PARTON DISTRIBUTIONS FOR THE LHC

STEFANO FORTE UNIVERSITÀ DI MILANO

MCWS, FRASCATI, FEBRUARY 19, 2008

AN ONGOING EFFORT



HERALHC: CERN, MAY 26-30, 2008 (LAST MEETING) PDF4LHC: CERN, FEBRUARY 22-23, 2008 (FIRST MEETING)

PARTONS FOR LHC:

THE ACCURATE COMPUTATION OF PHYSICAL PROCESS AT A HADRON COLLIDER REQUIRES GOOD KNOWLEDGE OF PARTON DISTRIBUTIONS OF THE NUCLEON

FACTORIZATION



IN ORDER TO EXTRACT THE RELEVANT PHYSICS SIGNAL,

WE NEED TO KNOW THE PARTON DISTRIBUTIONS AND THEIR UNCERTAINTY

- IS THIS ASPECT OF LHC PHYSICS UNDER CONTROL?
- WILL LHC TEACH US SOMETHING ABOUT QCD TOO?

SUMMARY

• DETERMINING PDFs

- factorization
- disentangling PDFs
- the standard approach

• CURRENT ISSUES & LHC NEEDS

- old puzzles and current solutions
- the state of the art
- LHC: issues and progress

• OPEN PROBLEMS AND NEW IDEAS

- theoretical issues
- the problem of PDF uncertainties
- the nural monte carlo

DETERMINING PDFs

FACTORIZATION I: DEEP-INELASTIC SCATTERING

STRUCTURE FUNCTIONS...



 $\lambda_l \rightarrow$ lepton helicity

 $\lambda_p \rightarrow \text{proton helicity}$

Lepton fractional energy loss: $y = \frac{p \cdot q}{p \cdot k}$; gauge boson virtuality: $q^2 = -Q^2$ Bjorken x: $x = \frac{Q^2}{2p \cdot q}$ lepton-nucleon CM energy: $s = \frac{Q^2}{xy}$; virtual boson-nucleon CM energy $W^2 = Q^2 \frac{1-x}{x}$;

$$\frac{d^2 \sigma^{\lambda_p \lambda_\ell}(x, y, Q^2)}{dx dy} = \frac{G_F^2}{2\pi (1 + Q^2/m_W^2)^2} \frac{Q^2}{xy} \left\{ \left[-\lambda_\ell y \left(1 - \frac{y}{2} \right) x F_3(x, Q^2) + (1 - y) F_2(x, Q^2) \right] \right\}$$

$$+y^{2}xF_{1}(x,Q^{2})\right]-2\lambda_{p}\left[-\lambda_{\ell}y(2-y)xg_{1}(x,Q^{2})-(1-y)g_{4}(x,Q^{2})-y^{2}xg_{5}(x,Q^{2})\right]$$

	PARITY CONS.	PARITY VIOL.
UNPOL.	F_1, F_2	F_3
POL.	g_1	g_4,g_5

...AND PARTON DISTRIBUTIONS



STRUCTURE FUNCTION=Hard coeff. (partonic structure function) \otimes parton distn.

$$F_2^{\mathrm{NC}}(x,Q^2) = x \sum_{\text{flav. } i} e_i^2(q_i + \bar{q}_i) + \alpha_s \left[C_i[\alpha_s] \otimes (q_i + \bar{q}_i) + C_g[\alpha_s] \otimes g\right]$$

 q_i quark, \bar{q}_i antiquark, g gluon

LEADING PARTON CONTENT (up to $O[\alpha_s]$ corrections)

$$q_{i} \equiv q_{i}^{\uparrow\uparrow} + q_{i}^{\uparrow\downarrow} \qquad \Delta q_{i} \equiv q_{i}^{\uparrow\uparrow} - q_{i}^{\uparrow\downarrow}$$

$$NC \qquad F_{1}^{\gamma, Z} = \sum_{i} e_{i}^{2} (q_{i} + \bar{q}_{i}) \qquad g_{1}^{\gamma, Z} = \sum_{i} e_{i}^{2} (\Delta q_{i} + \Delta \bar{q}_{i})$$

$$CC \qquad F_{1}^{W^{+}} = \bar{u} + d + s + \bar{c} \qquad g_{1}^{W^{+}} = \Delta \bar{u} + \Delta d + \Delta s + \Delta \bar{c}$$

$$CC \qquad -F_{3}^{W^{+}}/2 = \bar{u} - d - s + \bar{c} \qquad g_{5}^{W^{+}} = \Delta \bar{u} - \Delta d - \Delta s + \Delta \bar{c}$$

$$F_{2} = 2xF_{1} \qquad g_{4} = 2xg_{5}$$

 $W^+ \to W^- \Rightarrow u \leftrightarrow d, c \leftrightarrow s;$ more combinations using Isospin: $p \to n \Rightarrow u \leftrightarrow d$

FACTORIZATION II: HADRONIC PROCESSES $\sigma_X(s, M_X^2) = \sum_{a,b} \int_{x_{\min}}^1 dx_1 dx_2 f_{a/h_1}(x_1) f_{b/h_2}(x_2) \hat{\sigma}_{q_a q_b \to X} \left(x_1 x_2 s, M_X^2 \right)$ LEAD. ORD. = $\sigma_0 \sum_{a,b} \int_{\tau}^1 \frac{dx}{x} f_{a/h_1}(x) f_{b/h_2}(\tau/x) \equiv \sigma_0 \mathcal{L}(\tau) \Rightarrow \mathcal{L}$ parton lumi



- Hadronic c.m. energy: $s = (p_1 + p_2)^2$
- Momentum fractions $x_{1,2} = \sqrt{\frac{\hat{s}}{s}} \exp \frac{1}{s}$ Lead. Ord. $\hat{s} = M^2$
- Partonic c.m. energy: $\hat{s} = x_1 x_2 s$
- Invariant mass of final state X (dilepton, Higgs,...): $M_W^2 \Rightarrow$ scale of process

• Scaling variable
$$\tau = \frac{M_X^2}{s}$$

- $\hat{\sigma}_{q_a q_b \to X} = \sigma_0 C\left(x, \alpha_s(M_H^2)\right); \quad C\left(x, \alpha_s(M_H^2)\right) = \delta(1-x) + O(\alpha_s)$
- $\sigma_X(s, M^2) = \sigma_0 \sum_{a,b} \int_{x_{\min}}^1 dx_1 dx_2 f_{a/h_1}(x_1) f_{b/h_2}(x_2) \delta(x_1 x_2 x \tau) C\left(x, \alpha_s(M_H^2)\right)$ = $\sigma_0 \sum_{a,b} \int_{x_2}^1 \frac{dx_1}{x_1} \int_{\tau}^1 \frac{dx_2}{x_2} f_{a/h_1}(x_1) f_{b/h_2}(x_2) C\left(\frac{\tau}{x_1 x_2}, \alpha_s(M_H^2)\right)$

EXAMPLE: DRELL-YAN $\sigma_X \to M^2 \frac{d\sigma}{dM^2}$; $\sigma_0 = \frac{4}{9}\pi \alpha \frac{1}{s}$

LHC: KINEMATICS AND PHYSICAL PROCESSES



PARTON FITS

DATA → PARTON DISTRIBUTIONS

STRATEGY:

- CHOOSE SET OF OBSERVABLES (DIS, DRELL-YAN, W PRODUCTION...) & COM-PUTE THEM IN PERT. THEORY
- CHOOSE A SET OF BASIS PARTON DISTRIBU-TIONS (SINGLET, VALENCE, SEA...)
- FIT THE OBSERVABLES WITH THE PDFS AS FREE PARAMETERS

DATA INCLUDED IN CTEQ5 PARTON FIT



TASKS:

- STRUCTURE FUNCTION (OR XSECT) IS A CON-VOLUTION OVER x OF PARTON DISTNS. AND PERTURBATIVE CROSS SECTION \rightarrow MUST DECONVOLUTE
- EACH STRUCTURE FUNCTION (OR XSECT) IS A LINEAR COMBINATION OF MANY PARTON DISTNS ($2N_f$ QUARKS + 1 GLUON)
 - \rightarrow MUST COMBINE DIFFERENT PROCESSES
- DATA GIVEN AT VARIOUS SCALES, WANT PAR-TON DISTNS. AS FCTN OF x AT COMMON SCALE Q^2
 - \rightarrow MUST EVOLVE
- TH UNCERTAINTIES: HIGHER ORDERS, RE-SUMMATIONS, HEAVY QUARK THRESHOLDS, NUCLEAR CORRECTIONS, HIGHER TWIST,

DISENTANGLING QUARKS FROM ANTIQUARKS

 γ^* DIS only measures $q + \bar{q}$ combination!

DRELL-YAN p/d ASYMMETRY



DISENTANGLING STRANGENESS γ^* SCATTERING VS. W^{\pm} SCATTERING:

IN NC, CHARGED LEPTON DIS, ONLY MEASURE COMBINATION $\sum_{i} e_i^2 (q_i + \bar{q}_i)$

- ONLY C-EVEN $q_i + \bar{q}_i$
- ONLY FIXED COMBINATION $\frac{4}{9}(u+\bar{u}) + \frac{1}{9}(d+\bar{d}) + \frac{1}{9}(s+\bar{s})$

IN NEUTRINO DIS, CAN DISENTANGLE INDIVIDUAL PDFS BY LINEAR COMBINATION: AT LO





ERRORS ON PDFS AT A NUFACT

COMPARED TO A PURE DIS FIT

DETERMINING THE GLUON

EVOLUTION:

SINGLET SCALING VIOLATIONS

 $\frac{d}{dt}F_2^s(N,Q^2) = \frac{\alpha_s(Q^2)}{2\pi} \left[\gamma_{qq}(N)F_2^s + 2n_f\gamma_{qg}(N)g(N,Q^2) \right] + O(\alpha_s^2)$ $F_2(N,Q^2) \equiv \int_0^1 dx \, x^{N-1}F_2(x,Q^2); \qquad \gamma_{ij}(N) \equiv \int_0^1 dx \, x^{N-1}P_{ij}(x,Q^2)$ $LARGE/SMALL X \Leftrightarrow LARGE/SMALL N$

AT LARGE x



THE STANDARD APPROACH: FUNCTIONAL PARTON FITTING

• CHOOSE A FIXED FUNCTIONAL FORM:

- MRST: 24 parms., some fixed \rightarrow 15 parms.

$$xq(x,Q_0^2) = A(1-x)^{\eta}(1+\epsilon x^{0.5}+\gamma x)x^{\delta}, \quad x[\bar{u}-\bar{d}](x,Q_0^2) = A(1-x)^{\eta}(1+\gamma x+\delta x^2)x^{\delta}.$$

$$xg(x,Q_0^2) = A_g(1-x)^{\eta_g} (1+\epsilon_g x^{0.5} + \gamma_g x) x^{\delta_g} - A_-(1-x)^{\eta_-} x^{-\delta_-}$$

- CTEQ: 20 PARMS.

$$x f(x, Q_0) = A_0 x^{A_1} (1-x)^{A_2} e^{A_3 x} (1+e^{A_4} x)^{A_5}$$

with independent params for combinations $u_v \equiv u - \bar{u}$, $d_v \equiv d - \bar{d}$, g, and $\bar{u} + \bar{d}$, $s = \bar{s} = 0.2 (\bar{u} + \bar{d})$ at Q_0 ; NORM. FIXED BY SUM RULES

- Alekhin: 17 parms.

$$xu_{\rm V}(x,Q_0) = \frac{2}{N_{\rm u}^{\rm V}} x^{a_{\rm u}} (1-x)^{b_{\rm u}} (1+\gamma_2^{\rm u} x); \quad xu_{\rm S}(x,Q_0) = \frac{A_{\rm S}}{N_{\rm S}} \eta_{\rm u} x^{a_{\rm S}} (1-x)^{b_{\rm S} {\rm u}}$$

$$xd_{\rm V}(x,Q_0) = \frac{1}{N_{\rm d}^{\rm V}} x^{a_{\rm d}} (1-x)^{b_{\rm d}}; \quad xd_{\rm S}(x,Q_0) = \frac{A_{\rm S}}{N^{\rm S}} x^{a_{\rm S}} (1-x)^{b_{\rm S}},$$

$$xs_{\rm S}(x,Q_0) = \frac{A_{\rm S}}{N^{\rm S}} \eta_{\rm s} x^{a_{\rm S}} (1-x)^{(b_{\rm Su}+b_{\rm Sd})/2}; \quad xG(x,Q_0) = A_{\rm G} x^{a_{\rm G}} (1-x)^{b_{\rm G}} (1+\gamma_1^{\rm G}\sqrt{x}+\gamma_2^{\rm G}x),$$

- EVOLVE TO DESIRED SCALE & COMPUTE PHYSICAL OBSERVABLES
- DETERMINE BEST-FIT VALUES OF PARAMETERS
- DETERMINE ERROR BY PROPAGATION OF ERROR ON PARMS ('HESSIAN METHOD') OR BY PARM. SCANS ('LAGRANGE MULTIPLIER METHOD')

HOW WELL DOES IT WORK?



HOW WELL DOES IT WORK? DIS+DY ONLY (Alekhin 2003-2006)



HOW WELL DOES IT WORK? GLOBAL FITS (MRST-CTEQ 2002-2006)

CTEQ ERROR BAND & MRST/CTEQ CURVE



MRST GLUON ERROR BAND &CTEQ/MRST CURVE

Uncertainty of gluon from Hessian method



-FEW PERCENT ERROR ON VALENCE & GLUE -OTHER PDFS: ERROR NOT WELL CONTROLLED KNOWN ISSUES & LHC NEEDS

CASE STUDY I: THE CDF LARGE E_T JETS CDF 1995

- DISCREPANCY BETWEEN QCD CALCULATION AND CDF JET DATA (1995)
- EVIDENCE FOR QUARK COMPOSITENESS?
- BUT NO INFO ON PARTON UNCERTAINTY \Rightarrow RESULT STRONGLY DEPENDS ON GLUON AT $x \ge 0.1$



DISCREPANCY REMOVED IF JET DATA INCLUDED IN THE FIT NEW CTEQ FIT (1996)







$\begin{array}{l} \textbf{CASE STUDY II: THE NUTEV ANOMALY} \\ \textbf{THE PASCHOS-WOLFENSTEIN RATIO: DATA...} \\ \textbf{NuTeV 2001} \quad \sin^2 \theta_W(\text{OS}) = 0.2272 \pm 0.0013(\text{stat}) \pm 0.0009(\text{syst}) \pm 0.0002(M_t, M_H) \end{array}$

NuTeV 2001 $\sin^2 \theta_W(OS) = 0.2272 \pm 0.0013(\text{stat}) \pm 0.0009(\text{syst}) \pm 0.0002(M_t, M_H)$ Global Fit 2003 $\sin^2 \theta_W(OS) = 0.2229 \pm 0.0004$

...VS. THEORY

$$R^{-} = \frac{\sigma_{NC}(\nu) - \sigma_{NC}(\bar{\nu})}{\sigma_{CC}(\nu) - \sigma_{CC}(\bar{\nu})}$$

= $\left(\frac{1}{2} - \sin^{2}\theta_{W}\right) + 2\left[\frac{(u - \bar{u}) - (d - \bar{d})}{u - \bar{u} + d - \bar{d}} - \frac{s - \bar{s}}{u - \bar{u} + d - \bar{d}}\right] \times \left[\left(\frac{1}{2} - \frac{7}{6}\sin^{2}\theta_{W}\right) + \frac{4}{9}\frac{\alpha_{s}}{2\pi}\left(\frac{1}{2} - \sin^{2}\theta_{W}\right) + O(\alpha_{s}^{2})\right] + O(\delta(u - d)^{2}, \,\delta s^{2})$

u,d...denote momentum fractions carried by corresp. quark flavors NUTEV RESULT OBTAINED NEGLECTING:

- ISOSPIN VIOLATION \rightarrow ISOSPIN KNOWN TO BE GOOD TO $\sim .1\%$
- STRANGE ASYM. \rightarrow EXPECT SEA TO BE FLAVOUR/ANTIFLAVOUR SYMMETRIC
- QCD CORRECTIONS → TINY (ONLY ENTER THROUGH SYM. VIOLATING TERMS)

ISOSPIN VIOLATION

QED EFFECTS LEAD TO ISOSPIN VIOLATION: $u - \bar{u}$ radiate more photons than $d - \bar{d}$: $\frac{d}{dt}q_i \propto e_i^2 q_i$ \Rightarrow MORE PHOTON MOMENTUM IN PROTON THAN NEUTRON

 $\Rightarrow |u(x) - \bar{u}(x)| < |d(x) - \bar{d}(x)| \text{ at large } x$



- SIGN OF EFFECT AS REQUIRED TO EXPLAIN NUTEV
- SIZE OF EFFECTS WITH REASON-ABLE ASSUMPTIONS ABOUT 1/2 OF NUTEV ANOMALY
- THEORETICAL RESULTS AGREES WITH FIT IF ISOSPIN VIOLATION AL-LOWED

MRST 2005: "QED" PARTON SET

STRANGENESS ASYMMETRY

Q: ARE WE SURE THAT MOMENTUM FRACTION $s - \bar{s} = 0$?

A: MEASURE IT!: CHARM IS COPIOUSLY PRODUCED IN $W^+ + s \rightarrow c$

easily tagged through dimuon signal, 2nd muon from subsequent c decay

 \Rightarrow ACCURATE EXTRACTION OF THE STRANGE DISTRIBUTION $Q^2=20GeV^2$



CCFR/NUTEV $s - \bar{s}$ determination

5000 ν & 1500 $\bar{\nu}$ DIMUON EVENT SAMPLE: ASSUMED PARM.: $s(x) = \kappa \frac{\bar{u}(x) + \bar{d}(x)}{2} (1-x)^{\alpha}$ NEGATIVE $s - \bar{s}$ AT SMALL x \Rightarrow MOM. FRACT. $s - \bar{s} = -0.003 \pm 0.001$ NUTEV ANOMALY WORSE!

HOWEVER, BPZ GLOBAL FIT TO NEUTRINO INCLUSIVE DIS (Barone et al 2003) \Rightarrow POSITIVE (TINY) ASYMMETRY

COMBINING INCLUSIVE AND EXCLUSIVE INFORMATION

CTEQ DEDICATED DIMUON ANALYSIS

- $\int_0^1 (s(x) \bar{s}(x)) dx = 0$ in proton \Rightarrow Either $s(x) - \bar{s}(x)$ has a node or it vanishes everywhere
- $[s(x) \bar{s}(x)] < 0$ for small $x \lesssim 0.05$ constrained by dimuon
- LARGE x REGION WEIGHS MORE IN MOMENTUM FRACTION
- POSITIVE MOM. FRACTION $s \bar{s} \approx 0.02$:



STRANGE QUARK PDF FITTED IN FORTHCOMING MRST & CTEQ SETS

CASE STUDY III W PRODUCTION @ LHC: THE GOOD NEWS



- W/Z RAPDITY SPECTRA & TOTAL CROSS SECTIONS: $\sim 15\%$ PRE-HERA ACCURACY $\sim 3 - 5\%$ POST-HERA ACCURACY
- GOOD AGREEMENT BETWEN DIFFERENT PDF SETS

PDF Set	$\sigma(W^+).B(W^+ \to l^+\nu_l)$	$\sigma(W^-).B(W^- \to l^- \bar{\nu}_l)$	$\sigma(Z).B(Z \to l^+l)$
ZEUS-S NO HERA	$10.63 \pm 1.73~\mathrm{NB}$	$7.80 \pm 1.18~\mathrm{NB}$	$1.69\pm0.23~\mathrm{NB}$
ZEUS-S	$12.07\pm0.41~\mathrm{NB}$	$8.76\pm0.30~\mathrm{NB}$	$1.89\pm0.06~\mathrm{nb}$
CTEQ6.1	$11.66\pm0.56~\mathrm{NB}$	$8.58\pm0.43~\mathrm{NB}$	$1.92\pm0.08~\mathrm{nb}$
MRST01	$11.72\pm0.23~\mathrm{NB}$	$8.72\pm0.16~\mathrm{NB}$	$1.96\pm0.03~\mathrm{nb}$

W PRODUCTION @ LHC: THE NOT SO GOOD NEWS



- NEW (CTEQ6.5) PARTON SET IN-CLUDES HQ MATCHING
- EFFECT OF IMPROVED HW MASS FELT MOSTLY IN SMALL x QUARK SUPPRESSION OF CHARM \Rightarrow EN-HANCEMENT OF LIGHT SEA
- IN COMPARISON TO PREVIOUS (CTEQ76.1), SIGNIFICANT CHANGE OF u, d QUARK DISTNS AT $x \sim 0.01$
- W, Z TOTAL XSECT NO LONGER AGREES WITH MRST THOUGH MRST INCLUDES HQ MATCHING!
- EFFECT OF INTRINSIC CHARM (IC) MINOR

WHERE DO WE STAND NOW? WHAT WE HAVE LEARNT

- LIGHT QUARK STRUCTURE IN "VALENCE" REGION $0.1 \le x \le 0.5$ (old fixed target dis data)
- SINGLET AND GLUON AT SMALL $x < 10^{-2}$ (HERA)
- SEA ASYMMETRY AT MEDIUM $x \sim 0.1 \div 0.2$ (Drell-Yan)
- HINTS ON STRANGENESS (neutrinos)

WHAT WE ARE STILL MISSING

- GLUONS AT LARGE x (cfr large E_T jet problem)
- NONSINGLET & VALENCE AT SMALL x
- DETAILED INFO ON STRANGENESS (cfr NuTeV problem)
- INFO ON HEAVY QUARKS (cfr small x W xsect problem)

IS IT A PROBLEM?

EXAMPLE: LACK OF KNOWLEDGE OF LARGE x GLUON LIMITS DISCOVERY POTENTIAL FOR EXTRA DIMENSIONS

UPPER LIMIT ON COMPACTIFICATION SCALE FROM DIJET CROSS SECTIONS FROM 100 FB^{-1} AT LHC Ferrag (ATLAS, 2006)

	2	4	6	
	extra dimensions	extra dimensions	extra dimensions	
THEORETICALLY	5 TeV	5 TeV	5 TeV	
INCLUDING PDF UNCERTAINTIES	$< 2~{ m TeV}$	< 3 TeV	< 4 TeV	

CROSS-SECTION IN FIXED p_t BINS

10

10

10

10

10

10

10

10

-14



quark+gluon

GG CHANNEL





SOLUTIONS: LARGE E_T JETS @ LHC DETERMINING THE GLUON AT LARGE x

UNCERTAINTY IN THE GLUON GREATLY REDUCED PROVIDED SYSTEMATICS CAN BE KEPT AT FEW PERCENT LEVEL



D. Clements (Atlas 2006)

SOLUTIONS: W ASYMMMETRY @ LHC DETERMINING QUARKS AT SMALL x

- W production at LHC probes $x \sim 10^{-2}$
- W^{\pm} asymmetries sensitive to \bar{u}/\bar{d}
- \Rightarrow IF SMALL *x* BEHAVIOUR IS NOT AS CURRENTLY ASSUMED ("REGGE"), W^{\pm} ASYMMETRY CHANGES BY UP TO FACTOR 5!



E. Perez (CMS 2006)

SOLUTIONS: W DISTRIBUTION @LHC PRECISION PHYSICS @ SMALL x

- MRST03 \Rightarrow BEST FIT VS. MRST02 \Rightarrow VERY SMALL *x* HERA DATA NOT INCLUDED
- DIFFERENCE IN ASYMETRY SEEN AFTER FEW HOURS OF RUNNING!



A. Cooper-Sarkar (Atlas 2006)

CURRENT ISSUES AND NEW IDEAS

NNLO CORRECTIONS

- NNLO SPLITTING FUNCTIONS KNOWN (Moch, Vermaseren and Vogt, 2004)
- NNLO HARD XSECTS AVAILABLE FOR DIS, DY, W AND HIGGS PRODUCTION (INCL.)



• ALEKHIN, MRST NNLO FITS AVAILABLE

- EFFECT OF NNLO CORRECTIONS AROUND 5-10 %
- MUCH LARGER IN SPECIFIC KIN. REGIONS (SMALL x, LARGE x)
- SOME OPEN TH. ISSUES (HQ MATCHING...)

HEAVY QUARKS

HOW CAN ONE ACCOUNT FOR HEAVY FLAVOURS (CHARM, BEAUTY...)? SIMPLE OPTION: (CTEQ6, Alekhin) CHARM PDF VANISHES BELOW THRESHOLD, INCLUDED ALONG OTHER PDFS ABOVE THRESHOLD \Rightarrow EFFECTIVELY, $m_c \approx 0$ FOR $Q^2 > Q_{th}^2$.

- HQ PDF GENERATED DYNAMICALLY BY PERTURBATIVE EVOLUTION (HQ PAIR-PRODUCED BY RADIATION FROM GLUONS)
- TREATMENT NOT ACCURATE IN $Q^2 \approx Q_{th}^2$ REGION

MORE REFINED TREATMENT OF THRESHOLD: [Collins, Tung et al (ACOT) 1986-2006]



- INCLUDED IN LATEST MRST & CTEQ RELEASES
- CAN HAVE UP TO 10% EFFECT ON SMALL x light sea
- NO CONSENSUS ON MATCHING PRESCRIPTION (NO NNLO)

THRESHOLD (LARGE x) RESUMMATION

AT $O(\alpha_s^n)$, $O\left[ln^{2m}(1-x)\right]$ CONTRIBUTIONS:

- \Rightarrow PERT. TH. UNSTABLE AT LARGE x (C.M. ENERGY \sim FINAL STATE MASS)
 - DIS: SIZABLE ONLY @ VERY LARGE x, WHERE XSECT & PDF TINY (Corcella, Magnea 2005)
 - DY: CAN HAVE SIZABLE EFFECTS, ESPECIALLY ON RAP. DISTN. (CMNT, 1994; Bolzoni 2006)
 - NOT INCLUDED IN CURRENT FITS



HIGH ENERGY (SMALL x) RESUMMATION

AT $O(\alpha_s^n)$, $O\left[ln(\frac{1}{x})^n\right]$ CONTRIBUTIONS:

 \Rightarrow PERT. TH. UNSTABLE AT SMALL x (C.M. ENERGY >> FINAL STATE MASS)

x_{cut} :	0	0.0002	0.001	0.0025	0.005	0.01
# DATA POINTS	2097	2050	1961	1898	1826	1762
$\chi^2(x > 0)$	2267					
$\chi^2(x > 0.0002)$	2212	2203				
$\chi^2(x > 0.001)$	2134	2128	2119			
$\chi^2(x > 0.0025)$	2069	2064	2055	2040		
$\chi^2(x > 0.005)$	2024	2019	2012	1993	1973	
$\chi^2(x > 0.01)$	1965	1961	1953	1934	1917	1916
Δ_i^{i+1}	0.	19 0.10	0.24	0.28	0.02	

DATA-THEORY AGREEMENT FOR EVOLUTION OF F_2 IMPROVES IF SMALL x DATA REMOVED (MRST 2003) χ^2 improves with fixed # of pts (same row)

- CONSIDERABLE PROGRESS IN FULL RESUMMA-TION OF SMALL x SPLITTING FUNCTIONS IN GLUON SECTOR (Ciafaloni, Colferai, Salam, Staśto; Altarelli, Ball, S.F.)
 1.2
- FULLY RESUMMED RESULTS AVAILABLE FOR DIS (ABF 2008)
- RESUMMED PERTURBATIVE EXPANSION STABLE
- RESUMMATION AS LARGE AS NNLO, OPPOSITE SIGN AT $x \sim 10^{-4}$
- NOT YET INCLUDED IN PARTON FITS





PDF UNCERTAINTIES: WHAT'S THE PROBLEM?

- FOR A SINGLE QUANTITY, WE QUOTE 1 SIGMA ERRORS: VALUE \pm ERROR
- FOR A PAIR OF NUMBERS, WE QUOTE A 1 SIGMA ELLIPSE
- FOR A FUNCTION, WE NEED AN "ERROR BAR" IN A SPACE OF FUNCTIONS

MUST DETERMINE THE PROBABILITY DENSITY (MEASURE) $\mathcal{P}[f_i(x)]$ IN THE SPACE OF PARTON DISTRIBUTION FUNCTIONS $f_i(x)$ (*i*=quark, antiquark, gluon)

EXPECTATION VALUE OF $\sigma[f_i(x)] \Rightarrow$ FUNCTIONAL INTEGRAL

$$\left\langle \sigma\left[f_{i}(x)\right]\right\rangle = \int \mathcal{D}f_{i}\,\sigma\left[f_{i}(x)\right]\,\mathcal{P}[f_{i}],$$

MUST DETERMINE AN INFINITE–DIMENSIONAL OBJECT FROM A FINITE SET OF DATA POINTS


- ALEKHIN VS. MRST/CTEQ \rightarrow predictions for associate Higgs W production @ LHC do not agree within respective errors
- MRST vs. CTEQ GLUONS DO NOT AGREE WITHIN RESPECTIVE ERRORS

ARE MORE DATA ENOUGH TO RESOLVE THE DISCREPANCIES?

INCOMPATIBLE DATA?

E866 DY DATA DISAGREE WITH DIS DATA:

 $\sigma_{DY} \sim q(x_1)q(x_2)$ disagrees with DIS quark at same x and Q^2 σ_W PREDICTION UNSTABLE

ONE σ ERROR BAND FOR PHYSICAL PREDICTIONS BASED ON DIFFERENT UNDERLYING DATASETS DISAGREE





CTEQ 2004

ALEKHIN 2005

CONSERVATIVE SOLUTION:

SELECT COHERENT SET OF DATA: ALEKHIN PARTONS

- ONLY DIS + SUBSET OF DY DATA INCLUDED
- $\Delta \chi^2 = 1$ provides good 1- σ curves



- ALEKHIN 2003-2006 PARTON UNCERTAINTIES COMPARABLE TO CTEQ6
- ERROR ON σ_W COMPARABLE TO CTEQ, MRST

CONSERVATIVE SOLUTION:

SELECT COHERENT SET OF DATA: ALEKHIN PARTONS

- ONLY DIS + SUBSET OF DY DATA INCLUDED
- $\Delta \chi^2 = 1$ provides good 1- σ curves



- ALEKHIN 2003-2006 PARTON UNCERTAINTIES COMPARABLE TO CTEQ6
- ERROR ON σ_W COMPARABLE TO CTEQ, MRST
- LOTS OF MISSING INFORMATION (E.G. LARGE x GLUON)

STANDARD SOLUTION: CTEQ TOLERANCE CRITERION

- Determine eigenvectors of χ^2 paraboloid
- DETERMINE 90% C.L. FOR EACH EXPT. ALONG EACH EIGENVECTOR
- DETERMINE MOST RESTRICTIVE INTERVAL ABOUT GLOBAL MINIMUM (TOLERANCE)

TOLERANCE PLOT FOR 4TH EIGENVEC.





MINIMUM χ_i^2 VS GLOBAL χ^2

STANDARD SOLUTION: CTEQ TOLERANCE CRITERION

- DETERMINE EIGENVECTORS OF χ^2 PARABOLOID
- DETERMINE 90% C.L. FOR EACH EXPT. ALONG EACH EIGENVECTOR
- DETERMINE MOST RESTRICTIVE INTERVAL ABOUT GLOBAL MINIMUM (TOLERANCE)

 $\Delta \chi^2 = 100$



MINIMUM χ^2_i VS GLOBAL χ^2



(CTEQ6, 2002-2007)

• STUDY CORRELATION BETWEEN DATA & PDFS

- REMOVE TROUBLESOME DATA BY CUTTING LOW x, LOW Q^2 ("CONSERVATIVE PARTONS", MRST 2003)
- RESULTS UNSTABLE \rightarrow MISSING INFO (MRST)
- OR STABLE WITH "PROPER" ASSUMPTIONS (CTEQ)





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- RESULTS UNSTABLE \rightarrow MISSING INFO (MRST)
- OR STABLE WITH "PROPER" ASSUMPTIONS (CTEQ)
- IS A STABLE RESULT RELIABLE?





EFFECT OF HQ MATCHING



THE HERA-LHC BENCHMARK: AN IMPASSE

- HERA-LHC BENCHMARK PARTONS OBTAINED FROM NC DIS DATA ONLY, $Q^2 > 9 \ {\rm GeV}^2$
- FITTED WITH RESPECTIVE METHODS BY A & MRST
- ALL ERRORS DETERMINED AT ONE σ



GLUON AND d_V

MRST VS. BENCH



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GLUON AND d_V

- IT IS UNSURPRIZING THAT CENTRAL VALUES DEPEND STRONGLY ON THE DATASET
- BUT IT IS VERY WORRISOME THAT THE RESULT WITH THE FULL DATA SET IS NOT WITHIN THE ERROR BAND OF THE RESULT FROM A DATA SUBSET

THE NEURAL MONTE CARLO THE NNPDF COLLABORATION

(2004: Del Debbio, SF, Latorre, Piccione, Rojo; 2007: +Ball, Guffanti, Ubiali)

BASIC IDEA: USE NEURAL NETWORKS AS UNIVERSAL UNBIASED INTERPOLANTS



THE PROJECT AND ITS STRUCTURE OF THE CODE

- ABOUT 20000 LINES OF CODE, ABOUT 200 MODULES/ROUTINES
- OBJECT-ORIENTED STRUCTURE, SVN
- FULL DOCUMENTATION AVAILABLE AT

http://sophia.ecm.ub.es/nnpdf/







IN A STANDARD FIT, ONE LOOKS FOR MINIMUM χ^2 WITH GIVEN FINITE PARM.

- IF THE BASIS IS TOO LARGE, THE FIT NEVER CONVERGES
- IF THE BASIS IS TOO SMALL, THE FIT IS BIASED

Q: HOW CAN ONE BE SURE THAT THE COMPROMISE IS UNBIASED? IN A NEURAL FIT, SMOOTHNESS DECREASES AS FIT QUALITY IMPROVES:



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PROPER LEARNING

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OVERLEARNING

A: STOP THE FIT BEFORE OVERLEARNING SETS IN!

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OVERLEARNING

A: STOP THE FIT BEFORE OVERLEARNING SETS IN! COULD BE DONE WITH STANDARD PARAMETRIZATIONS, BUT VERY INEFFICIENTLY

MINIMIZE BY GENETIC ALGORITHM: AT EACH GENERATION, THE χ^2 EITHER UNCHANGED OR DECREASING

- DIVIDE THE DATA IN TWO SETS: TRAINING AND VALIDATION
- MINIMIZE THE χ^2 OF THE DATA IN THE TRAINING SET
- AT EACH ITERATION, COMPUTE THE χ^2 FOR THE DATA IN THE VALIDATION SET (NOT USED FOR FITTING)
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GO!

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TOO LATE!

STATUS OF THE PROJECT NONSINGLET FIT

- FIT TO $F_2^p F_2^d$ DIS data from BCDMS & NMC (about 500 datapoints)
- DETERMINATION OF $u + \bar{u} (d + \bar{d})$ (ISOTRIPLET) PARTON DISTRIBUTION
- PUBLISHED IN JHEP 0703:039,2007

FULL FIT

- FIT TO AVAILABLE UNPOLARIZED DIS DATA: ELECTRON AND NEUTRINO BEAMS, NC & CC SCATTERING, PROTON AND DEUTERIUM TARGETS FROM 12 EXPERIMENTS (ABOUT 3000 DATAPOINTS)
- DETERMINATION OF A SET OF FIVE INDEP. PDFS (SAME AS MRST, CTEQ)
- TO APPEAR BEFORE THE SPRING 2008

RESULTS & COMPARISON TO OTHER APPROACHES (NONSINGLET FIT)

NLO RESULTS: THE STRUCTURE FUNCTION $F_2^{NS}(x, Q^2)$ vs x at $Q^2 = 15 \text{ GeV}^2$ vs Q^2 at x = 0.15



- COMPATIBLE WITH EXISTING FITS WITHIN ERROR (even when they disagee with each other)
- UNCERTAINTY MUCH LARGER IN EXTRAPOLATION BUT ALSO IN DATA REGION (note no other global fit data constrain $q_{\rm NS}$)
- CENTRAL FIT DISAGREES WITH EXISTING FITS IN VALENCE REGION $0.1 \leq x \leq 0.3$

RESULTS (SINGLET FIT) PRELIMINARY: 25 REPLICAS



CONCLUSION

AT LHC, WE NEED PRECISION PHYSICS

FOR DISCOVERY PHYSICS



Higgs decay in $e^+e^- + 2$ jets at CMS



THE BAYESIAN MONTE CARLO (GIELE, KOSOWER, KELLER 2001)

- generate a Monte-Carlo sample of fcts. with "reasonable" prior distn. (e.g. an available parton set) \rightarrow representation of probability functional $\mathcal{P}[f_i]$
- calculate observables with functional integral
- update probability using Bayesian inference on MC sample: better agreement with data \rightarrow more functions in sample
- iterate until convergence achieved

PROBLEM IS MADE FINITE-DIMENSIONAL BY THE CHOICE OF PRIOR, BUT RESULT DO NOT DEPEND ON THE CHOICE IF SUFFICIENTLY GENERAL HARD TO HANDLE "FLAT DIRECTIONS" (Monte Carlo replicas which lead to same agreement with data); COMPUTATIONALLY VERY INTENSIVE; DIFFICULT TO ACHIEVE INDEP. FROM PRIOR

RESULT: FERMI PARTONS



 F_2^{singlet} and gluon ratios fermi/MRST

ONLY SUBSET OF DATA FITTED (H1, E665, BCDMS DIS DATA)

GOOD AGREEMENT WITH TEVATRON W XSECT TROUBLE WITH VALUE OF α_s

WHAT ARE NEURAL NETWORKS?



MULTILAYER FEED-FORWARD NETWORKS

- Each neuron receives input from neurons in preceding layer and feeds output to neurons in subsequent layer
- Activation determined by weights and thresholds

$$\xi_i = g\left(\sum_j \omega_{ij}\xi_j - \theta_i\right)$$

• Sigmoid activation function $g(x) = \frac{1}{1 + e^{-\beta x}}$

JUST ANOTHER SET OF BASIS FUNCTIONS!

A 1-2-1 NN:
$$\xi_1^{(3)}(\xi_1^{(1)}) = \frac{1}{\substack{\theta_1^{(3)} - \frac{\omega_{11}^{(2)}}{1+e^{\theta_1^{(2)} - \xi_1^{(1)}\omega_{11}^{(1)}} - \frac{\omega_{12}^{(2)}}{1+e^{\theta_2^{(2)} - \xi_1^{(1)}\omega_{21}^{(1)}}}}$$

ANY FUNCTION CAN BE REPRESENTED BY A SUFFICIENTLY BIG NEURAL NETWORK LESS PARAMETERS \rightarrow SMOOTHER FUNCTIONS

MONTE CARLO DATA GENERATION

- BCDMS+ NMC PROTON & DEUTERON F_2 DATA (FULL CORRELATED SYSTEMATICS AVAILABLE), TAKEN AT 4 BEAM ENERGIES
- ON TOP OF STAT. ERRORS, 4 SYSTEMATICS + 1 NORMALIZATION (NMC) OR 6 SYSTEMATICS + 1 ABSOLUTE & 2 RELATIVE NORMALIZATIONS (BCDMS), WITH VARIOUS FORMS OF CORRELATION (FULL, OR FOR EACH TARGET, OR FOR EACH BEAM ENERGY)

GENERATE DATA ACCORDING TO A MULTIGAUSSIAN DISTRIBUTION

$$F_{i}^{(art)(k)} = (1 + r_{5}^{(k)} \sigma_{N}) \sqrt{1 + r_{i,6}^{(k)} \sigma_{N_{t}}} \sqrt{1 + r_{i,7}^{(k)} \sigma_{N_{b}}} \left[F_{i}^{(exp)} + \frac{r_{i,1}^{(k)} f_{b} + r_{i,2}^{(k)} f_{i,s} + r_{i,3}^{(k)} f_{i,r}}{100} F_{i}^{(exp)} + r_{i,s}^{(k)} \sigma_{s}^{i} \right]$$

$$r \text{ univariate gaussian random nos., one } r_{i,s} \text{ for each data, but single } r_{i,j} \text{ for all correlated data}$$

$$r_{\text{Central values}} = F_{\text{Froton}} = C_{\text{Correlations}}$$



SCATTER PLOT ART. VS. EXP. FOR 10 (RED) 100 (GREEN) AND 1000 (BLUE) REPLICAS

NEED 1000 REPLICAS TO REPRODUCE CORRELATIONS TO PERCENT ACCURACY

PERTURBATIVE EVOLUTION

- PARAMETRIZE INITIAL PDFS AS A FUNCTION OF \boldsymbol{x}
- DETERMINE GREEN'S FUNCTION FOR ALTARELLI-PARISI EVOLUTION $\Gamma(x, \alpha_s (Q^2), \alpha_s (Q_0^2))$ (note it is a distribution)
- DETERMINE EVOLVED PDF AS $q(x,Q^2) = Gq(x,Q_0^2) + \int_x^1 \frac{dy}{y} \Gamma^{(+)}(y,\alpha_s(Q^2),\alpha_s(Q_0^2)) q\left(\frac{x}{y},Q_0^2\right)$
- GREEN FUNCTION CAN BE INTERPOLATED OR COMPUTED ON A GRID AND STORED
- EVOLUTION AND INTERPOLATION FULLY BENCHMARKED

TRAINING...

- EACH NEURAL NET IS FITTED TO A PSEUDODATA REPLICA BY MINIMIZING ITS χ^2
- MINIMIZATION THROUGH GENETIC ALGORITHM + REWEIGHTING OF EXPERIMENTS
- QUALITY OF FIT MEASURED BY χ^2 OF AVERAGE OF NN COMPARED TO DATA

 χ^2 OF BEST FIT

 χ^2 OF BEST FIT VS. AVERAGE χ^2



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 χ^2 OF BEST FIT VS. AVERAGE χ^2



- IF NO STOPPING IMPLEMENTED, χ^2 OF THE AVERAGE DECREASES AS A FUNCTION OF AVERAGE χ^2 OF REPLICAS
- At best fit, average χ^2 of replicas ~ 2 ; χ^2 of average to data ~ 1

THE DATA

NONSINGLET

FULL



STOPPING I

- EACH NEURAL NET IS FITTED TO A PSEUDODATA REPLICA BY MINIMIZING THE χ^2 TO SUBSET OF DATA (TRAINING SET)
- FIT STOPS WHEN THE χ^2 OF THE REMAINING DATA STARTS TO GROW (VALIDATION SET)



STOPPING FOR THE χ^2 OF ONE REPLICA (FULL FIT)

STOPPING II

AFTER STOPPING CRITERION IMPLEMENTED (NONSINGLET FIT)DISTRIBUTION OF χ^2 AT STOPPINGDISTRIBUTION OF TRAINING LENGTHS



- POISSONIAN DISTRIBUTION OF TRAINING LENGTHS
- BEST FIT $\chi^2 = 0.75$ (BCDMS: 0.75, NMC: 0.72): EXPT. ERRORS SOMEWHAT OVERESTIMATED?
STABILITY (NONSINGLET FIT)

CAN CHECK STABILITY BY COMPARING RESULTS IF THE WHOLE PROCEDURE IS REPEATED WITH A DIFFERENT SET OF REPLICAS

DEFINE R.M.S. DISTANCE $\langle d[q] \rangle = \sqrt{2}$

$$\left\langle \left\langle \frac{\left(\langle q_i \rangle_{(1)} - \langle q_i \rangle_{(2)} \right)^2}{\sigma^2 [q_i^{(1)}] + \sigma^2 [q_i^{(2)}]} \right\rangle_{dat} \right\rangle_{dat}$$

NOTE $\sigma \Rightarrow$ ERROR ON AVERAGE = (ERROR ON q_i)/ \sqrt{N} \Rightarrow TESTS BOTH ACCURACY OF CENTRAL VALUE & ERRORS

SELF-STABILITY: DIFFERENT SETS OF 100 REPLICAS

$\langle d\left[q\right] \rangle_{\mathrm{dat}}$	0.96
$\left\langle d\left[q ight] ight angle _{\mathrm{extra}}$	0.99
$\left\langle d\left[\sigma_{q}\right] \right\rangle_{\mathrm{dat}}$	0.88
$\left\langle d\left[\sigma_{q}\right] ight angle_{\mathrm{extra}}$	0.97

CHANGE OF ARCHITECTURE: 2-4-3-1 VS. 2-5-3-1

$\langle d[q] \rangle_{\text{dat}}$	0.9
$\left\langle d\left[q ight] ight angle _{\mathrm{extra}}$	0.9
$\langle d\left[\sigma_{q}\right] \rangle_{\mathrm{dat}}$	0.9
$\left \left\langle d\left[\sigma_q\right] \right\rangle_{\text{extra}} \right $	1.4

DISTANCE COMPUTED FOR 14 POINTS LINEARLY SPACED IN THE DATA REGION ($0.05 \le x \le 0.75$) & 14 POINTS LOG SPACED IN THE EXTRAPOLATION REGION ($10^{-3} \le x \le 10^{-2}$)

RESULTS:



THE NONSINGLET QUARK PDF $q^{\rm NS}(x,Q^2)$

LO, NLO & NNLO NLO vs. NNLO



- quality of fit (χ^2) same at LO, NLO, NNLO
- NLO & NNLO agree within one σ NNLO terms negligible within errors
- LO & NLO agree within three σ

NLO terms absorbed in b.c.

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NLO terms absorbed in b.c.

LOW $\alpha_s = 0.116$



VARIATION OF α_s



- quality of fit (χ^2) unchanged with $\alpha_s = 0.118 \pm 0.002$
- all fits agree within one σ $\Rightarrow \alpha_s$ cannot be determined with good accuracy

NEURAL INFORMATION HANDLING III INCOMPATIBLE DATA

- FOR PROTON FITS, CONVERGENCE ACHIEVED, BUT $E^{(0)}\gtrsim 1.4$ even w. very long training
- for NMC data $E^{(0)} \gtrsim 1.6$ (training with all data)
- for NMC data $E^{(0)} \gtrsim 2.2$ (training with NMC only)
- ALL OTHER STATISTICAL INDICATORS OK



NEURAL NET DISCARDS INCONSISTENT DATA & PROVIDES GOOD FIT TO THE REST