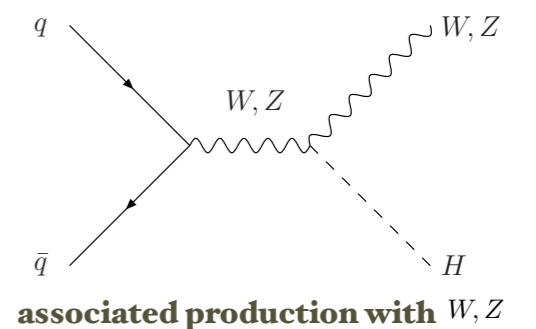
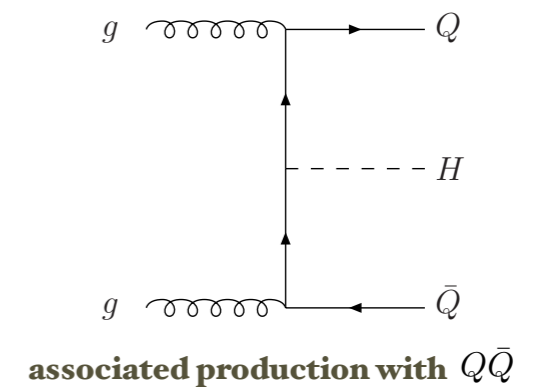
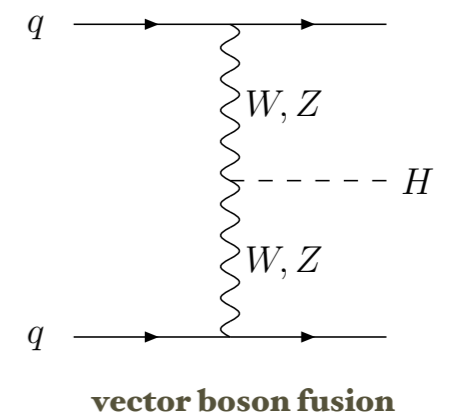
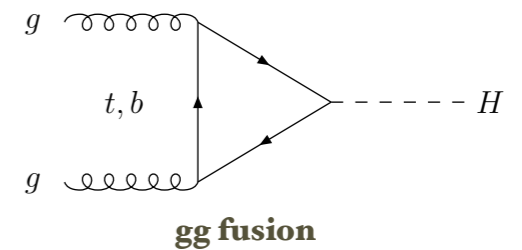
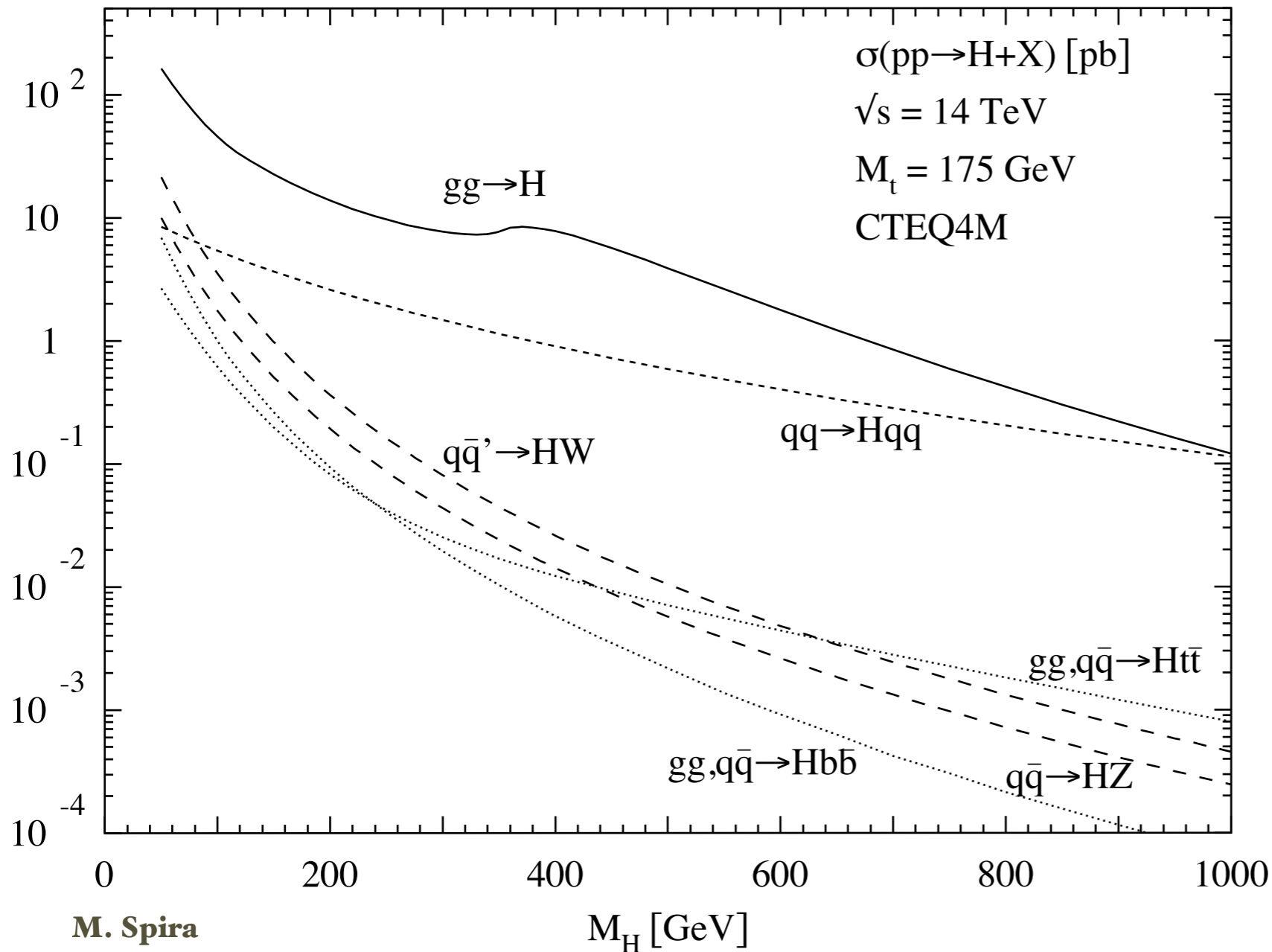


HNNLO: a MC program for Higgs boson production at hadron colliders

Massimiliano Grazzini (INFN, Firenze)

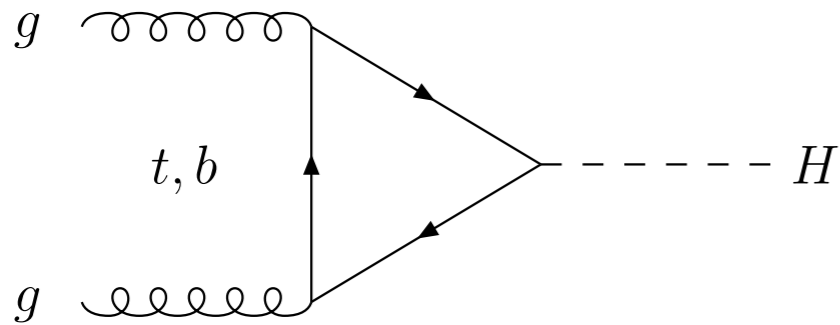
MCWS, february 18 2008

Higgs production at the LHC



Large gluon luminosity \longrightarrow gg fusion is the dominant production channel over the whole range of M_H

gg fusion



The Higgs coupling is proportional to the quark mass → top-loop dominates

NLO QCD corrections to the total rate computed more than 15 years ago and found to be large

A. Djouadi, D. Graudenz, M. Spira, P. Zerwas (1991)

They increase the LO result by about 80-100 % !

They are well approximated by the large- m_{top} limit

S.Dawson (1991)

(differences range from 1 to 4 % for $M_H < 200$ GeV) M.Kramer, E. Laenen, M.Spira(1998)

NNLO corrections to σ_H^{tot} computed in the large m_{top} limit

S. Catani, D. De Florian, MG (2001)

R.Harlander, W.B. Kilgore (2001,2002)

C. Anastasiou, K. Melnikov (2002)

V. Ravindran, J. Smith, W.L.Van Neerven (2003)

Effect ranges from 15 to 20 % for $M_H < 200$ GeV

Effects of soft-gluon resummation: additional +6 %

S. Catani, D. De Florian,

P. Nason, MG (2003)

Nicely confirmed by computation of soft terms at N³LO

S. Moch, A. Vogt (2006)

EW two-loop effects also known

U. Aglietti et al. (2004)

G. Degrandi, F. Maltoni (2004)

Up to now only total cross sections but...more exclusive observables are needed !

- $H + 1$ jet: NLO corrections known

D. de Florian, Z. Kunszt, MG (1999)
J. Campbell, K.Ellis (MCFM)

- $H + 2$ jet: NLO corrections recently computed

J. Campbell, K.Ellis, G. Zanderighi (2006)

→ background for VBF

All these predictions are obtained in the large- m_{top} limit

→ (it is a good approximation for small transverse momenta of the accompanying jets)

Del Duca et al. (2001)

NNLO corrections to $gg \rightarrow H$ computed for arbitrary cuts for $H \rightarrow \gamma\gamma$ → FEHIP

C. Anastasiou,
K. Melnikov, F. Petrello(2005)



It was the first fully exclusive NNLO calculation for a physically interesting process but....



If you are interested in distributions you need to do a single run for each bin
→ requires a lot of CPU time !

The optimal solution would be to have a *parton-level event generator*

With such a program one can apply arbitrary cuts and obtain the desired distributions in the form of bin histograms

→ this is what is typically done at NLO with the *subtraction method*

Quite an amount of work has been done in the last few years towards a general extension of the subtraction method to NNLO

D. Kosower (1998,2003,2005)

S. Weinzierl (2003)

S. Frixione, MG (2004)

A. & T. Gehrmann, N. Glover (2005)

G, Somogyi, Z. Trocsanyi, V. Del Duca
(2005, 2007)

Up to now results obtained for $e^+e^- \rightarrow 2$ jets

A. & T. Gehrmann, N. Glover (2004)

S. Weinzierl (2006)

and now for $e^+e^- \rightarrow 3$ jets

A. & T. Gehrmann, N. Glover, G. Heinrich (2007)

NEW:

HNNLO

S. Catani, MG (2007)

We propose a new version of the subtraction method to compute higher order QCD corrections to a specific class of processes in hadron collisions (vector boson, Higgs boson production, vector boson pairs.....)

We compute the NNLO corrections to $gg \rightarrow H$ implementing them in a fully exclusive parton level generator including all the relevant decay modes

 encompasses previous calculations in a single stand-alone numerical code
it makes possible to apply arbitrary cuts

Strategy: start from NLO calculation of H+jet(s) and observe that as soon as the transverse momentum of the Higgs $q_T \neq 0$ one can write:

$$d\sigma_{(N)NLO}^H|_{q_T \neq 0} = d\sigma_{(N)LO}^{H+jets}$$

Define a counterterm to deal with singular behaviour at $q_T \rightarrow 0$

But.....

the singular behaviour of $d\sigma_{(N)LO}^{H+\text{jet}(s)}$ is well known from the resummation program of large logarithmic contributions at small transverse momenta

G. Parisi, R. Petronzio (1979)

J. Collins, D.E. Soper, G. Sterman (1985)

S. Catani, D. de Florian, MG (2000)

→ choose $d\sigma^{CT} \sim d\sigma^{(LO)} \otimes \Sigma^H(q_T/Q)$

where

$$\Sigma^H(q_T/Q) \sim \sum_{n=1}^{\infty} \left(\frac{\alpha_S}{\pi}\right)^n \sum_{k=1}^{2n} \Sigma^{H(n;k)} \frac{Q^2}{q_T^2} \ln^{k-1} \frac{Q^2}{q_T^2}$$

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Then the calculation can be extended to include the $q_T = 0$ contribution:

$$d\sigma_{(N)NLO}^H = \mathcal{H}_{(N)NLO}^H \otimes d\sigma_{LO}^H + [d\sigma_{(N)LO}^{H+\text{jets}} - d\sigma_{(N)LO}^{CT}]$$

where I have subtracted the truncation of the counterterm at (N)LO and added a contribution at $q_T = 0$ to restore the correct normalization

HNNLO implements three decay channels:

- $H \rightarrow \gamma\gamma$ (higgsdec = 1)
 - $H \rightarrow WW \rightarrow l\nu l\nu$ (higgsdec = 2)
 - $H \rightarrow ZZ \rightarrow 4l$
 - $H \rightarrow e^+e^-\mu^+\mu^-$ (higgsdec = 31)
 - $H \rightarrow e^+e^-e^+e^-$ (higgsdec = 32)
- ➔ includes appropriate interference contribution

The user can choose the cuts and plot the required distributions by modifying the cuts.f and plotter.f subroutines

Results: $gg \rightarrow H \rightarrow \gamma\gamma$

S. Catani, MG (2007)

Use cuts as in CMS TDR

$$p_T^{\min} > 35 \text{ GeV}$$

$$p_T^{\max} > 40 \text{ GeV}$$

$$|y| < 2.5$$

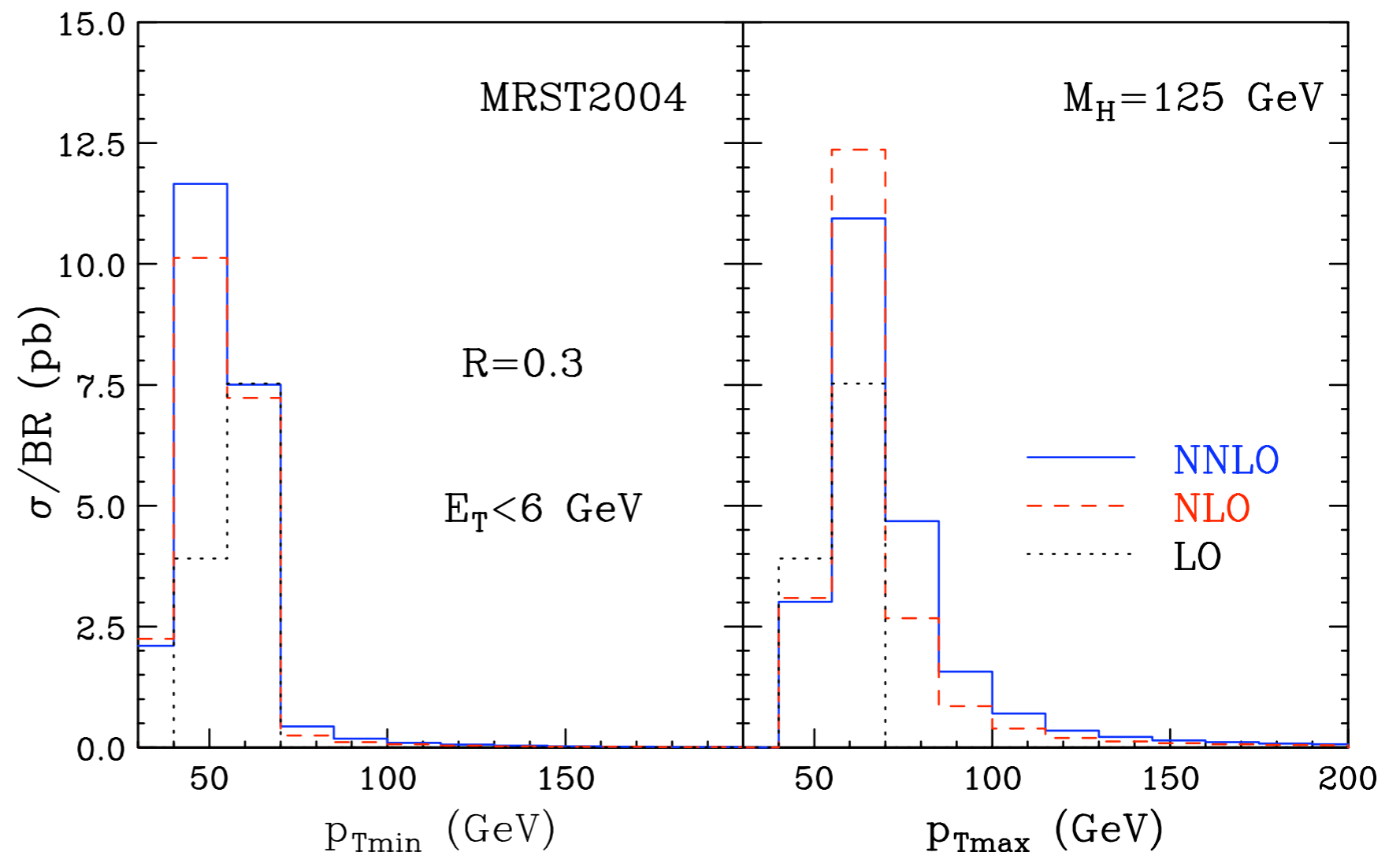
Photons should be isolated: total transverse energy in a cone of radius $R = 0.3$ should be smaller than 6 GeV

corresponding distributions

note perturbative instability when

$$p_T \rightarrow M_H/2$$

We find good agreement with FEHIP



Results: $gg \rightarrow H \rightarrow \gamma\gamma$

Use cuts as in CMS TDR

$$p_T^{\min} > 35 \text{ GeV}$$
$$p_T^{\max} > 40 \text{ GeV} \quad |y| < 2.5$$

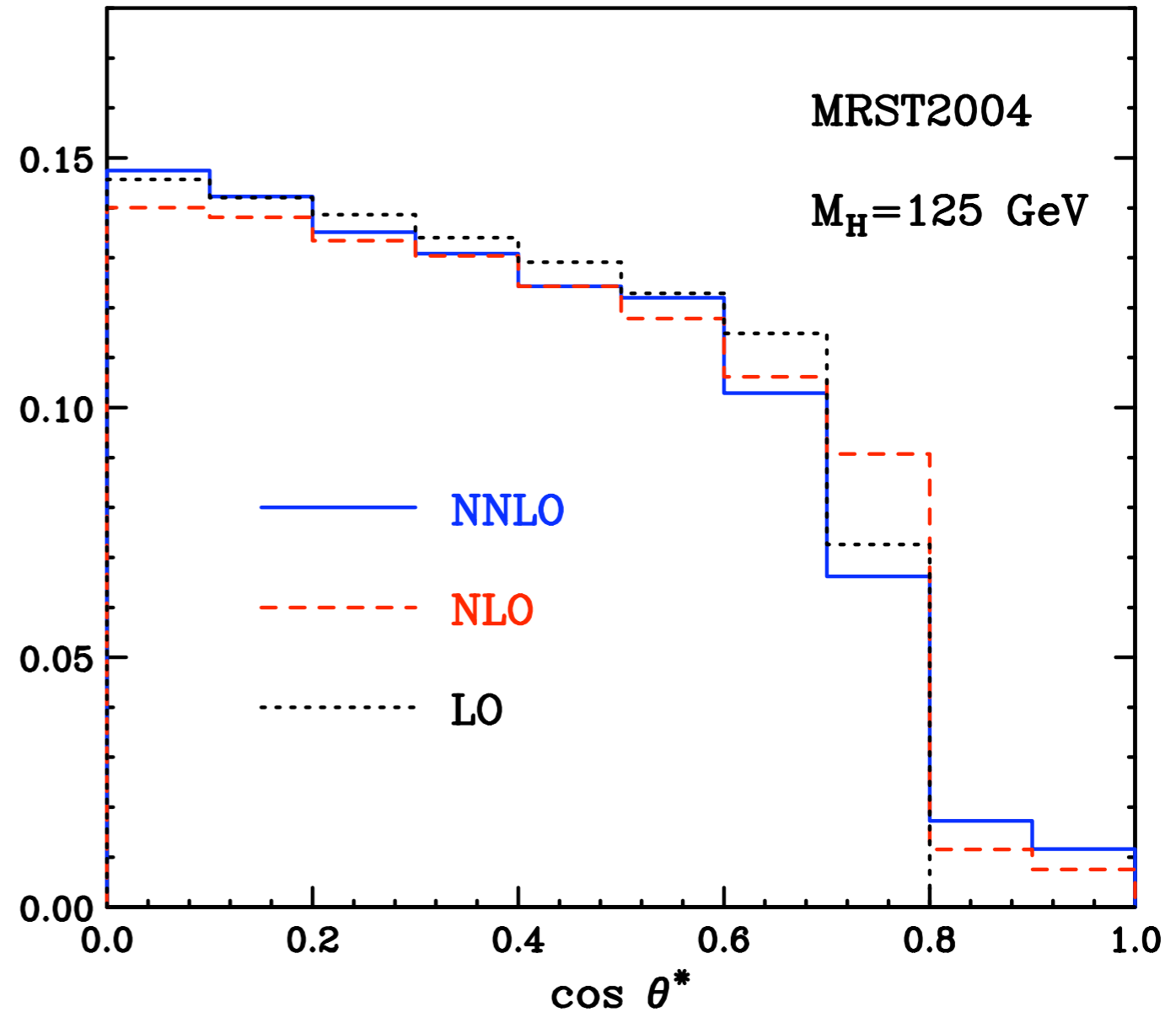
Photons should be isolated:
total transverse energy in a
cone of radius $R = 0.3$ should
be smaller than 6 GeV

define $\cos \theta^*$ distribution

θ^* **polar angle of one of the
photons in the Higgs rest frame
(used by ATLAS)**

note upper bound on $\cos \theta^*$ at LO

→ again perturbative instability
beyond LO !



Results: $gg \rightarrow H \rightarrow WW \rightarrow l\nu l\nu$

MG (2007)

Use *preselection cuts* as in Davatz. et al (2003)

see also C.Anastasiou, G. Dissertori, F. Stockli (2007)

$$p_T^l > 20 \text{ GeV}$$

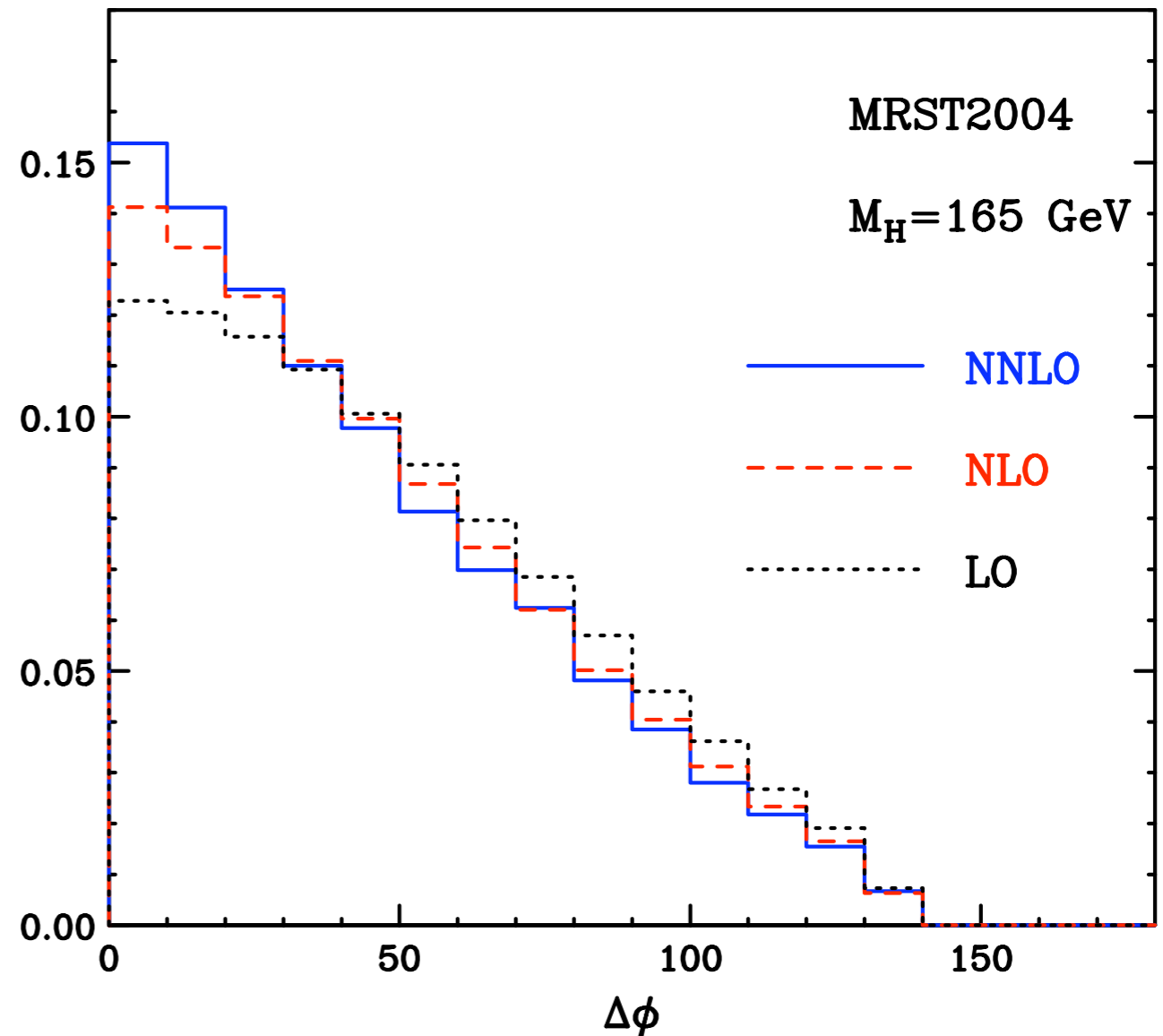
$$|y_l| < 2$$

$$p_T^{\text{miss}} > 20 \text{ GeV}$$

$$\Delta\phi < 135^\circ$$

$$m_{ll} < 80 \text{ GeV}$$

**normalized $\Delta\phi$
distribution**



The distributions appears to be steeper when going from LO to NLO and from NLO to NNLO

Use now *selection cuts* as in Davatz. et al (2003)

$$p_T^{\min} > 25 \text{ GeV} \quad m_{ll} < 35 \text{ GeV} \quad \Delta\phi < 45^\circ$$

$$35 \text{ GeV} < p_T^{\max} < 50 \text{ GeV} \quad |y_l| < 2 \quad p_T^{\text{miss}} > 20 \text{ GeV}$$

Results for

$$p_T^{\text{veto}} = 30 \text{ GeV}$$

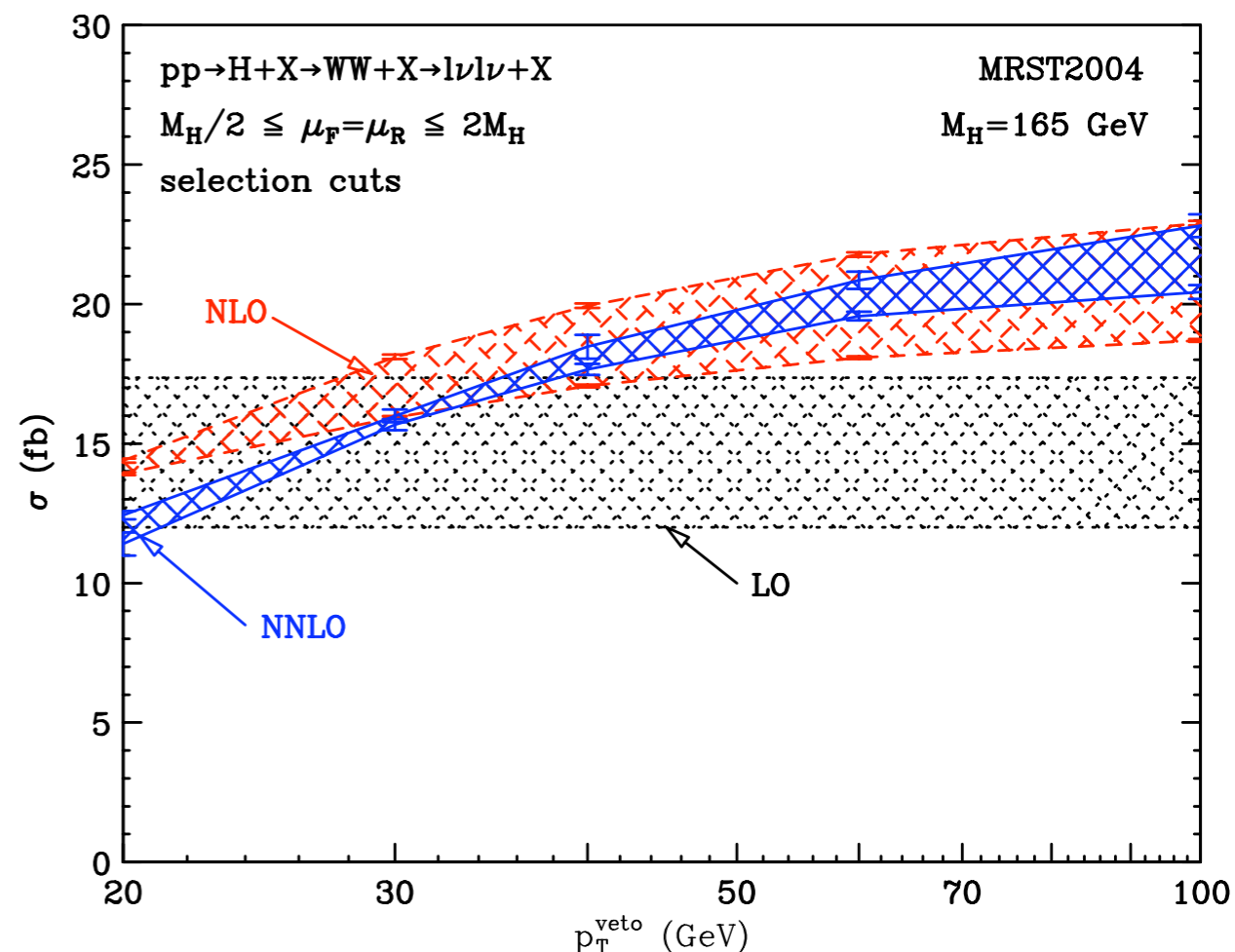
σ (fb)	LO	NLO	NNLO
$\mu_F = \mu_R = M_H/2$	17.36 ± 0.02	18.11 ± 0.08	15.70 ± 0.32
$\mu_F = \mu_R = M_H$	14.39 ± 0.02	17.07 ± 0.06	15.99 ± 0.23
$\mu_F = \mu_R = 2M_H$	12.00 ± 0.02	15.94 ± 0.05	15.68 ± 0.20

➔ **Impact of higher order corrections strongly reduced by selection cuts**

The NNLO band overlaps with the NLO one for $p_T^{\text{veto}} \gtrsim 30 \text{ GeV}$

The bands do not overlap for $p_T^{\text{veto}} \lesssim 30 \text{ GeV}$

NNLO efficiencies found in good agreement with MC@NLO



Results: $gg \rightarrow H \rightarrow ZZ \rightarrow e^+e^-e^+e^-$

MG (2007)

Inclusive cross sections:

σ (fb)	LO	NLO	NNLO
$\mu_F = \mu_R = M_H/2$	2.457 ± 0.001	4.387 ± 0.006	4.82 ± 0.03
$\mu_F = \mu_R = M_H$	2.000 ± 0.001	3.738 ± 0.004	4.52 ± 0.02
$\mu_F = \mu_R = 2M_H$	1.642 ± 0.001	3.227 ± 0.003	4.17 ± 0.01

$$K_{NLO} = 1.87$$

$$K_{NNLO} = 2.26$$

Consider the *selection cuts* as in the CMS TDR: $|y| < 2.5$

$$p_{T1} > 30 \text{ GeV} \quad p_{T2} > 25 \text{ GeV} \quad p_{T3} > 15 \text{ GeV} \quad p_{T4} > 7 \text{ GeV}$$

Isolation: total transverse energy in a cone of radius R around each lepton should fulfill $E_T < 0.05 p_T$

For each e^+e^- pair, find the closest (m_1) and next to closest (m_2) to m_Z

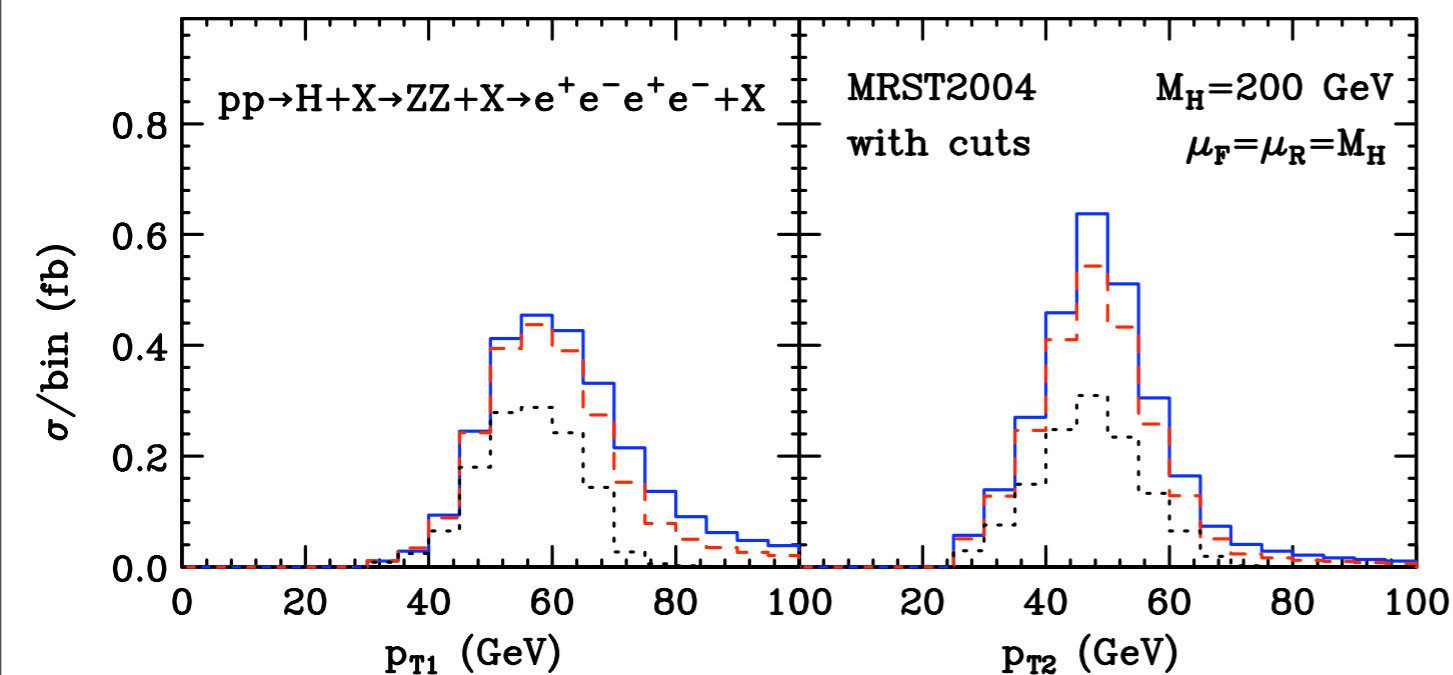
→ $81 \text{ GeV} < m_1 < 101 \text{ GeV}$ and $40 \text{ GeV} < m_2 < 110 \text{ GeV}$

The corresponding cross sections are:

σ (fb)	LO	NLO	NNLO
$\mu_F = \mu_R = M_H/2$	1.541 ± 0.002	2.764 ± 0.005	2.966 ± 0.023
$\mu_F = \mu_R = M_H$	1.264 ± 0.001	2.360 ± 0.003	2.805 ± 0.015
$\mu_F = \mu_R = 2M_H$	1.047 ± 0.001	2.044 ± 0.003	2.609 ± 0.010

$$K_{NLO} = 1.87$$

$$K_{NNLO} = 2.22$$

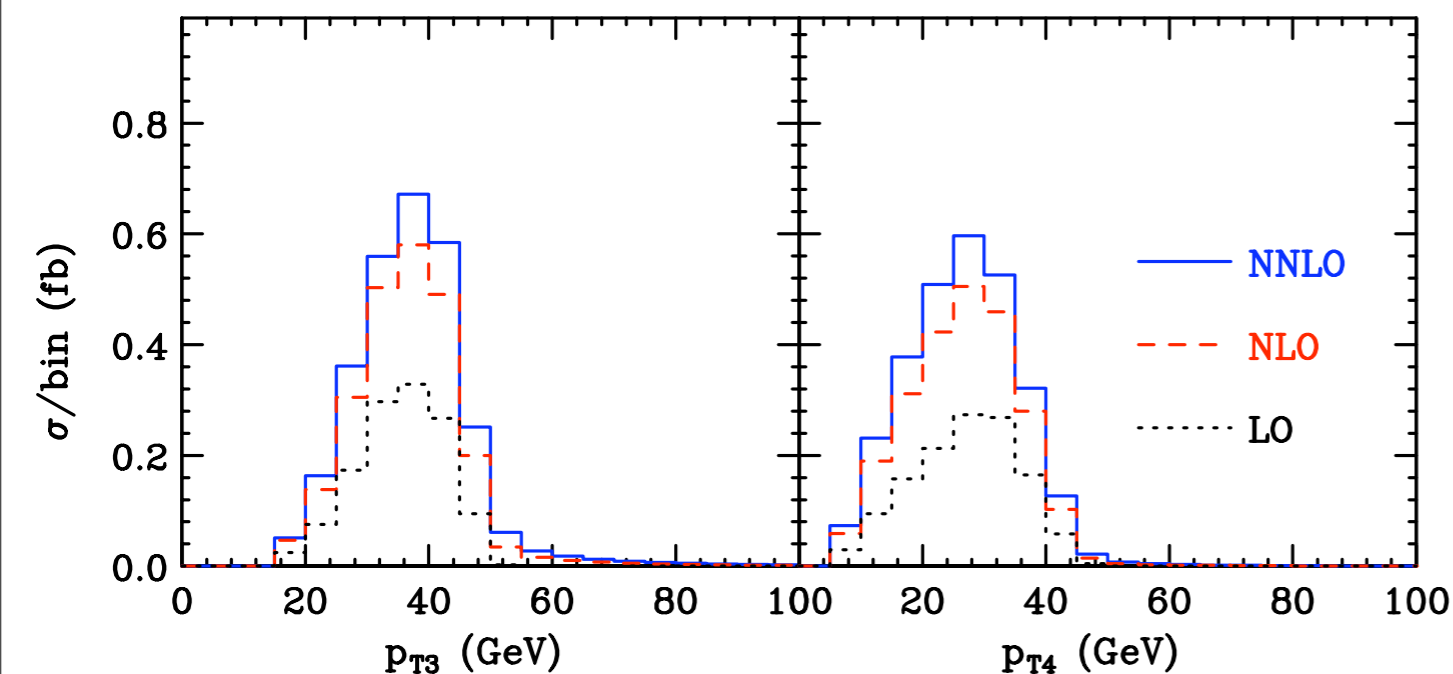


in this case the cuts are mild and do not change significantly the impact of higher order corrections

Note that at LO

$$p_{T1}, p_{T2} < M_H/2$$

$$p_{T3} < M_H/3 \quad p_{T4} < M_H/4$$



Behaviour at the kinematical boundary is smooth



**No instabilities
beyond LO**

Summary

- **HNNLO** is a numerical program to compute Higgs boson production through gluon fusion in pp or $p\bar{p}$ collisions at LO, NLO, NNLO

- It implements all the relevant decay modes of the Higgs boson:

$$H \rightarrow \gamma\gamma \quad H \rightarrow WW \rightarrow l\nu l\nu \quad H \rightarrow ZZ \rightarrow 4l$$

- It allows the user to apply arbitrary cuts on the final state photon/leptons and the associated jet activity, and to obtain the required distributions in the form of bin histograms
- These features should make our program a useful tool for Higgs studies at the Tevatron and the LHC
- Public version can be downloaded from

<http://theory.fi.infn.it/grazzini/codes.html>