183E-11	6.130E-11	-3.365E-11	5.189E-10
232	P	2212	1	156	24	0	0	0.40	0.78	1.0	1.6	0.94	4.156E-11	6.029E-11	-4.205E-11	5.250E-10
233	NBAR	-2112	1	156	24	0	0	-0.13	-0.35	-0.9	1.3	0.94	4.168E-11	6.021E-11	-4.217E-11	5.249E-10
234	PI-	-211	1	157	9	0	0	0.14	0.34	286.9	286.9	0.14	4.660E-13	8.237E-12	1.748E-09	1.749E-09
235	PI+	211	1	157	9	0	0	-0.14	-0.34	624.5	624.5	0.14	4.056E-13	8.532E-12	2.462E-09	2.462E-09
236	P	2212	1	158	9	0	0	-1.23	-0.26	0.9	1.8	0.94	-4.815E-11	1.893E-11	7.520E-12	3.252E-10
237	DLTABR--	-2224	197	158	9	319	320	0.94	0.35	1.6	2.2	1.23	-4.817E-11	1.900E-11	7.482E-12	3.252E-10
238	PI0	111	1	159	9	0	0	0.74	-0.31	-27.9	27.9	0.13	-1.889E-10	9.893E-11	-2.123E-09	2.157E-09
239	RHO0	113	197	159	9	321	322	0.73	-0.88	-19.5	19.5	0.77	-1.888E-10	9.859E-11	-2.129E-09	2.163E-09
240	K+	321	1	160	9	0	0	0.58	0.02	-11.0	11.0	0.49	-1.890E-10	9.873E-11	-2.135E-09	2.169E-09
241	KL_1-	-10323	197	160	9	323	324	1.23	-1.50	-50.2	50.2	1.57	-1.890E-10	9.879E-11	-2.132E-09	2.166E-09
242	K-	-321	1	161	24	0	0	0.01	0.22	1.3	1.4	0.49	4.250E-11	6.333E-11	-2.746E-11	5.211E-10
243	PI0	111	1	161	24	0	0	0.31	0.38	0.2	0.6	0.13	4.301E-11	6.282E-11	-2.751E-11	5.210E-10

HEP experiments feed this kind of output through their detector simulation software, and use it to determine efficiencies for signal detection, and perform background estimates. Analysis strategies are set up using these simulated data.

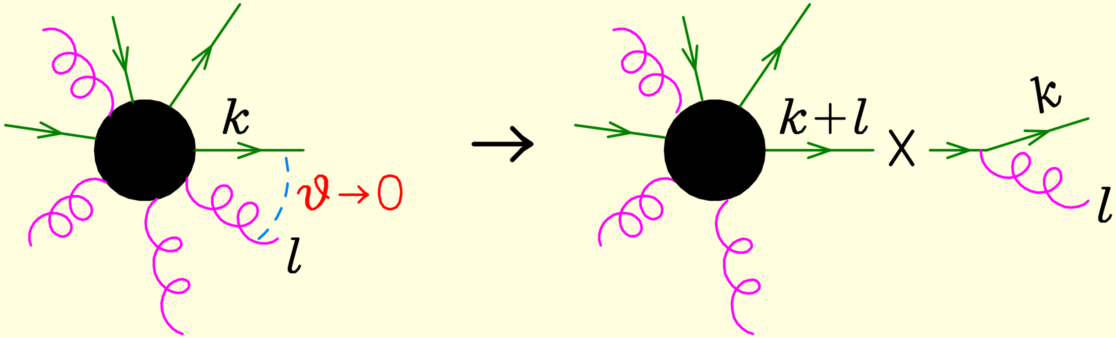
Summarizing:

- In HEP (i.e. collider physics) not many questions can be answered without a Shower Monte Carlo (SMC). Heavily used since 1980's
- SMC's are forever (well, as long as HEP lives). Even if QCD was solved exactly, it is unlikely that complex, high energy phenomena will be described better than in SMC models.
- After LEP, QCD testing is less important. With LHC, QCD modeling is a primary issue.
- SMC models have long been neglected in theoretical physics:
  Emphasis on QCD tests required more transparent theoretical methods.
  After LEP, QCD testing is less important.
  With LHC, QCD modeling is a primary issue: recent SMC revival.
- Thinking in terms of Shower algorithms gives us an easy to grasp, intuitive understanding of complex QCD phenomena (and a practical way to verify our ideas).

Shower basics: Collinear factorization

QCD emissions are enhanced near the collinear limit

Cross sections factorize near collinear limit



$$|M_{n+1}|^2 d\Phi_{n+1} \Longrightarrow |M_n|^2 d\Phi_n \frac{\alpha_s}{2\pi}\frac{dt}{t} P_{q,qg}(z)\, dz \frac{d\phi}{2\pi}$$

t : virtuality (or p_T^2, or $E^2\theta^2$)

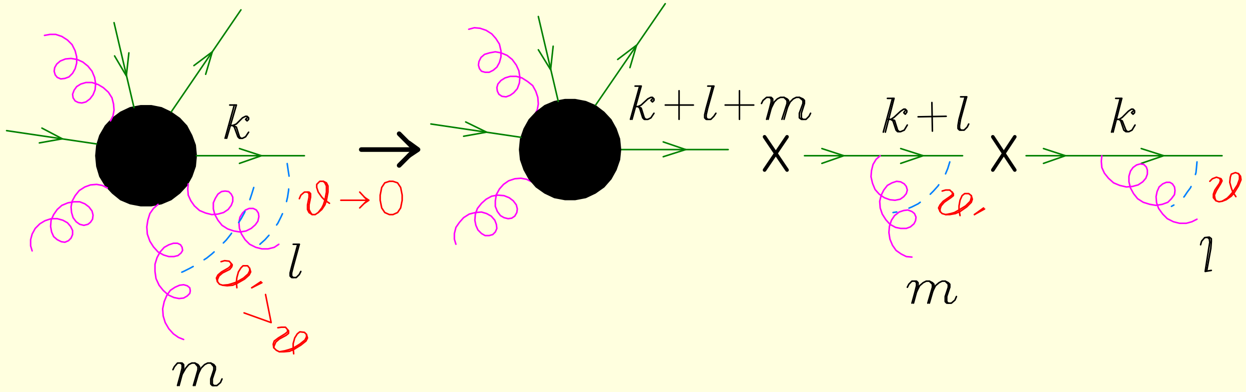
$z = k^0/(k^0 + l^0)$: energy (or $p_\parallel$, or p^+) fraction of quark

$P_{q,qg}(z) = C_F \frac{1+z^2}{1-z}$: Altarelli – Parisi splitting function

(ignore $z \to 1$ IR divergence for now)

If another gluon becomes collinear, iterate the previous formula:

$\theta', \theta \to 0$
with $\theta' > \theta$



$$|M_{n+1}|^2 d\Phi_{n+1} \Longrightarrow |M_{n-1}|^2 d\Phi_{n-1} \times \frac{\alpha_s}{2\pi} \frac{dt'}{t'} P_{q,qg}(z')\, dz' \frac{d\phi'}{2\pi}$$

$$\times \frac{\alpha_s}{2\pi} \frac{dt}{t} P_{q,qg}(z)\, dz \frac{d\phi}{2\pi} \theta(t'-t)$$

Collinear partons can be described by a factorized integral ordered in t.
For m collinear emissions:

$$\int_{\theta_{\min}} d\theta_1 \int_{\theta_1} d\theta_2 \ldots \int_{\theta_{m-1}} d\theta_m \;\; \propto \; \frac{\log^m \frac{1}{\theta_{\min}^2}}{m!} \approx \frac{\log^m \frac{Q^2}{\Lambda^2}}{m!}, \;\; \Lambda \approx \Lambda_{\rm QCD}$$

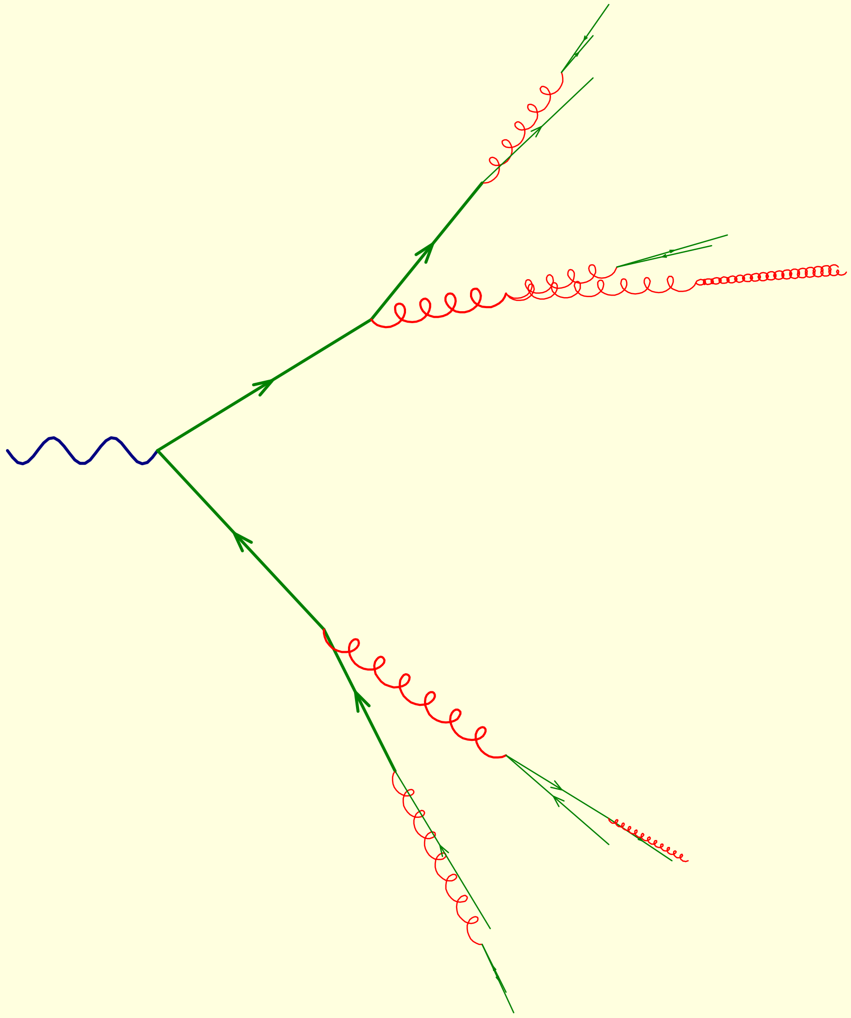
Typical dominant configuration at very high Q^2

(Example: $\gamma^\star \to \mathrm{hadrons}$)
Besides $q \to qg$, also $g \to gg$, $g \to q\bar{q}$ come into play.

Typical configurations: intermediate angles of order of geometric average of upstream and downstream angles.

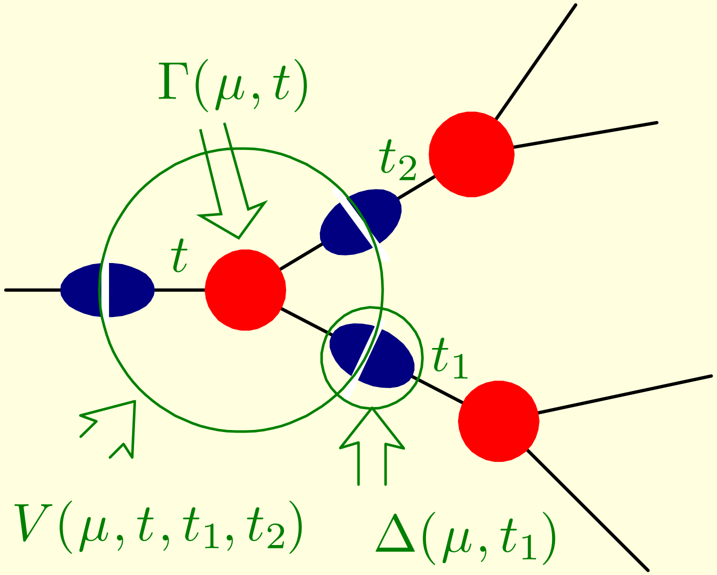
Each angle is $\mathcal{O}(\alpha_s)$ smaller than its upstream angle, and $\mathcal{O}(\alpha_s)$ bigger than its downstream angle.

As relative momenta become smaller α_s becomes bigger, and this picture breaks down.



For a consistent description:

include virtual corrections to same LL approximation



Effective (RG invariant) splitting vertex:

$$V^2(\mu,t,t_1,t_2) = \Gamma^2(\mu,t)\Delta(\mu,t)\Delta(\mu,t_1)\Delta(\mu,t_2)$$

Choosing $\mu = t$ (using $\Delta(t,t) \approx 1$)

$$V^2(\mu,t,t_1,t_2) = V^2(t,t,t,t)\,\Delta(t,t_1)\Delta(t,t_2)$$

$V(t,t,t,t)$ is the three level vertex with $\alpha \to \alpha(t)$.
The form $\Delta(t,t_1)$ follows from RG arguments.

In fact: $\Delta_i(t,t_1) = \exp\left[-\sum_{(jk)} \int_{t_1}^{t} \frac{dt'}{t'} \int dz \, \frac{\alpha_s(t')}{2\pi} P_{i,jk}(z) \right]$ Sudakov form factor

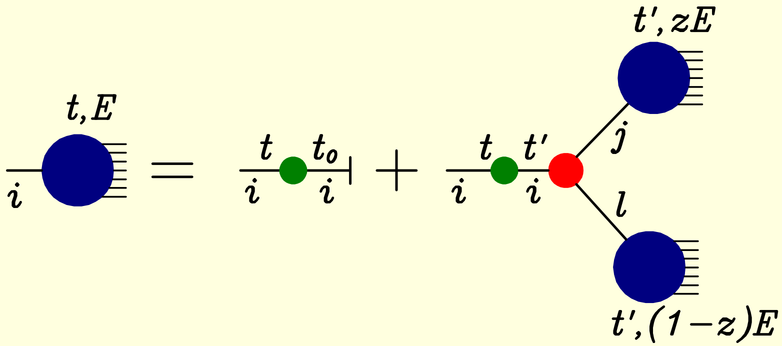
consistent with KLN cancellation of IR singularities, and with RG.

Final Recipe

- Consider all tree graphs.
- Assign ordered hardness parameters t to each vertex.
- Include a factor
$$\frac{dt}{t}\,dz\,\frac{\alpha_s(t)}{2\pi}P_{i,jk}(z)$$
at each vertex $i \to jk$.
- Include a factor $\Delta_i(t_1, t_2)$ to each internal line with a parton i, from hardness t_1 to hardness t_2.
- Include a factor $\Delta_i(t, t_0)$ on final lines (t_0: IR cutoff)

recipe can be written as (omitting angular dependencies)

$$\mathcal{S}_i(t,E) = \Delta_i(t,t_0)\mathcal{S}_i(t_0) + \sum_{(jl)} \int_{t_0}^{t} \frac{\alpha_S(t')}{2\pi}\frac{d\,t'}{t'} P_{i,jl}(z)\,d\,z\frac{d\phi}{2\pi}\Delta_i(t,t')\,\mathcal{S}_j(t',zE)\,\mathcal{S}_l(t',(1-z)E)$$

Graphically: 

It also satisfies the differential equation:

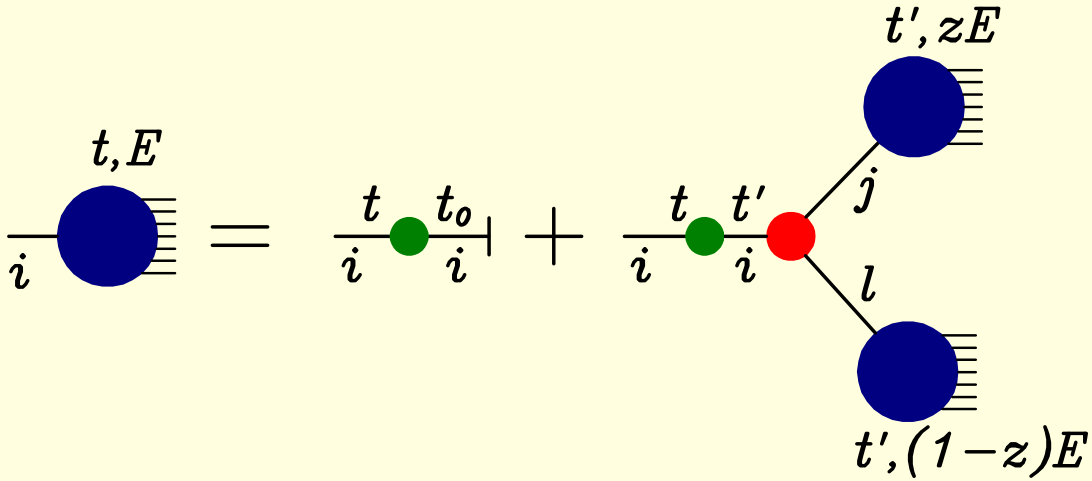
$$t\frac{\partial\mathcal{S}_i(t,E)}{\partial t} = \sum_{(jl)} \int_0^1 \frac{\alpha_S(t)}{2\pi} P_{i,jl}(z)\,d\,z\frac{d\phi}{2\pi}\mathcal{S}_j(t,zE)\,\mathcal{S}_l(t,(1-z)E) - \mathcal{S}_i(t,E)\sum_{(jl)}\int_0^1 d\,z\,\frac{\alpha_S(t)}{2\pi}P_{i,jl}(z)$$

Easy to show now that $\mathcal{S}_i^{\rm inc}(t,E) = \mathcal{O}_{\rm inc}\cdot\mathcal{S}_i(t,E) = 1$ (KLN cancellation!)

Introducing suitable observables one can easily prove the evolution equations for fragmentation functions.

Collinear radiation from initial state can be treated similarly.
One can derive a recipe in the presence of initial state radiation.
One can derive the evolution equations for parton densities.

But, (most important) the recipe can be easily implemented as a computer code!



The probability of the first branching is independent of subsequent branchings because of KLN cancellation. It is given by

$$\frac{\alpha_S(t')}{2\pi}\frac{d\,t'}{t'}P_{i\,,jk}(z)\,d\,z\frac{d\phi}{2\pi}\Delta_i(t,t')$$

Integrating in $d\,z, d\phi$, summing over jk, the t' distribution is

$$\Delta_i(t,t')\,\frac{\alpha_S(t')}{2\pi}\frac{d\,t'}{t'}\int\sum_{(jk)}P_{i\,,jk}(z)\,d\,z = d\,\Delta_i(t,t')$$

i.e., the distribution is uniform in the Sudakov form factor!

Shower Algorithm

- Generate a uniform random number $0 < r < 1$;
- Solve the equation $\Delta_i(t, t') = r$ for t';
- If $t' < t_0$ stop here (final state line);
- generate z, jk with probability $P_{i,jk}(z)$, and $0 < \phi < 2\pi$ uniformly;
- restart from each branch, with hardness parameter t'.

Elementary example

Simulate a radioactive source with emission probability p in unit time.

Probability distribution for first emission:

$$P(t)\,dt = \lim_{n\to\infty}\left(1 - p\frac{t}{n}\right)^n p\,dt = e^{-pt}\,p\,dt = -d(e^{-pt})$$

uniform in $0 < e^{-pt} < 1$. Monte Carlo implementation for $t_0 < t < t_f$:

- generate a random number $0 < r < 1$
- solve the equation $e^{-p(t-t_0)} = r$ for t
- if $t > t_f$ stop
- Continue setting $t_0 = t$.

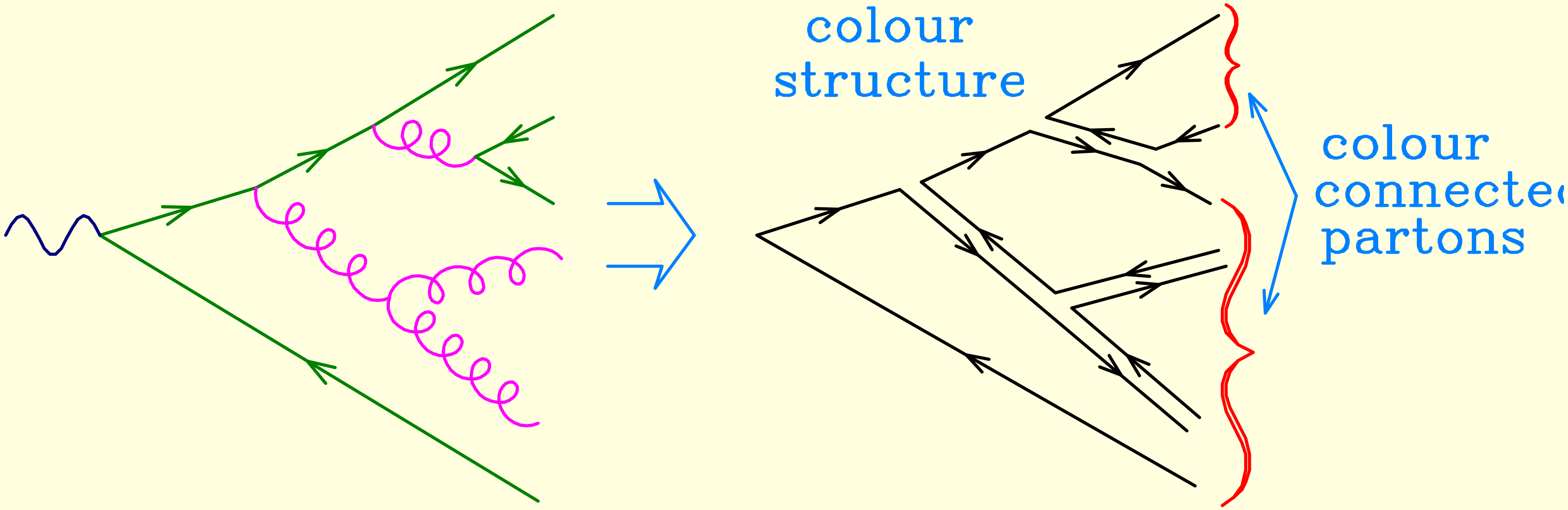
Notice: Virtual corrections = no-emission probability (easy to teach!)

COLOUR AND HADRONIZATION

SMC's assign colour labels to partons.

Only colour connections are recorded (as in large N limit).

Initial colour assigned according to hard cross section.



Colour assignements are used in the hadronization model.

Most popular models: Lund String Model, Cluster Model.

In all models, color singlect structures are formed out of colour connected partons, and are decayed into hadrons preserving enery and momentum.

Implementation

- COJETS Odorico (1984)
- ISAJET Page+Protopopescu (1986)
- FIELDAJET Field (1986)
- JETSET Sjöstrand (1986)
- PYTHIA Bengtsson+Sjöstrand (1987), Skands+Sjöstrand
- Ariadne Lönnblad (1991)
- HERWIG Marchesini+Webber (1988)
  Marchesini+Webber+Abbiendi+Knowles+Seymour+Stanco (1992)
- SHERPA Gleisberg+Hoche+Krauss+Schalicke+Schumann+Winter (2004)

Accuracy: soft divergences and double log region

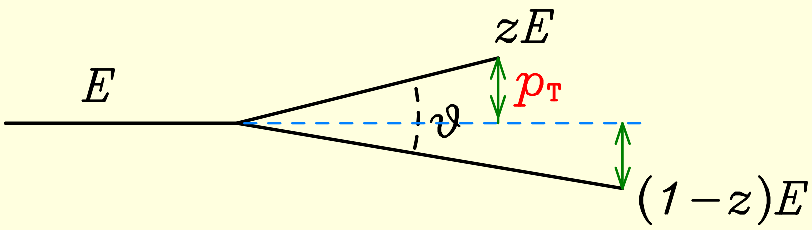
$z \to 1$ ($z \to 0$) region problematic: for $z \to 1$: $P_{qq}, P_{gg} \propto \frac{1}{1-z}$

Choice of hardness variable makes a difference

virtuality: $t \equiv E^2 z(1-z) \overbrace{\theta^2}^{1-\cos\theta}$

p_T^2: $t \equiv E^2 z^2 (1-z)^2 \, \theta^2$

angle: $t \equiv E^2 \, \theta^2$

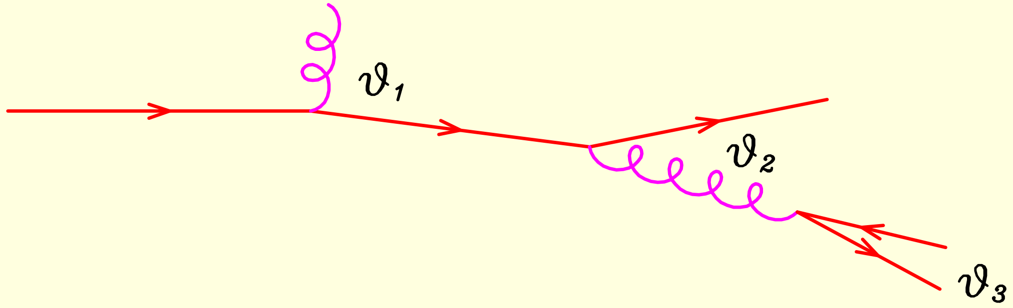


$$\underbrace{\int \frac{dt}{t} \int_{\sqrt{t}/E}^{1-\sqrt{t}/E} \frac{dz}{1-z}}_{\text{virtuality}: z(1-z) > t/E^2} \approx \frac{\log^2 \frac{t}{E^2}}{4}; \underbrace{\int \frac{dt}{t} \int_{t/E^2}^{1-t/E^2} \frac{dz}{1-z}}_{p_T^2: z^2(1-z)^2 > t/E} \approx \frac{\log^2 \frac{t}{E^2}}{2}; \underbrace{\int \frac{dt}{t} \int_0^1 \frac{dz}{1-z}}_{\text{angle}} \approx \log t \log \Lambda$$

Sizeable difference in double log structure!

Angular ordering is the correct choice (Mueller 1981)

$\frac{d\theta}{\theta}\frac{\alpha_s(p_T^2)}{2\pi}P(z)\mathrm{dz}$



$\theta_1 > \theta_2 > \theta_3 \ldots$

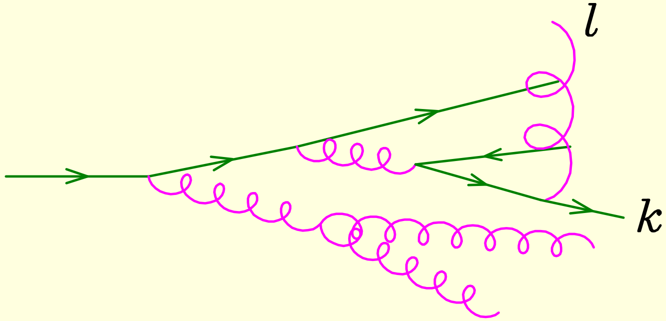
$p_T^2 = E^2 z^2 (1-z)^2\, \theta^2$

$\alpha_s(p_T^2)$ for a correct treatment of charge renormalization in soft region.

$$\Delta_i(t,t') = \exp\left[-\int_{t'}^{t}\frac{dt}{t}\int_{\sqrt{\frac{t_0}{t}}}^{1-\sqrt{\frac{t_0}{t}}} dz\,\frac{\alpha_s(p_T^2)}{2\pi}\sum_{(jk)} P_{i,jk}(z)\right]$$

$$\approx \exp\left[-\frac{c_i}{4\pi b_0}\left\{\log\frac{t}{\Lambda^2}\log\frac{\log\frac{t}{\Lambda^2}}{\log\frac{t_0}{\Lambda^2}} - \log\frac{t}{t_0}\right\}_{t'}^{t}\right] \qquad (c_q = C_F,\, c_g = 2C_A)$$

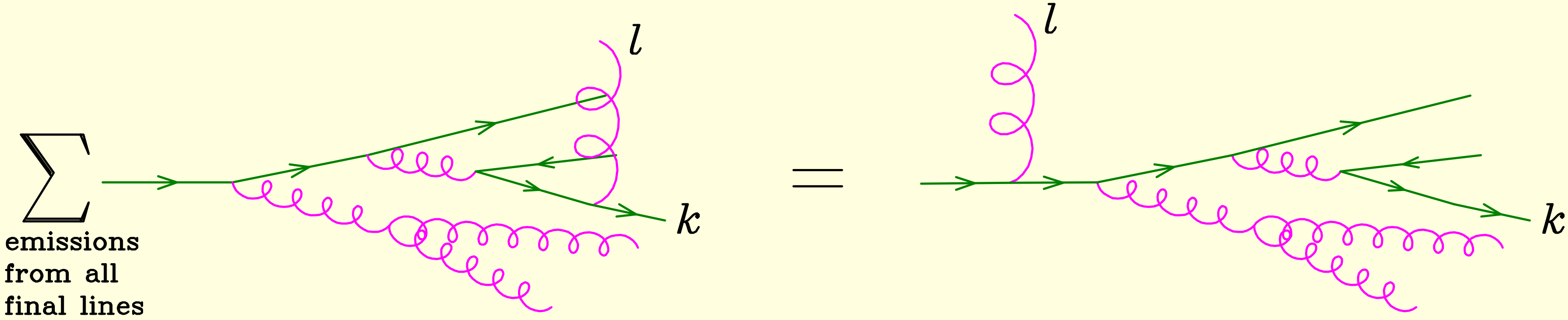
Sudakov dumping stronger than any power of t.



With virtuality ordering:

Soft emissions give small virtuality.
At end of shower, large amount of unrestricted (all angles) soft radiation

But soft gluons emitted at large angles from final state partons add coherently!



large angle, high energy: already ordered in angle
large angle, small energy: should be reordered by angle

Thus: order in angle

Accuracy

	Collinear	Soft-Collinear	Soft-large N_c	Soft
PYTHIA	Leading	Partial	No	No
HERWIG	Leading	Leading	No	No
ARIADNE	Partial	Partial	Leading	No
PYTHIA 6.4	Partial	Partial	Leading	No
SHERPA	Leading	Partial	No	No

One can realistically aim at:

Leading Collinear, Leading double log, Leading soft in large N_c limit

(Soft effects for finite N_c require matrix exponentiation in the Sudakov FF)

Not much progress in shower accuracy since the 80's.

New developements

- Interfacing ME (Matrix-Elements) generators with Parton Showers (CKKW matching (Catani, Krauss, Küen, Webber), MLM matching)
- Interfacing NLO calculations to Parton Showers (MC@NLO (Frixione, Webber), POWHEG (PN))

Several other approaches have appeared:

- Kramer, Mrenna, Soper ($e^+e^- \to 3\,\mathrm{partons}$)
- Shower by antenna factorization (Frederix,Giele,Kosower,Skands) (toy implementation for $H \to gg$)
- Shower by Catani-Seymour dipole factorization (Schumann)
- Shower with quantum interference (Nagy,Soper)
- Shower by Soft Collinear Effective Theory (Bauer,Schwartz)

Until now, complete results for hadron colliders only from MC@NLO and POWHEG

NLO+Shower

LO-ME good for shapes; uncertain absolute normalizations.

$$\alpha_s^n(2\mu) \approx \alpha_s^n(\mu)(1 - b_0\alpha_s(\mu)\log(4))^n \approx \alpha_s(\mu)(1 - n\alpha_s(\mu))$$

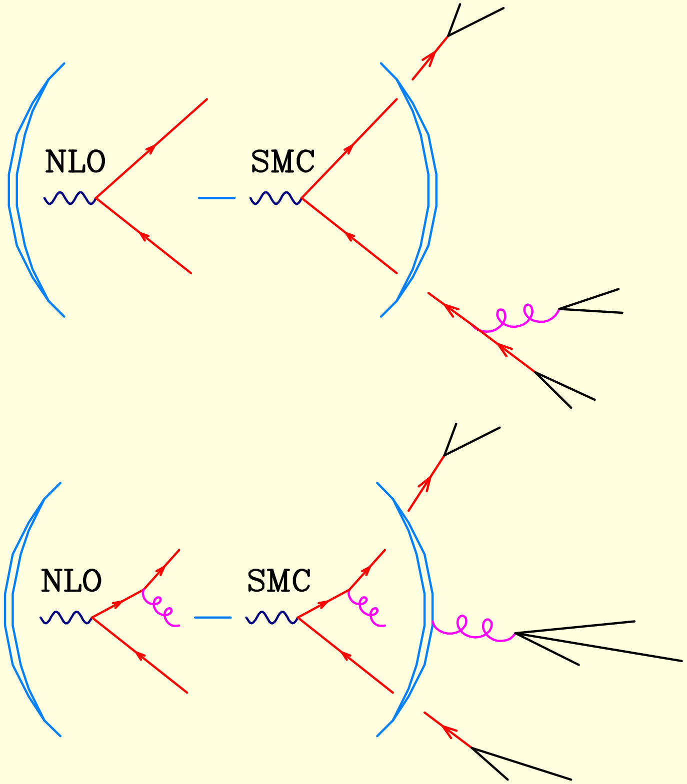
For $\mu = 100\,\mathrm{GeV}$, $\alpha_s = 0.12$;
Normalization uncertainty:

$W + 1J$	$W + 2J$	$W + 3J$
$\pm 12\%$	$\pm 24\%$	$\pm 36\%$

To improve on this, need to go to NLO

- Positive experience with NLO calculations at LEP, HERA, Tevatron (we TRUST perturbative QCD after LEP!)
- NLO results are cumbersome to use: typically made up of an n body (Born+Virtual+Soft and Collinear remnants) and $n + 1$ body (real emission) terms, both divergent (finite only when summed up).
- Merging NLO with shower: a natural extension of present approaches

MC@NLO (2002, Frixione+Webber)



Add difference between exact NLO and approximate (MC) NLO.

- Must use MC kinematics
- Difference should be regular (if the MC is OK)
- Difference may be negative

Several collider processes already there:
Vector Bosons, Vector Bosons pairs,
Higgs, Single Top.
Heavy Quarks (with P.N.)

POWHEG

Positive Weight Hardest Emission Generator

Method to generate the hardest emission first, with NLO accuracy, and independently of the SMC (P.N. 2004).

- SMC independent; no need of SMC expert; same calculation can be interfaced to several SMC programs with no extra effort
- SMC inaccuracies only affect next-to-hardest emissions; no matching problems
- As the name says, it generates events with positive weight

How it works (roughly)

In words: works like a standard Shower MC for the hardest radiation, with care to maintain higher accuracy.

Inclusive cross section $\Longrightarrow$ NLO inclusive cross section. Positive if $\mathrm{NL} < \mathrm{LO}$

$\Phi_n = \text{Born variables}$
$\Phi_r = \text{radiation vars.}$

$$\bar{B}(\Phi_n) = B(\Phi_n) + \underbrace{\left[\overbrace{V(\Phi_n)}^{\text{INFINITE}} + \overbrace{\int R(\bar{\Phi}_n, \Phi_r)\, d\Phi_r}^{\text{INFINITE}} \right]}_{\text{FINITE!}}$$

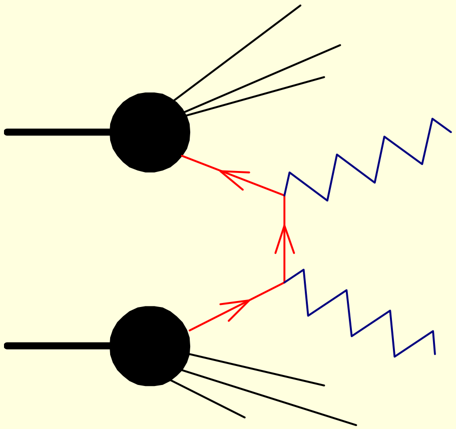
Sudakov form factor for hardest emission built from exact NLO real emission

$$\Delta_t = \exp\left[\underbrace{-\int \theta(t_r - t) \frac{R(\Phi_n, \Phi_r)}{B(\Phi_n)} d\Phi_r}_{\text{FINITE because of } \theta \text{ function}} \right]$$

with $t_r = k_T(\Phi_n, \Phi_r)$, the transverse momentum for the radiation.

First example: ZZ production in hadron collisions

(Ridolfi, P.N.)



- NLO known
  (Mele,Ridolfi, P.N.)
- Intermediate complexity
- Hadrons in initial state
- Similar to WZ, WW, $Q\bar{Q}$

$\bar{\Phi}$ and Φ_r variables

$\bar{\Phi}$ variables: choose $M_{\rm zz}$, $Y_{\rm zz}$ and θ, where

- $M_{\rm zz}$: invariant mass of the ZZ pair
- $Y_{\rm zz}$: rapidity of ZZ pair
- θ: go in the (longitudinally) boosted frame where $Y_{\rm zz} = 0$.
  go to the ZZ rest frame with a transverse boost
  In this frame θ is the angle of a Z to the longitudinal direction.

Φ_r variables:

- $x = M_{\rm zz}/s$, (s is the invariant mass of the incoming parton system)
  $x \to 1$ is the soft limit
- y: cosine of the angle of the radiated parton to the beam direction
  in the partonic CM frame.
- ϕ: radiation azimuth.

Few tricks to do it

$$\bar{B}(\Phi) = B(\Phi) + V(\Phi) + \int d\Phi_r \left[R(\Phi, \Phi_r) - C(\Phi, \Phi_r) \right]$$

Seems to need one Φ_r integrations to get weight of each Φ point.

In fact, write

$$\tilde{B}(\Phi, \Phi_r) = N\left[B(\Phi) + V(\Phi) \right] + R(\Phi, \Phi_r) - C(\Phi, \Phi_r), \qquad N = \frac{1}{\int d\Phi_r}.$$

so that

$$\bar{B}(\Phi) = \int \tilde{B}(\Phi, \Phi_r) d\Phi_r .$$

Use standard procedures (SPRING-BASES, Kawabata)
to generate unweighted events for $\tilde{B}(\bar{\Phi}, \Phi_r) d\Phi_r d\bar{\Phi}$.
discard Φ_r (same as integrating over it!).

$$\Delta(\Phi, p_T) = \exp\left[-\int \frac{R(\Phi, \Phi_r)}{B(\Phi)} \theta(k_T(\Phi, \Phi_r) - p_T) d\Phi_r \right],$$

Look for an upper bounding function;

$$\frac{R(\Phi, \Phi_r)}{B(\Phi)} \leq U(\Phi) = N \frac{\alpha_S(k_T)}{(1-x)(1-y^2)}$$

Generate x, y according to

$$\exp\left[-\int U(\Phi) \theta(k_T(\Phi, \Phi_r) - p_T) d\Phi_r \right]$$

accept the event with a probability

$$\frac{R(\Phi, \Phi_r)}{B(\Phi) U(\Phi)} \,.$$

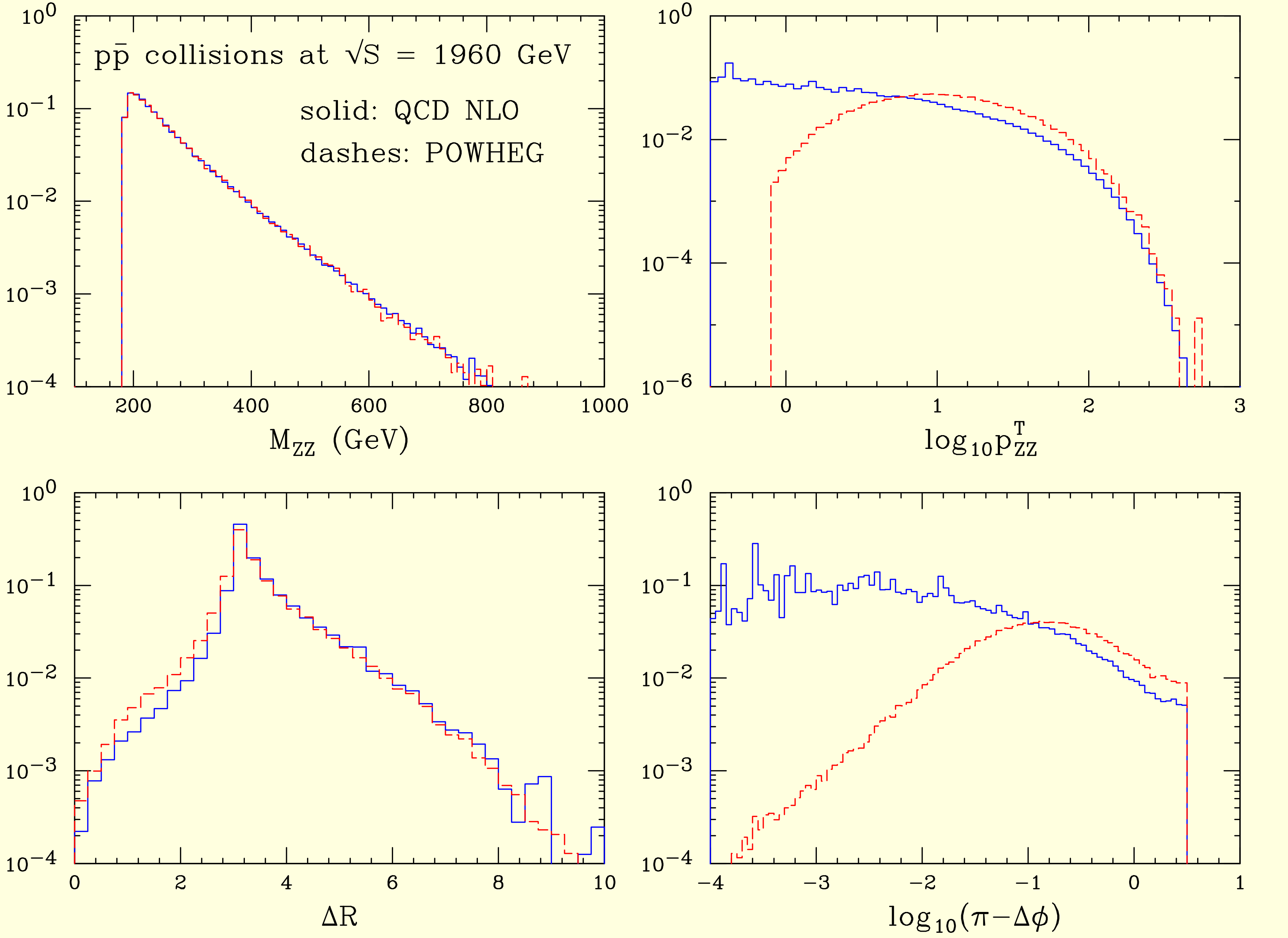
If the event is rejected generate a new one for smaller p_T, and so on
(This procedure reconstructs the exact emission probability).
In the ZZ case, an event is generated with about six calls ro $R(\Phi, \Phi_r)$.

Interfacing to SMC's

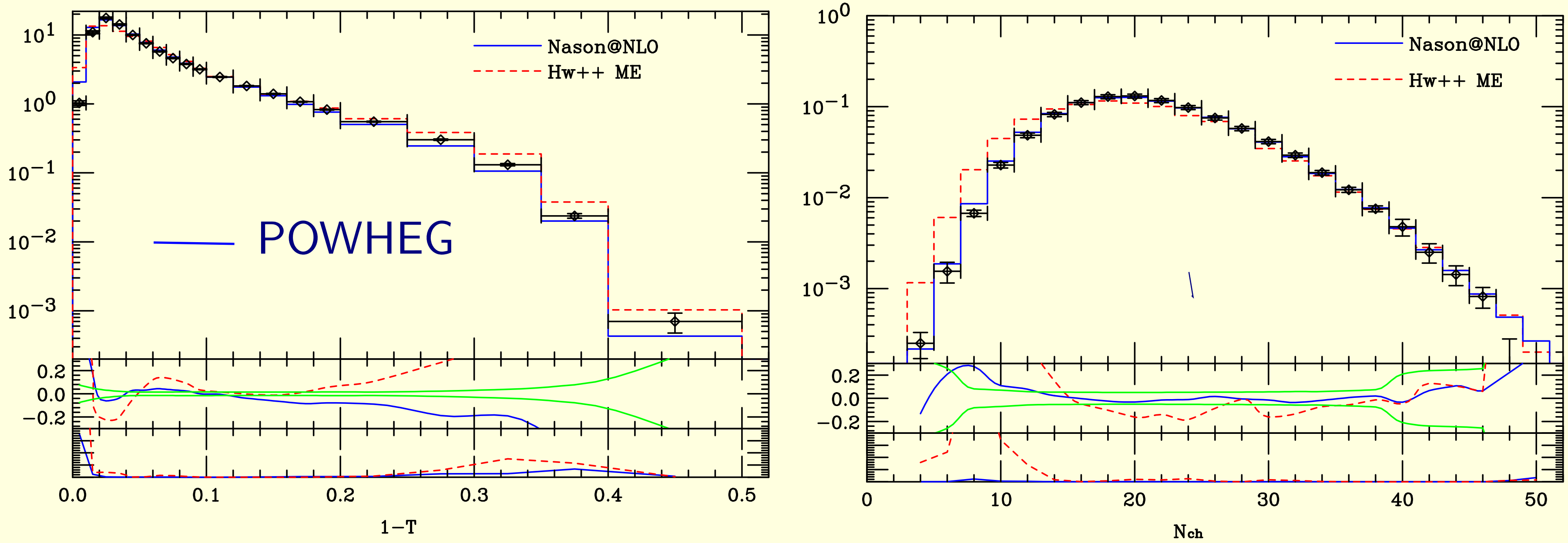
For a p_T ordered SMC, nothing else needs to be done.
Use the standard Les Houches Interface for User's Processes (LHI):
put partonic event generated by POWHEG on the LHI;
Run the SMC in the LHI mode.
The LHI provides a facility to pass the p_T of the event to the SMC (SCALUP).
As far as the hardest emission is concerned, POWHEG can reach:

- NLO accuracy of (integrated) shape variables
- Collinear, double-log, soft (large N_c) accuracy of the Sudakov FF.
  (In fact, corrections that exponentiates are obviously OK)

As far as subsequent (less hard) emissions, the output has the accuracy of the SMC one is using.

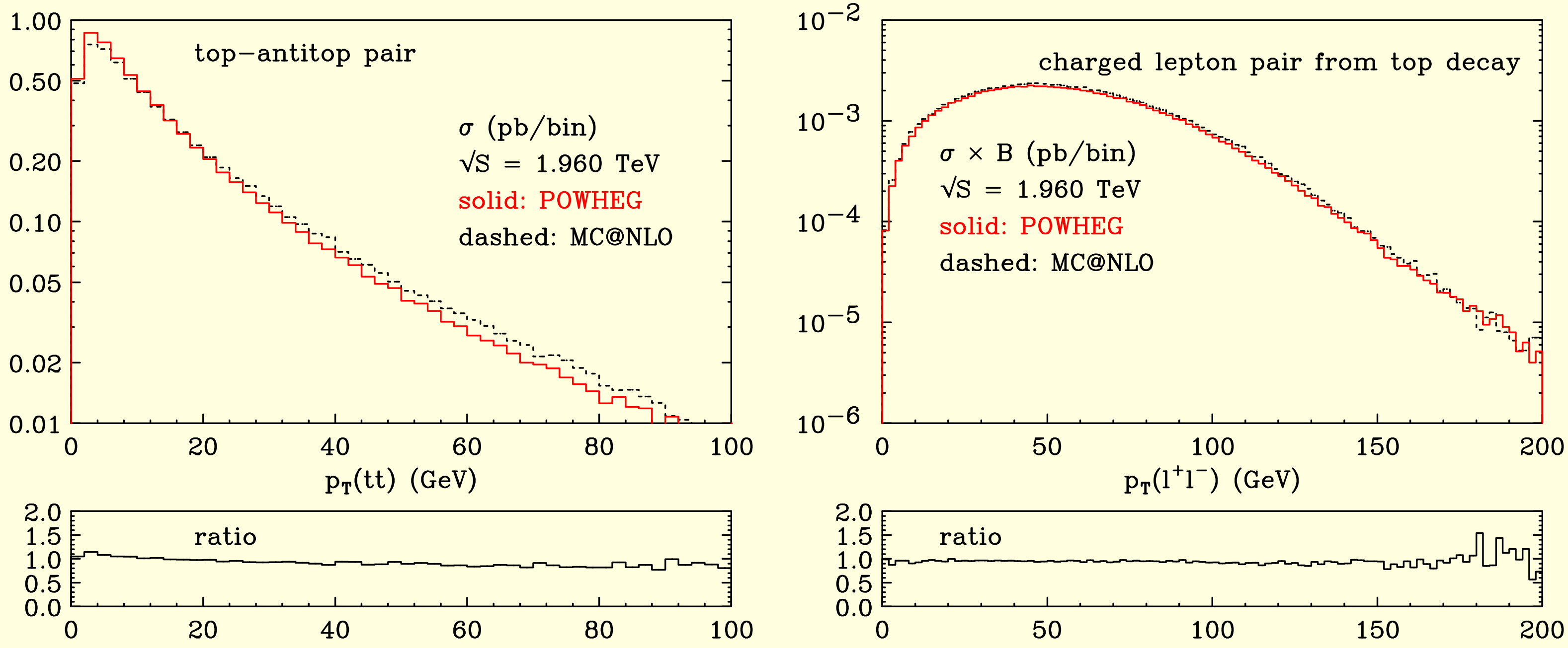

$p\bar{p}$ collisions at $\sqrt{S}$ = 1960 GeV
solid: QCD NLO
dashes: POWHEG
M_{ZZ} (GeV)
$\log_{10} p^T_{ZZ}$
ΔR
$\log_{10}(\pi - \Delta\phi)$

$e^+e^- \rightarrow$ Hadrons



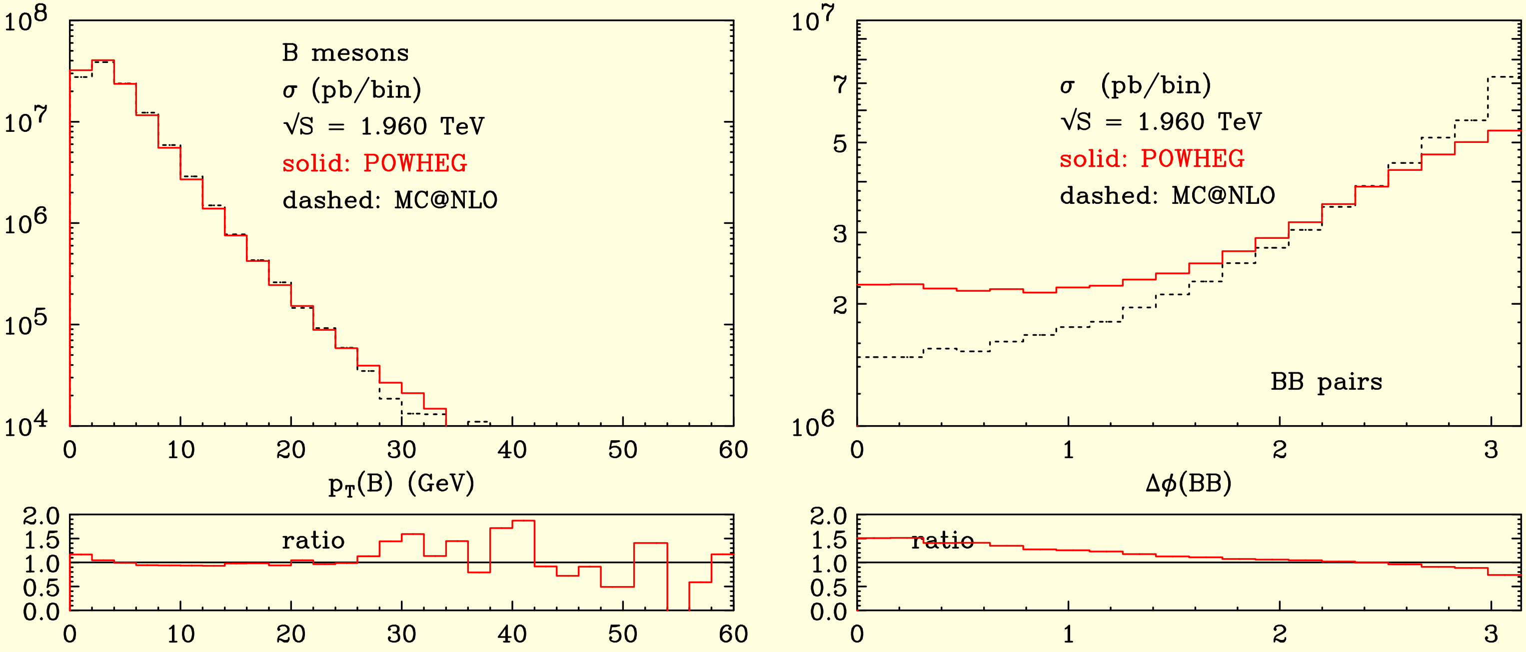
(Latunde-Dada,Gieseke,Webber, Dec.06),
fit to e^+e^- data: better than standard ME correction approach,
shown for the Thrust distribution and the charged multiplicity.

POWHEG and MC@NLO comparison: Top pair production



Good agreement for all observable considered
(differences can be ascribed to different treatment of higher order terms)

Bottom pair production



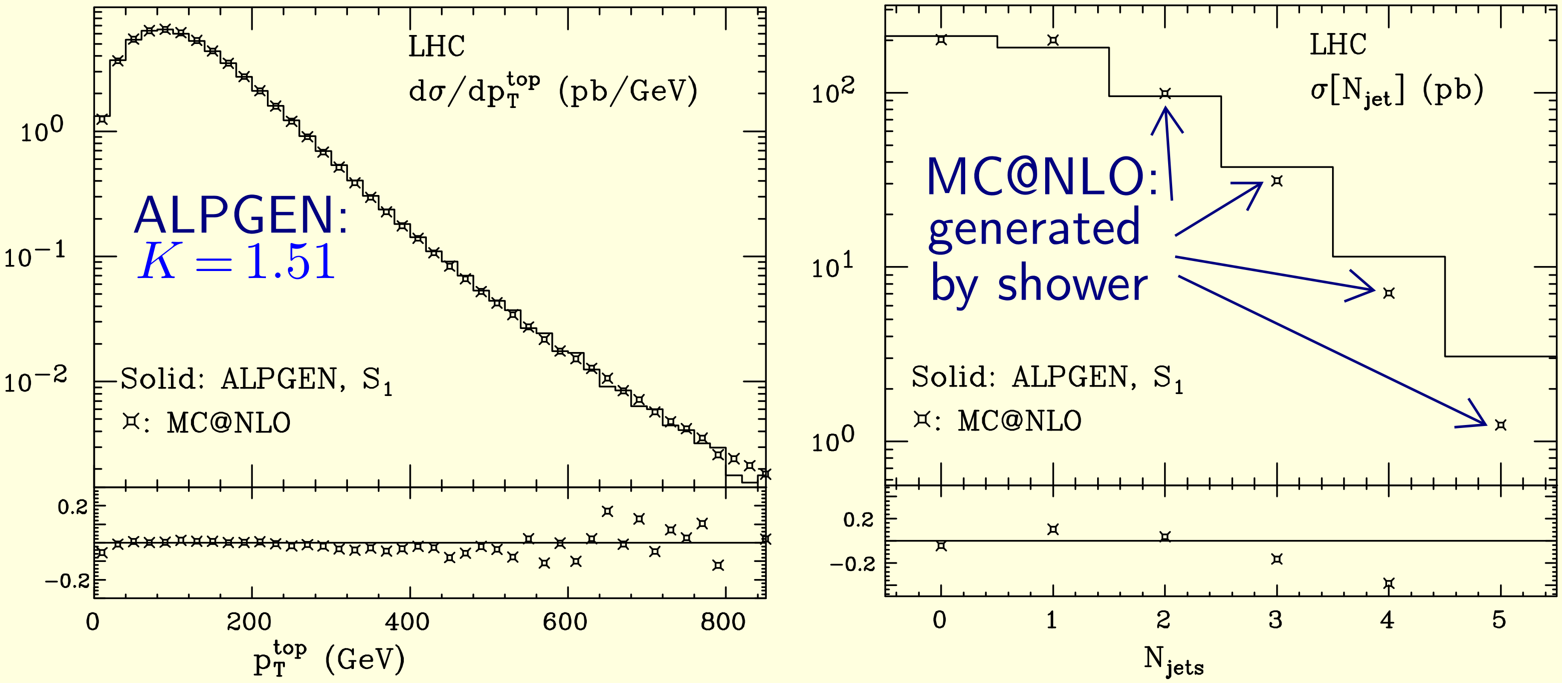
- Very good agreement For large scales (ZZ, $t\bar{t}$ production)
- Differences at small scales ($b\bar{b}$ at the Tevatron)
- POWHEG more reliable in extreme cases like $b\bar{b}, c\bar{c}$ at LHC (yields positive results, MC@NLO has problems with negative weights)

ALPGEN can generate samples of $t\bar{t} + n\,\mathrm{jets}$; can be compared to NLO+PS;

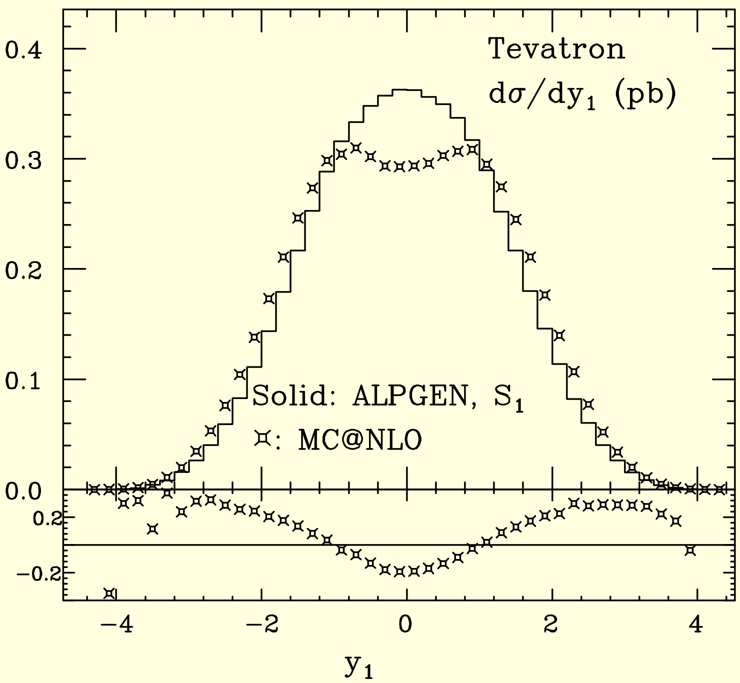
expect:

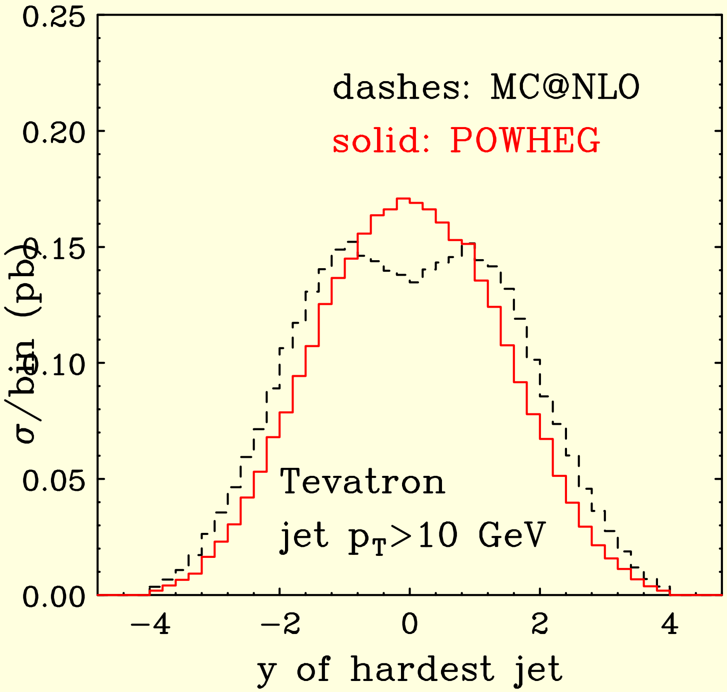
- Disadvantage: worse normalization (no NLO)
- Advantage: better high jet multiplicities (exact ME)

Comparison ALPGEN-MC@NLO carried out in detail
(Mangano, Moretti,Piccinini,Treccani, Nov.06)



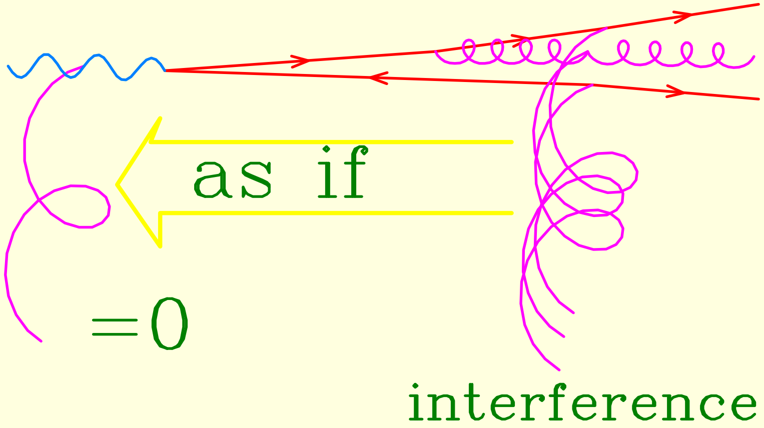
Results as expected but for 1 observable





POWHEG's distribution as in ALPGEN (i.e., no dip);
Notice: size of discrepancy can be attributed to different treatment of higher order terms. Is this "feature" really there?
New $pp \to t\bar{t} + \text{Jet}$ at NLO (Dittmaier, Uwer, Weinzierl) can help to solve this issue: further study needed!

For angular ordered SMC's (i.e. HERWIG):



Angular ordering accounts for soft gluon interference.

Intensity for photon jets $=0$

Intensity for gluon jets $=C_A$

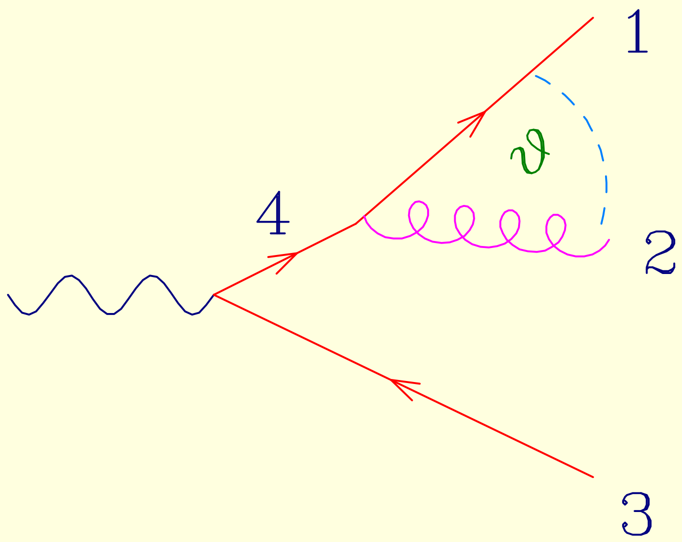
instead of $2C_F + C_A$

Consistent with a boosted jet pair, in the case of a photon jet.
In angular ordered SMC large angle soft emission is generated first.
Hardest emission (i.e. highest p_T) happens later.
Difficult to correct it explicitly.

Recipe for angular ordered showers

- Generate event with harderst emission
- Generate all subsequent emissions with a p_T veto equal to the hardest emission p_T
- Pair up the partons that are nearest in p_T
- Generate an angular ordered shower associated with the paired parton, stopping at the angle of the paired partons (truncated shower)
- Generate all subsequent (vetoed) showers

Example of truncated shower: e^+e^-



Nearby partons: 1,2

Truncated shower: 1,2 pair,

from maximum angle to θ

1 and 2 shower from θ to cutoff

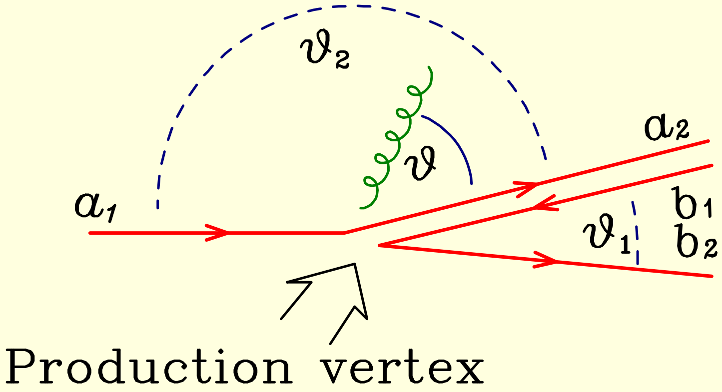
3 showers from maximum to cutoff

The truncated shower reintroduces coherent soft radiation from 1,2 at angles larger than θ (Angular ordered SMC's generate those earlier).

(No evidence of effects from their absence up to now)

Truncated shower are generally needed in angular ordered SMC's

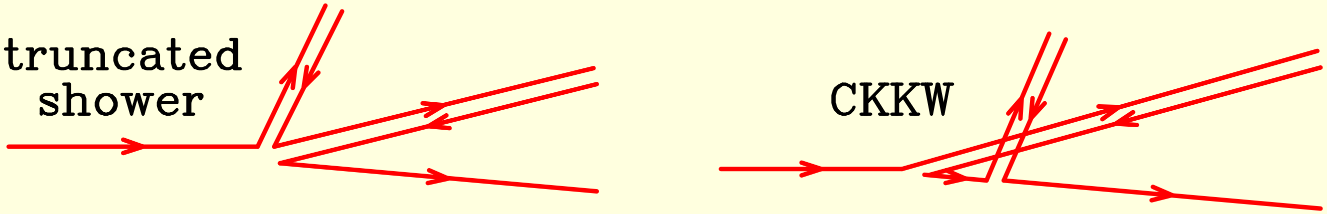
- Every time the shower is initiated by a relatively complex matrix element a truncated shower is needed
- CKKW mocks the effect of truncated shower with a trick (but it misses the correct colour flow)



Consider $e^+e^- \to q\,\bar{q}\,g$.
Assume θ_1 small. Consider gluon emission with angle $\theta \gg \theta_1$, $\theta \ll \theta_2$.
Coherence requires that the emission strength is C_F (gluon and quark coherently)

In HERWIG: initial angle for gluon radiation is θ_1 or θ_2 with a 50% probability. Thus (in the above region) strength is $C_A/2 \approx C_F$ (but only in the average!!)

In CKKW: radiation from gluon restricted to $\theta < \theta_1$, quark radiates with angle up to θ_2. Thus only the quark radiates in the above region, with strength C_F. However, the colour connection is incorrect! Large colour gap in CKKW!



So: coherent showers are always needed when doing ME-Shower matching with angular ordered showers.

Some topics on general formulation of POWHEG

Frixione, Oleari, P.N.

Extension to the general case only a matter of bookkeeping;
POWHEG is fully general, can be applied in any subtraction framework.

We look in details at POWHEG in

- the FKS (Frixione, Kunszt, Signer)
- the CS (Catani, Seymour) subtraction frameworks.

Flavour separation

There are several allowed flavour structures in the n body process. A flavour structure is a flavour assignment to the incoming and outgoing partons.
The B and V contributions are labelled by the flavour structure index f_b.

There are several allowed flavour structures in the $n+1$ body process.
Thus R is labelled by a flavour structure index f_r.
Each component R_{f_r} has several singularity regions. We thus write

$$R = \sum_{\alpha_r} R^{\alpha_r}$$

where each R^{α_r} has a specific flavour structure, and is singular in only one singular region. This partition of R is trivial to perform:

- FKS provides specific kinematic functions S_{α_r}, with $\sum_{\alpha_r} S_{\alpha_r} = 1$ that suppress all but one singular regions.
- in CS one can use instead $S_{\alpha_r} = C_{\alpha_r} / (\sum_{\alpha_r} C_{\alpha_r})$ where C_{α_r} are the dipole subtraction terms.

$\bar{B}$ carries an f_b index;
Sudakov FF also carries an f_b index:

$$\Delta^{f_b}(\Phi_n, p_T) = \exp\left\{ - \sum_{\alpha_r \in \{\alpha_r | f_b\}} \int \frac{\left[d\Phi_r \, R(\Phi_n, \Phi_r) \theta(k_T - p_T) \right]_{\alpha_r}}{B^{f_b}(\Phi_n)} \right\}$$

or

$$\Delta^{f_b}(\Phi_n, p_T) = \prod_{\alpha_r \in \{\alpha_r | f_b\}} \exp\left\{ - \sum \int \frac{\left[d\Phi_r \, R(\Phi_n, \Phi_r) \theta(k_T - p_T) \right]_{\alpha_r}}{B^{f_b}(\Phi_n)} \right\}$$

where

- $\{\alpha_r | f_b\}$ is the set of all singular regions having the underlying Born configuration with flavour structure f_b.
- $[\ldots]_{\alpha_r}$ means that everything inside is relative to the α_r singular term: thus R is R_{α_r}, the parametrization (Φ_n, Φ_r) is the one appropriate to the α_r singular region

The last expression is closer to typical SMC's, with each emission considered independently.

Caveats

- Singularities in B
- Zeros in B

Both cause problems, but they are easily fixed.
For example, zeros in B: further separate

$$R_{\alpha_r} = \frac{k_T^2}{k_T^2 + B} R_{\alpha_r} + \frac{B}{k_T^2 + B} R_{\alpha_r}$$

The first term in non-singular (can be generated directly without Sudakov), while in the second term the zero in B cancels in the Sudakov exponent.

Accuracy of the Sudakov Form Factor

POWHEG's Sudakov FF has the form (with $c \approx 1$)

$$\Delta_t = \exp\left[-\int_t^{Q^2} \frac{dk_T^2}{k_T^2} \frac{\alpha_S(c\, k_T^2)}{\pi} \left\{ A \log \frac{M^2}{k_T^2} + B \right\} \right]$$

We know that the NLL Sudakov form factor has the form

$$\Delta_t^{\rm NLL} = \exp\left[-\int_t^{Q^2} \frac{dk_T^2}{k_T^2} \frac{\alpha_S(k_T^2)}{\pi} \left\{ \left(A_1 + A_2 \frac{\alpha_s(k_T^2)}{\pi} \right) \log \frac{M^2}{k_T^2} + B \right\} \right]$$

provided the colour structure of the process is sufficiently simple ($\leqslant 3$ coloured legs). Can use this to fix c in POWHEG's Sudakov FF. (Suggested in (Catani, Webber, Marchesini, 1991) for HERWIG)

$\geqslant 4$ coloured legs: exponentiation only holds in LL, or LL + (NLL large N_c) if planar colour structures are suitably separated

Summarizing:

POWHEG Sudakov is: always LL accurate,

NLL accurate for $\leqslant 3$ coloured legs, NLL accurate in leading N_c in all cases.

Conclusions and Perspective

- POWHEG is a viable method for interfacing NLO and SMC
- It is easy to implement, does not require new NLO computations
- Does not require committtment to specific SMC implementations
- Its output is closer to traditional SMC's: positive weighted events
- To get it going, we will implement a number of processes: vector bosons and boson pairs, Higgs, Heavy Flavour, etc.
- We collect and publish material to make it easy for others to implement POWHEG with their NLO calculation