

Stefano Frixione

Update on MC@NLO

MCWS, Frascati, 24/10/2006

SF & B. Webber, JHEP 0206(2002)029 [[hep-ph/0204244](#)]

SF, P. Nason & B. Webber, JHEP 0308(2003)007 [[hep-ph/0305252](#)]

SF, E. Laenen, P. Motylinski & B. Webber, JHEP 0603(2006)092 [[hep-ph/0512250](#)]

MC@NLO 3.2 [hep-ph/0601192]

IPROC	IV	IL ₁	IL ₂	Spin	Process
-1350-IL				✓	$H_1 H_2 \rightarrow (Z/\gamma^* \rightarrow) l_{\text{IL}} l_{\text{IL}} + X$
-1360-IL				✓	$H_1 H_2 \rightarrow (Z \rightarrow) l_{\text{IL}} l_{\text{IL}} + X$
-1370-IL				✓	$H_1 H_2 \rightarrow (\gamma^* \rightarrow) l_{\text{IL}} l_{\text{IL}} + X$
-1460-IL				✓	$H_1 H_2 \rightarrow (W^+ \rightarrow) l_{\text{IL}}^+ \nu_{\text{IL}} + X$
-1470-IL				✓	$H_1 H_2 \rightarrow (W^- \rightarrow) l_{\text{IL}}^- \bar{\nu}_{\text{IL}} + X$
-1396				×	$H_1 H_2 \rightarrow \gamma^*(\rightarrow \sum_i f_i f_i) + X$
-1397				×	$H_1 H_2 \rightarrow Z^0 + X$
-1497				×	$H_1 H_2 \rightarrow W^+ + X$
-1498				×	$H_1 H_2 \rightarrow W^- + X$
-1600-ID					$H_1 H_2 \rightarrow H^0 + X$
-1705					$H_1 H_2 \rightarrow b\bar{b} + X$
-1706				×	$H_1 H_2 \rightarrow t\bar{t} + X$
-2000-IC				×	$H_1 H_2 \rightarrow t/\bar{t} + X$
-2001-IC				×	$H_1 H_2 \rightarrow t + X$
-2004-IC				×	$H_1 H_2 \rightarrow t + X$
-2600-ID	1	7		×	$H_1 H_2 \rightarrow H^0 W^+ + X$
-2600-ID	1	<i>i</i>		✓	$H_1 H_2 \rightarrow H^0 (W^+ \rightarrow) l_i^+ \nu_i + X$
-2600-ID	-1	7		×	$H_1 H_2 \rightarrow H^0 W^- + X$
-2600-ID	-1	<i>i</i>		✓	$H_1 H_2 \rightarrow H^0 (W^- \rightarrow) l_i^- \bar{\nu}_i + X$
-2700-ID	0	7		×	$H_1 H_2 \rightarrow H^0 Z + X$
-2700-ID	0	<i>i</i>		✓	$H_1 H_2 \rightarrow H^0 (Z \rightarrow) l_i l_i + X$
-2850		7	7	×	$H_1 H_2 \rightarrow W^+ W^- + X$
-2850		<i>i</i>	<i>j</i>	✓	$H_1 H_2 \rightarrow (W^+ \rightarrow) l_i^+ \nu_i (W^- \rightarrow) l_j^- \bar{\nu}_j + X$
-2860		7	7	×	$H_1 H_2 \rightarrow Z^0 Z^0 + X$
-2870		7	7	×	$H_1 H_2 \rightarrow W^+ Z^0 + X$
-2880		7	7	×	$H_1 H_2 \rightarrow W^- Z^0 + X$

Recent activities:

- ▶ Spin correlations in $t\bar{t}$ and single-top production
- ▶ Wt channel for single-top production
- ▶ Improvements to Higgs production
- ▶ Interface to HERWIG++ (ISR only)
- ▶ Dijet production

<http://www.hep.phy.cam.ac.uk/theory/webber/MCatNLO>

Spin correlations: definitions

In the production process

$$a + b \longrightarrow P(\longrightarrow d_1 + \cdots + d_n) + X$$

there are

- ▶ *Decay* s.c.: if there is a non-trivial dependence on $(d_i \cdot d_j)$
- ▶ *Production* s.c.: if there is a non-trivial dependence on $(d_i \cdot a)$, $(d_i \cdot b)$, $(d_i \cdot X)$

MC@NLO **always implements decay s.c.** through HERWIG

Production s.c. are available in v3.2 for dilepton, H , WH , ZH , W^+W^- processes

Production s.c. are now also included in $t\bar{t}$ and **single-top** processes

Production spin correlations

The standard way: compute the matrix elements for

$$a + b \longrightarrow (P \longrightarrow) d_1 + \cdots + d_n + X \quad \text{Full ME}$$

This full-ME strategy is implemented in MC@NLO for:

- ▶ Single- V production ($V = W, Z, \gamma, Z/\gamma$)
- ▶ VH production ($V = W, Z$)

For large-multiplicity final states this may not be convenient, since

- ▶ ME must be integrated and unweighted
- ▶ The integration time increases and the unweighting efficiency decreases (for MC@NLO, typically $\varepsilon=10\text{--}40\%$) by increasing the number of final-state particles

For W^+W^- , $t\bar{t}$ and t production we have implemented an alternative strategy:

hit-and-miss

Hit-and-miss

Whatever the behaviours of the decay products, the momenta of the decaying particles will not change

⇒ The full ME's must be bounded from above by the undecayed ME's, times a **suitable constant**. Find this bound and do hit-and-miss

Advantages

- ▶ Only the undecayed ME's will be integrated: no further loss of time
- ▶ Unweighting is a **two-step** procedure: first get the P 's momenta, then the d 's momenta with hit-and-miss

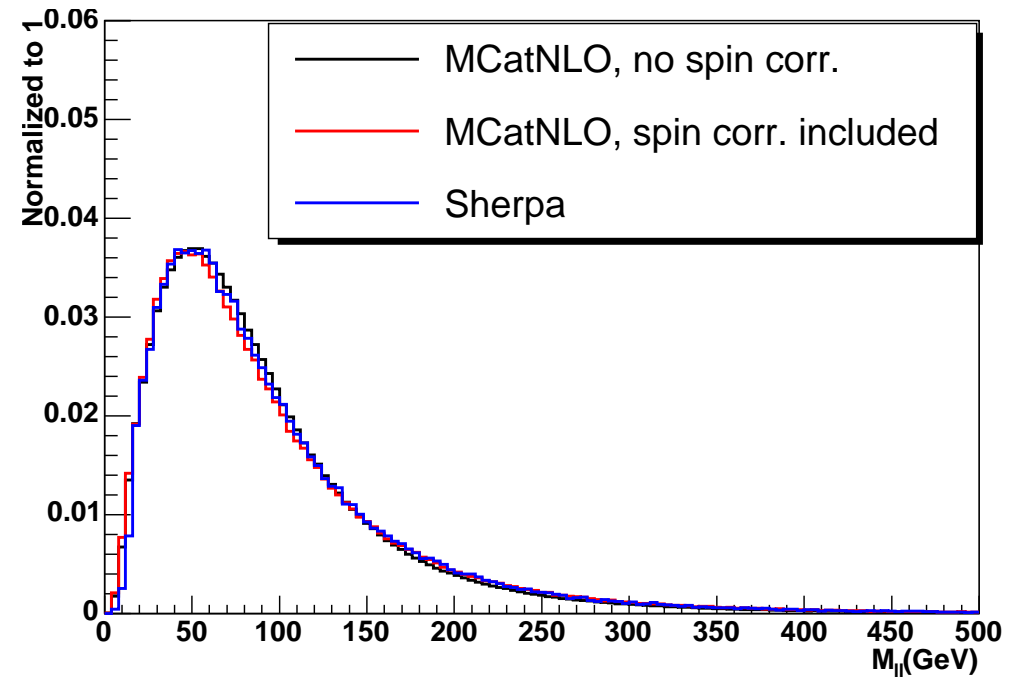
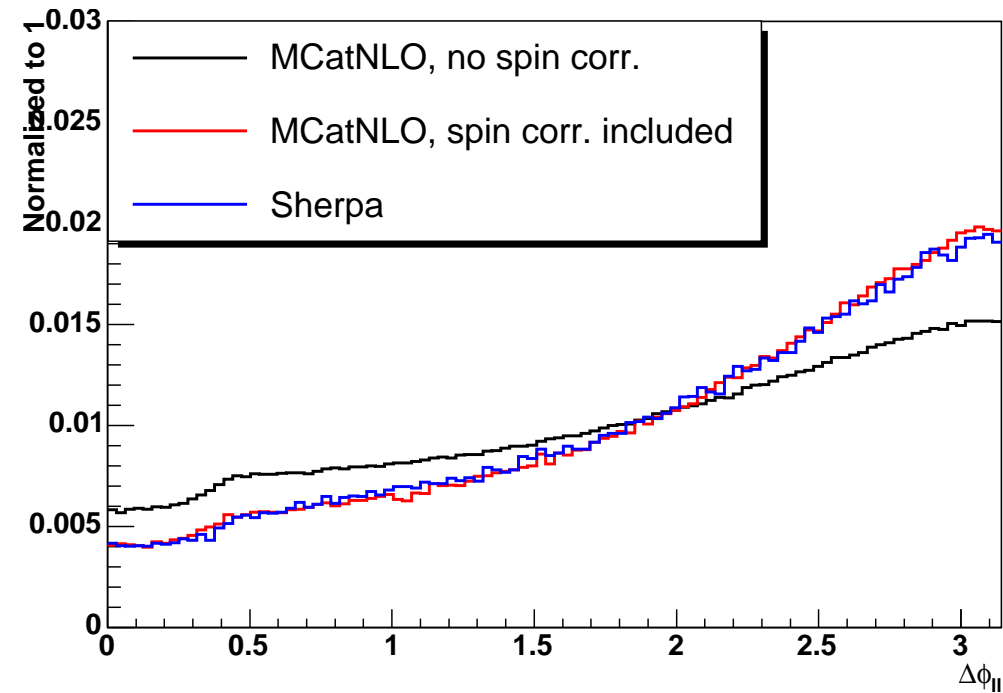
Vector bosons (tested and running)

$$\frac{d\sigma_{l_1\bar{l}_1\dots l_n\bar{l}_n}}{d\Phi_{2n+k}} \leq \left(\prod_{i=1}^n \frac{2 F_{V_i}^2 (V_{V_i l_i} + A_{V_i l_i})^2}{\Gamma_{V_i}^2} \right) \frac{d\sigma_{V_1\dots V_n}}{d\Phi_{n+k}}$$

Top (tests done – not yet released)

$$\frac{d\sigma_{b_1 l_1 \nu_1 \dots b_n l_n \nu_n}}{d\Phi_{3n+k}} \leq \left(\prod_{i=1}^n \frac{4g_W^4 |V_{tb}|^2 (k_{t_i} \cdot k_{l_i})(k_{b_i} \cdot k_{\nu_i})}{\left((q_i^2 - m_W^2)^2 + (m_W \Gamma_W)^2 \right) m_t^2 \Gamma_t^2} \right) \frac{d\sigma_{t_1 \dots t_n}}{d\Phi_{n+k}}$$

Results for W^+W^-

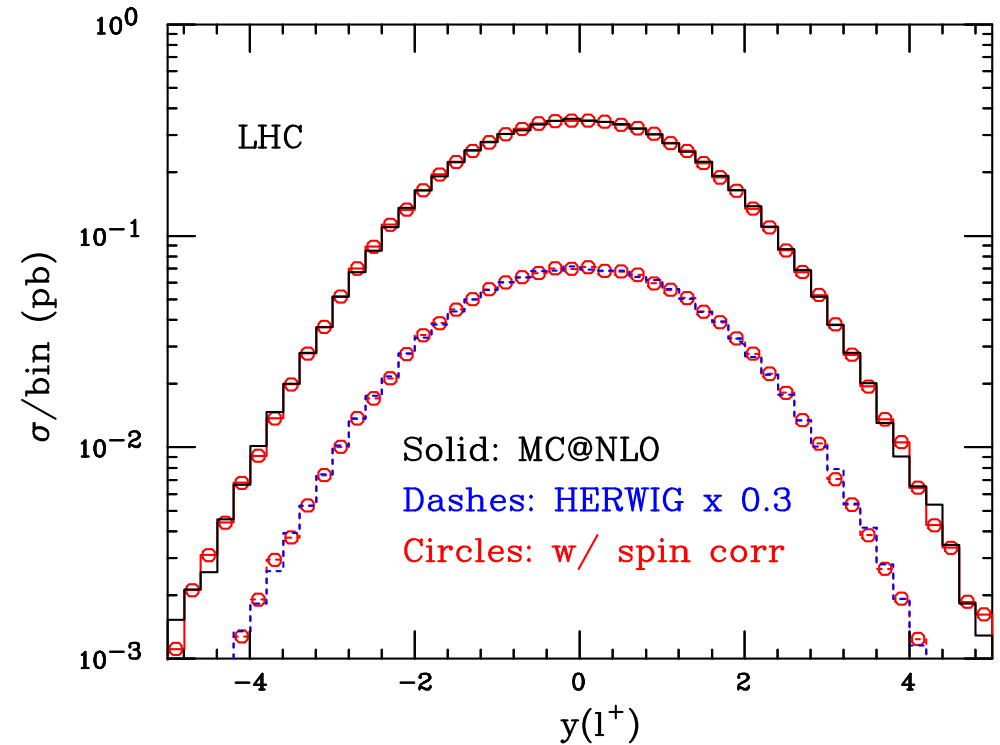
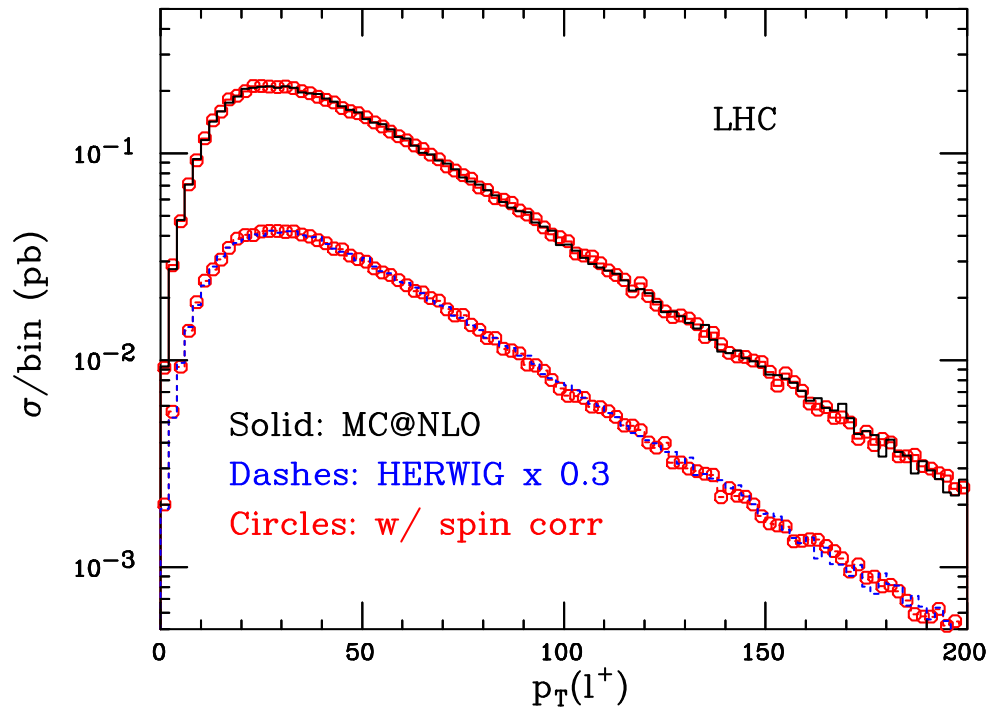


Plots: B. Quayle

- ▶ Virtual effects appear to be unimportant
- ▶ The effect of spin correlations is strictly dependent on the observable
- ▶ Released with v3.1

Thanks to Bill Quayle and Volker Drollinger for testing a preliminary version

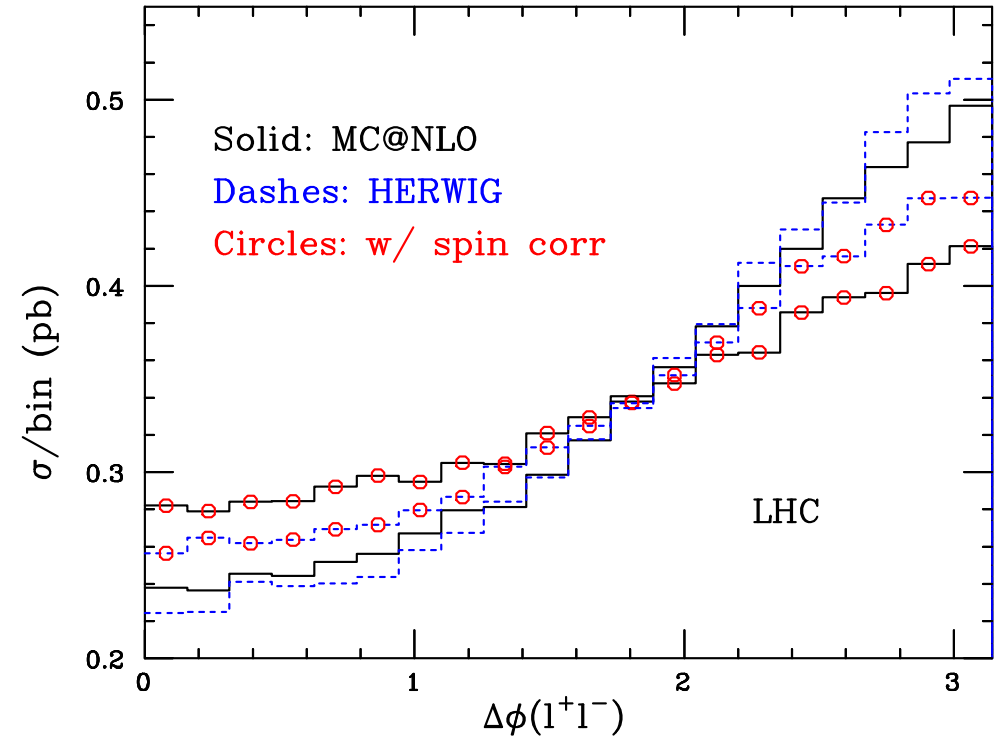
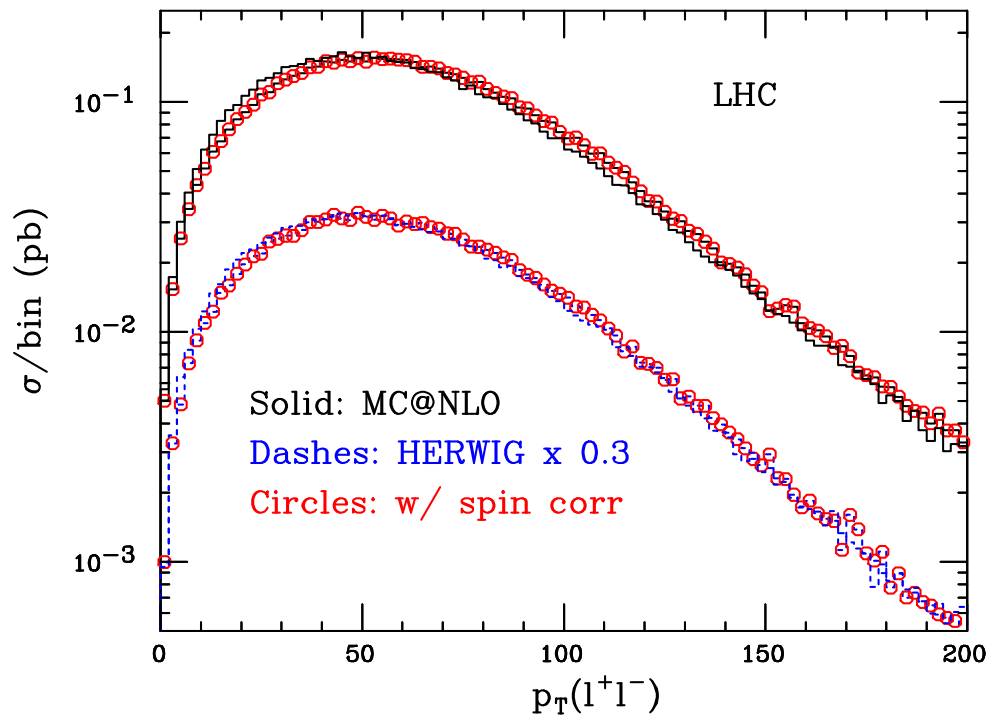
Spin correlations in $t\bar{t}$ I



- ▶ All single-inclusive distributions have this pattern
- ▶ Almost all correlations display a similar behaviour

“Large” is here comfortably small. Will this stay true after acceptance cuts?

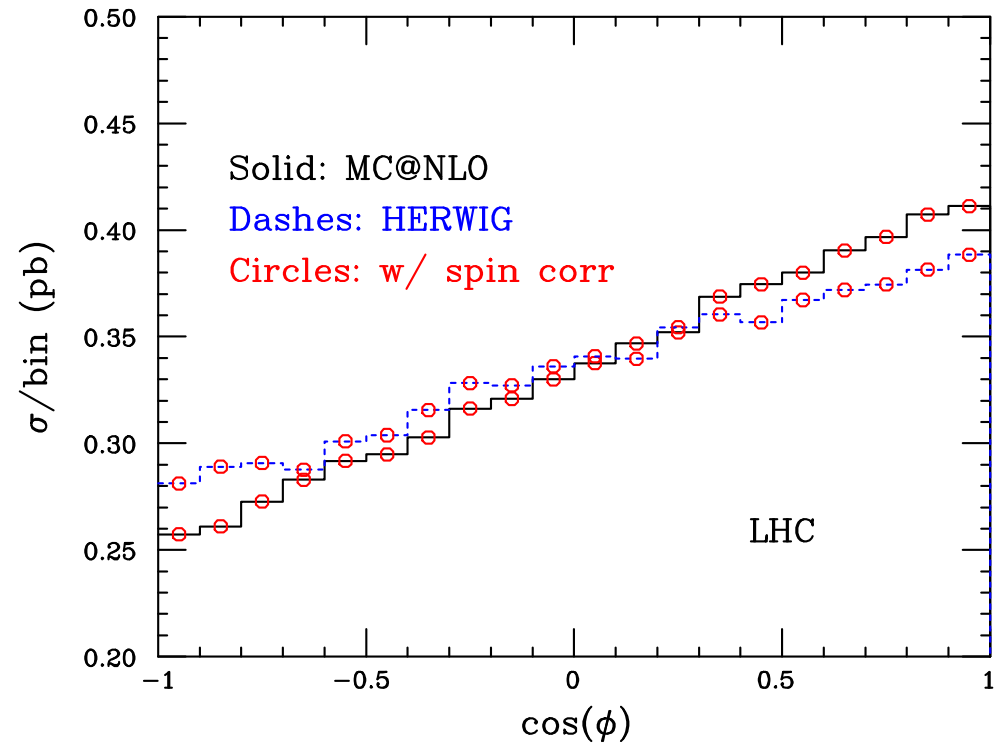
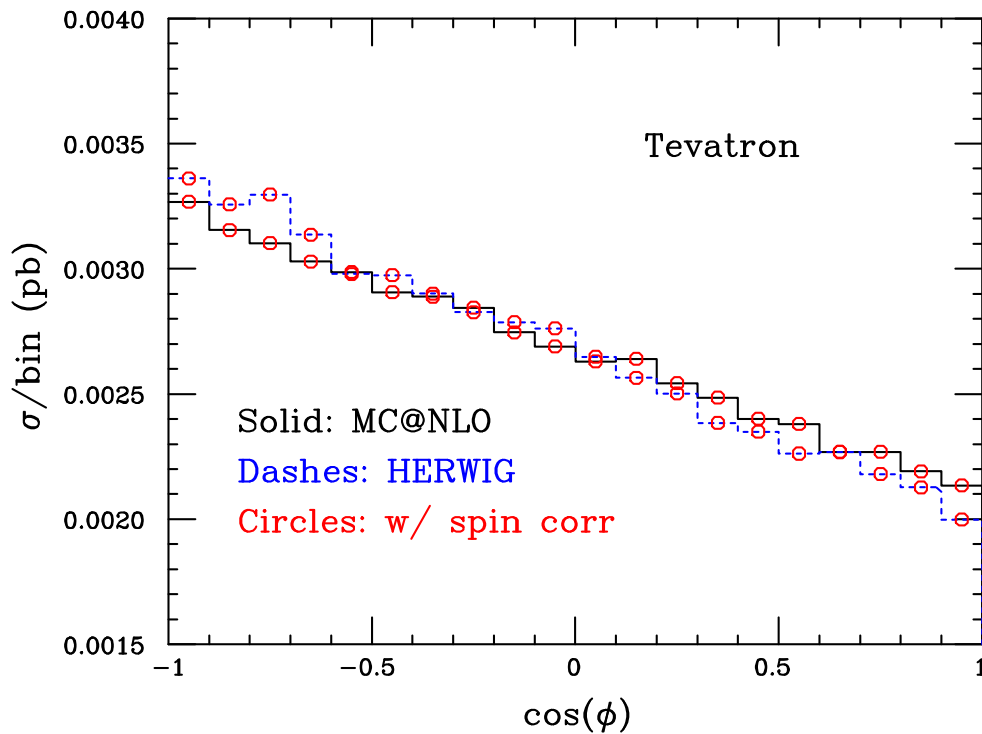
Spin correlations in $t\bar{t}$ II



- ▶ These are the only cases in which I've found non-negligible effects
- ▶ Spin correlations are not the whole story: for $\Delta\phi$, NLO effects are clearly visible

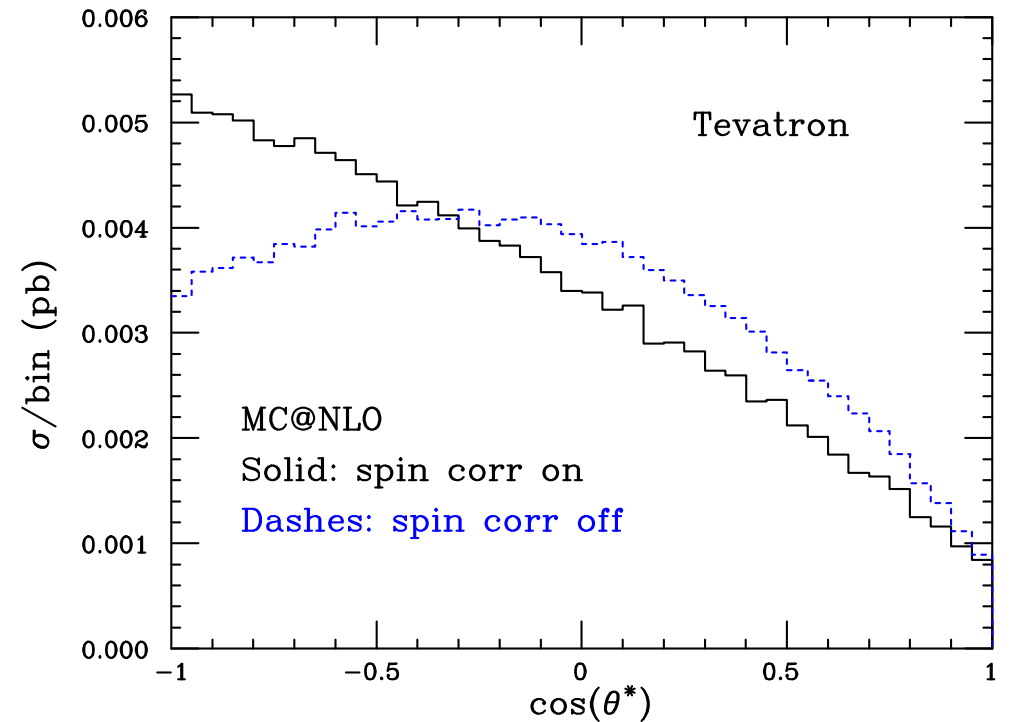
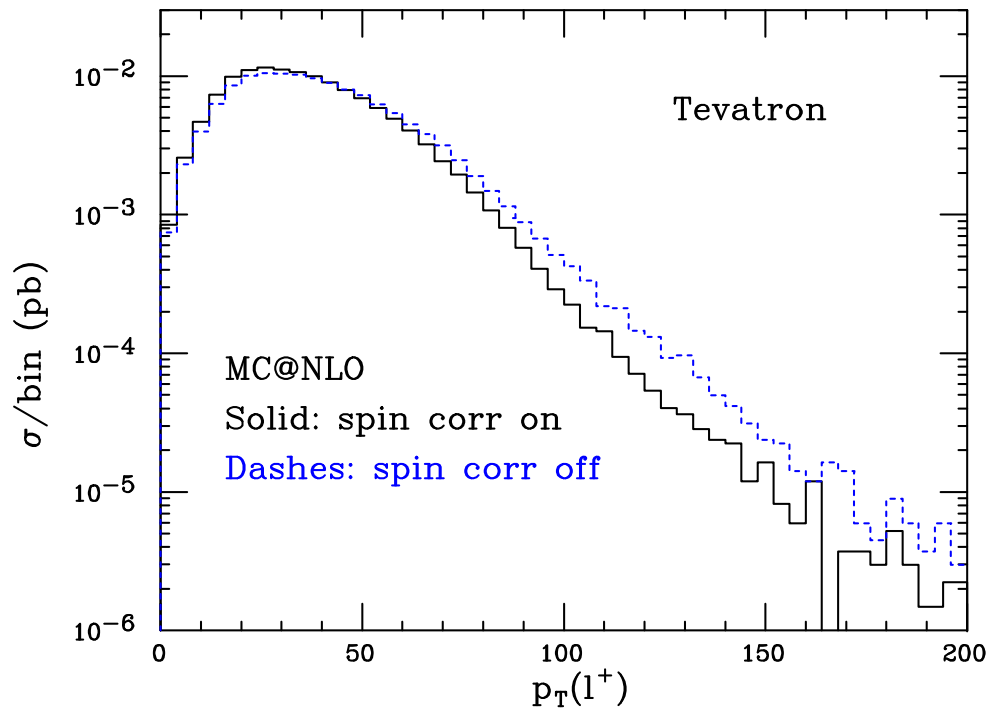
Observables can be designed to specifically target spin correlations \longrightarrow

Spin correlations in $t\bar{t}$ III



- ▶ NLO corrections again visible and in agreement with parton-level fixed-order results
- ▶ These kinds of observables are difficult to define in practice: need to know the rest frames of the $t\bar{t}$ system, of the t and of the \bar{t}

Spin correlations in single top



- ▶ For single-top, “large” is large indeed: the production proceeds through W exchange which effectively polarizes the top
- ▶ The effects are visible in single-inclusive observables (at variance with $t\bar{t}$)

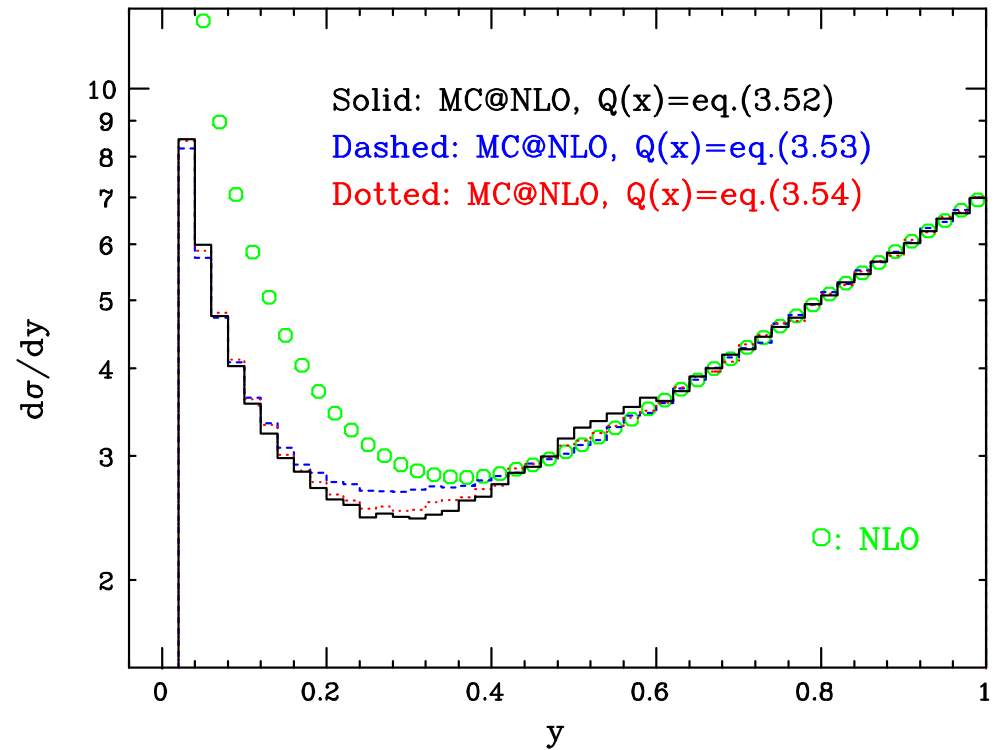
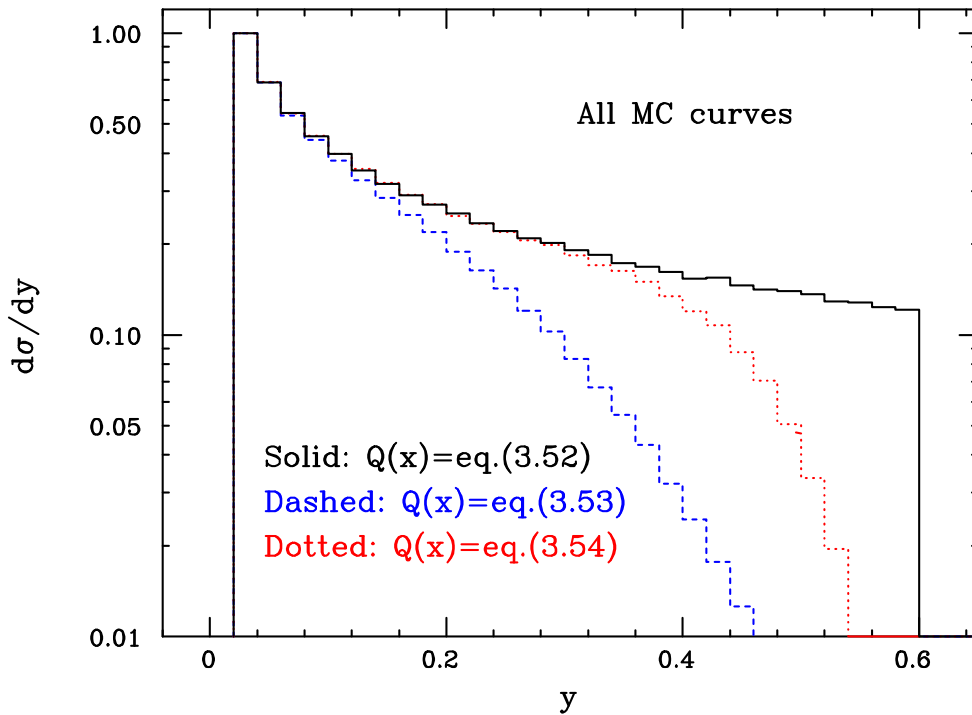
MC@NLO: reminder

■ MC@NLO generating functional (simplified notation)

$$\mathcal{F}_{\text{MC@NLO}} = \int_0^1 dx \left[\mathcal{F}_{\text{MC}}(S, x) \frac{\alpha_S [R(x) - BQ(x)]}{x} + \mathcal{F}_{\text{MC}}(S, 0) \left(B + \alpha_S V + \frac{\alpha_S B [Q(x) - 1]}{x} \right) \right]$$

- ▶ The form of $Q(x)$ is dictated by the parton shower MC@NLO is interfaced to. For HERWIG, it is a Θ function \longrightarrow dead zone
- ▶ We may, however, replace the Θ function in HERWIG with a smoother function, in order to reduce border effects. This can be done easily *without* modifying the code
- ▶ This also allows one to study matching ambiguity
- ▶ Never done in practice so far, border effects being invisible

Results (for a toy model)



1. $Q(x) = \Theta(x_{dead} - x);$

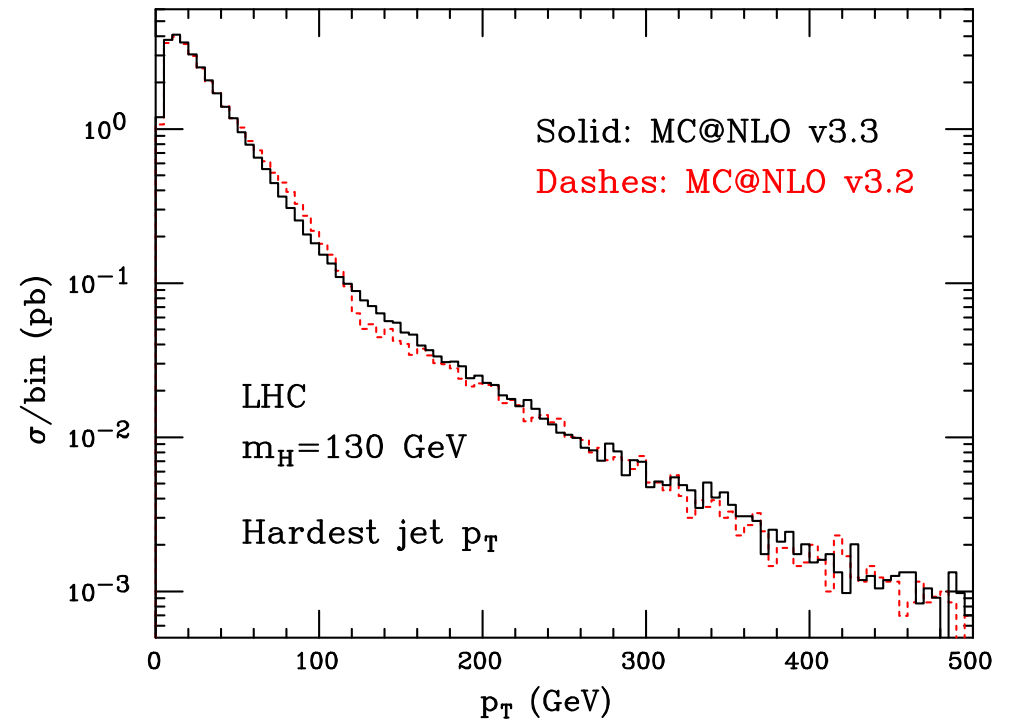
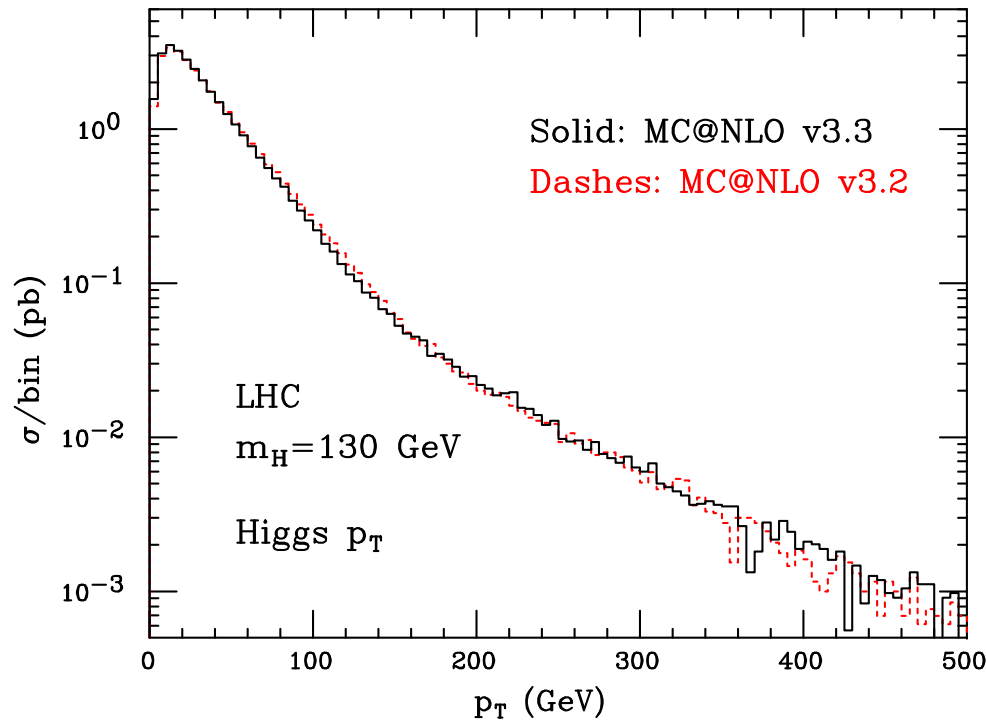
2. $Q(x) = \Theta(x_{dead} - x)G(x/x_{dead}),$ with $\alpha = 1, \beta = 1, c = 1;$

3. $Q(x) = \Theta(x_{dead} - x)G(x/x_{dead}),$ with $\alpha = 2, \beta = 1, c = 8.$

$$G(x) = \frac{c^2(1-x)^{2\beta}}{x^{2\alpha} + c^2(1-x)^{2\beta}}$$

■ Very smooth transition across the dead zone border (good control beyond NLO)

Border effects in Higgs production



- ▶ Pointed out by the Wisconsin group (ATLAS)
- ▶ Affects hardest-jet p_T
- ▶ New version stops HERWIG shower at $\alpha m_H \leq p_T \leq m_H/\sqrt{2}$, with p_T generated according to a probability function $\mathcal{P}(\alpha m_H) = 1$, $\mathcal{P}(m_H/\sqrt{2}) = 0$
- ▶ This also allowed us to change the scale of α_S in the MC counterterms \implies negative weights went down to 5% (were 8%)

On MC@NLO code

Time for the inclusion of a new process is spent:

- ◆ 80% for the pure-NLO computation
- ◆ 15% for MC counterterms and LHI-related code
- ◆ 5% debugging

The structure of the MC counterterms is modular

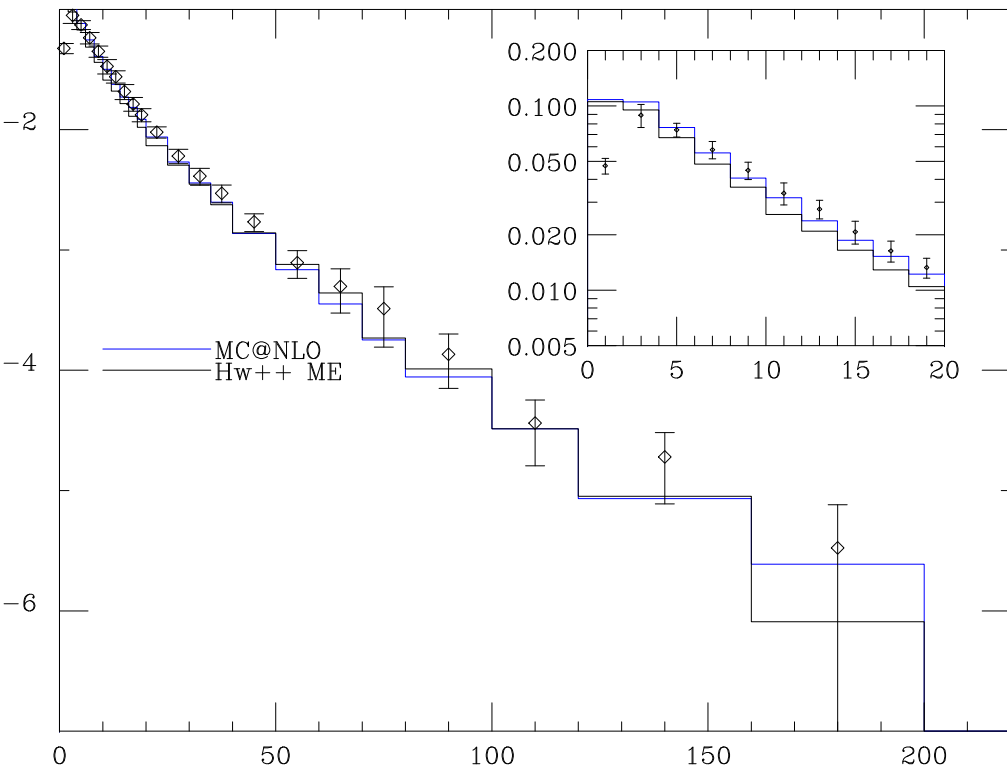
$$\mathcal{M}^{(\text{MC})} = \mathcal{K}^{(\text{MC})} \mathcal{M}^{(b)}$$

Kernels $\mathcal{K}^{(\text{MC})}$ now fully worked out for HERWIG

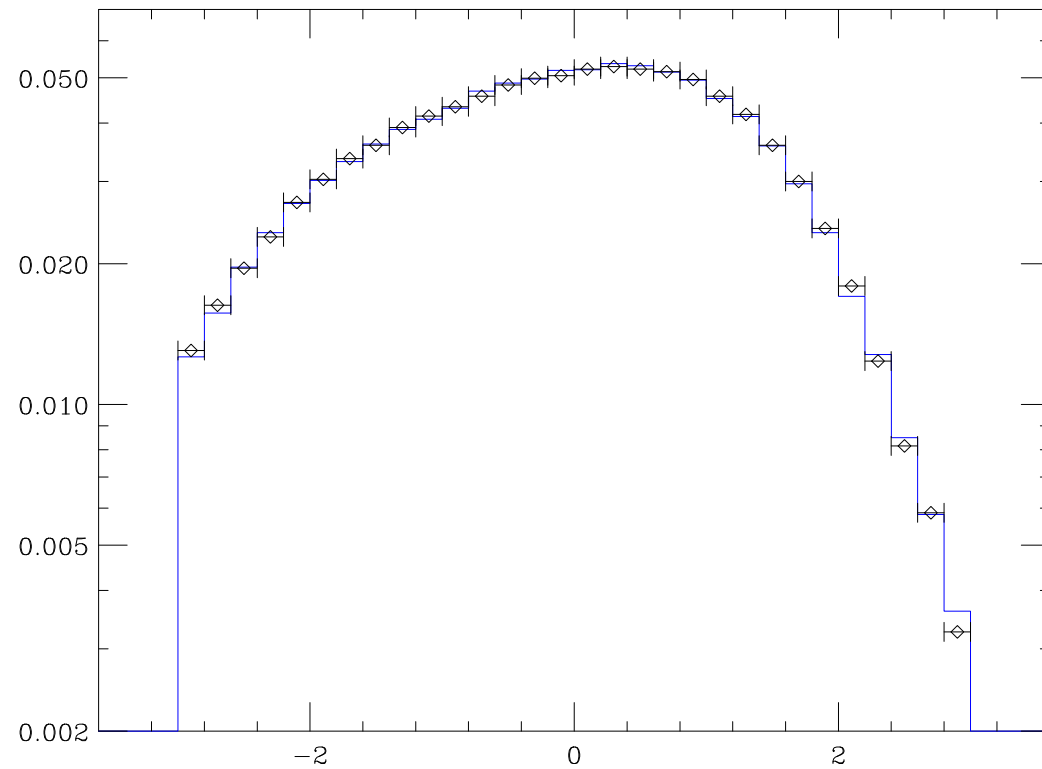
- Work in progress ([Seyi Latunde-Dada](#)) on the computation of $\mathcal{K}^{(\text{MC})}$ for HERWIG++. ISR now done

W production with MC@NLO/HERWIG++

pt of W from Run 1 data



Rapidity of e^+ from Run 2 data



- ▶ We know that old-style ME corrections distort the p_T spectrum
- ▶ We see that by adding the full NLO MEs one improves the agreement to data
- ▶ This is without specific tuning. Also, it is not yet known how to include a k_T -kick into HERWIG++ (affects lowest- p_T bins)

Conclusions

- ◆ The addition of spin correlations (to be officially released with the next version) adds interesting features in top physics – we are beginning to study phenomenology implications
- ◆ We have seen in the case of Higgs production that by limiting MC radiation one has beneficial effects. Presumably will have an impact on jet shapes in $t\bar{t}$ production (to be tested soon)
- ◆ This is a very well known technique in matched computations based on analytical resummations
- ◆ HERWIG++ has started producing results. More will follow